

Derivative of Unitary is not always $-iHU$

A S Muddu¹

¹Reliance Corporate Park, Reliance Industries Limited, Mumbai 400701, India

December 23, 2020

Abstract

$i\dot{U} = HU$ (or $\dot{U} = -iHU$) is the equation that is said to govern the evolution of a unitary matrix U given the Hamiltonian H of the system. This equation is said to hold true even if the Hamiltonian is time dependent. We show $iU = HU$ (or $\dot{U} = -iHU$) is the equation that is said to govern the evolution of a unitary matrix U given the Hamiltonian H of the system. This equation is said to hold true even if the Hamiltonian is time dependent. We show in this paper that $i\dot{U} = HU$ may not always hold for time dependent Hamiltonians. this paper that $i\dot{U} = HU$ may not always hold for time dependent Hamiltonians.

Introduction

The evolution of a unitary operator is given by $i\dot{U} = HU$. (where H is the Hamiltonian of the system). This equation is so common place that no one bothers to cite it. It is considered as a standard part of the curricula of certain courses, in physics books (see [1]) and even in allied engineering fields. For example see [2] (a course in Nuclear Engineering). Needless to say, the formula also has been extensively used in the literature. For example see Equation 35 in [3], Equation 25 in [4], Equation 7 in [5], Equation 4 in [6].

The organisation of next section is as follows:

In subsection A we show that for any matrix A , $[A, \dot{A}]$ is not necessarily equal to 0 . Subsection B deals with the derivative of e^A . From this we find the derivative of a unitary matrix in subsection C. We show in subsection D as to why in case of time independent Hamiltonian $i\dot{U} = HU$ more or less holds good.

Finally in section III, we synthesize the various lines of argument into a concluding paragraph.

Proof

A and its time derivative don't commute

The title can be mathematically paraphrased as

$[A, \dot{A}] \neq \mathbf{0}$. For any function $f(A)$ (matrix or scalar valued) $\dot{f} = \frac{\partial f}{\partial t}$. Say we have a matrix valued function

$$A(t) = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \quad (1)$$

where a, b, c, d are four distinct functions of time. We can differentiate $A(t)$ as follows [7], [8]

$$\dot{A} = \begin{bmatrix} \dot{a} & \dot{b} \\ \dot{c} & \dot{d} \end{bmatrix} \quad (2)$$

Thus,

$$\begin{aligned} [A, \dot{A}] &= A\dot{A} - \dot{A}A \\ &= \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} \dot{a} & \dot{b} \\ \dot{c} & \dot{d} \end{bmatrix} - \begin{bmatrix} \dot{a} & \dot{b} \\ \dot{c} & \dot{d} \end{bmatrix} \begin{bmatrix} a & b \\ c & d \end{bmatrix} \\ &= \begin{bmatrix} b\dot{c} - \dot{b}c & a\dot{b} - \dot{a}b + b\dot{d} - \dot{b}d \\ c\dot{a} - \dot{c}a + d\dot{c} - \dot{d}c & - (b\dot{c} - \dot{b}c) \end{bmatrix} \end{aligned} \quad (3)$$

Any of the elements of $[A, \dot{A}]$ are not necessarily zero at all times. Thus $[A, \dot{A}]$ is not necessarily zero.

Special cases where $[A, \dot{A}] = \mathbf{0}$

1. A is a constant function. This makes $\dot{A} = \mathbf{0}$. Because of this $[A, \dot{A}] = \mathbf{0}$
2. $A = \lambda(t)I \Rightarrow \dot{A} = \dot{\lambda}(t)I$

$$\begin{aligned} \therefore [A, \dot{A}] &= A\dot{A} - \dot{A}A \\ &= (\lambda I)(\dot{\lambda}I) - (\dot{\lambda}I)(\lambda I) \\ &= (\lambda\dot{\lambda}I) - (\dot{\lambda}\lambda I) \\ &= \mathbf{0} \end{aligned} \quad (4)$$

3. A is a diagonal matrix. This implies that $b, c = 0$. From Equation 3 we have $[A, \dot{A}] = \mathbf{0}$.

4. $A = tB$, where B is a constant matrix.

$$\begin{aligned} \therefore [A, \dot{A}] &= A\dot{A} - \dot{A}A \\ &= tBB - BtB \\ &= \mathbf{0} \end{aligned} \quad (5)$$

It can be easily shown that even if A is a skew Hermitian $[A, \dot{A}]$ is not necessarily zero. For the rest of discussion we only consider A such that $[A, \dot{A}] \neq \mathbf{0} \forall t$

In the case $[A, \dot{A}] = \mathbf{0}$, everything is true and wonderful.

Derivative of matrix exponential functions

Let us see what we get on differentiating A^2 . From the product rule, we have

$$(\dot{A}^2) = A\dot{A} + \dot{A}A \quad (6)$$

But, given the previous discussion

$$\begin{aligned} (\dot{A}^2) &\neq \dot{A}A + \dot{A}A \quad (\because [A, \dot{A}] \neq 0) \\ &\neq 2\dot{A}A \end{aligned} \quad (7)$$

Similarly,

$$(\dot{A}^3) \neq 3\dot{A}A \quad (8)$$

Instead, we obtain

$$(\dot{A}^3) = A\dot{A}A + A\dot{A}A + \dot{A}AA \quad (9)$$

For A^n , we get

$$(\dot{A}^n) = (A\dot{A}^{n-1}) \quad (10)$$

$$= A(\dot{A}^{n-1}) + \dot{A}A^{n-1} \quad (11)$$

Continuing down, we are left with

$$(\dot{A}^n) = \sum_{m=0}^{n-1} A^m \dot{A}A^{n-(m+1)} \quad (12)$$

We know that

$$e^A = \sum_{n=0}^{\infty} \frac{A^n}{n!} \quad (13)$$

On applying the above train of thought

$$(e^{\dot{A}}) = \sum_{n=0}^{\infty} \sum_{m=0}^{n-1} \frac{A^m \dot{A}A^{n-(m+1)}}{n!} \quad (14)$$

[9] and [10] have rewritten the above formula in a more pleasing format. The simplified form of Equation 2.1 from [10] in terms of our notation is as follows

$$(e^{\dot{A}}) = \int_0^1 e^{A(1-s)} \dot{A} e^{As} ds \quad (15)$$

Let us try to derive Equation 15 from Equation 14. The steps below are inspired from [9].

When we differentiate A^n we obtain a series in which each term is a permutation of a product of $n-1$ A 's and one \dot{A} . So the series in Equation 14 can be rewritten as:

$$(e^{\dot{A}}) = \sum_{p=0}^{\infty} \sum_{m=0}^{\infty} \frac{A^m \dot{A} A^p}{(m+p+1)!} \quad (\because n = m + p + 1) \quad (16)$$

$$= \sum_{p=0}^{\infty} \sum_{m=0}^{\infty} \frac{A^m \dot{A} A^p}{m! p!} \frac{m! p!}{(m+p+1)!} \quad (17)$$

We know that

$$\int_0^1 (1-s)^m s^p ds = \frac{m! p!}{(m+p+1)!} \quad (18)$$

Hence:

$$(e^{\dot{A}}) = \sum_{p=0}^{\infty} \sum_{m=0}^{\infty} \frac{A^m \dot{A} A^p}{m! p!} \int_0^1 (1-s)^m s^p ds \quad (19)$$

$$(e^{\dot{A}}) = \int_0^1 \sum_{p=0}^{\infty} \sum_{m=0}^{\infty} \frac{(1-s)^m A^m \dot{A} A^p s^p}{m! p!} ds \quad (20)$$

$$(e^{\dot{A}}) = \int_0^1 \sum_{p=0}^{\infty} \left(\sum_{m=0}^{\infty} \frac{(1-s)^m A^m}{m!} \right) \dot{A} \frac{A^p s^p}{p!} ds \quad (21)$$

$$(e^{\dot{A}}) = \int_0^1 \sum_{p=0}^{\infty} e^{(1-s)A} \dot{A} \frac{A^p s^p}{p!} ds \quad (22)$$

$$(e^{\dot{A}}) = \int_0^1 e^{(1-s)A} \dot{A} e^{As} ds \quad (23)$$

Since, $[A, \dot{A}] \neq \mathbf{0}$, $[e^A, \dot{A}] \neq \mathbf{0}$. Given that $[e^A, \dot{A}] \neq \mathbf{0}$, we can say that

$$\begin{aligned} e^{A(1-s)} \dot{A} e^{As} &\neq \dot{A} e^{A(1-s)} e^{As} \\ &\neq \dot{A} e^A \end{aligned} \quad (24)$$

Hence,

$$\begin{aligned} (e^{\dot{A}}) &\neq \int_0^1 \dot{A} e^A ds \\ (e^{\dot{A}}) &\neq \dot{A} e^A \end{aligned} \quad (25)$$

From Equation 15, 24 and 25 one can safely say that $(e^{\dot{A}})$ *can not be* written as Be^A (where $B = g(\dot{A})$ i.e B is a function of only \dot{A}).

Derivative of a unitary matrix

We can write $U = e^{\Omega(t)}$, where $\Omega(t)$ is a matrix valued function of time. By the virtue of its construction, $\Omega(t)$ is an skew Hermitian matrix. Thus from end of section , $[\Omega, \dot{\Omega}] \neq \mathbf{0}$

Since $(\dot{e}^A) \neq B(\dot{A})e^A$ as proved in end of section

$$\begin{aligned}\dot{U} &\neq B(\dot{\Omega})e^{\Omega} \quad \text{or} \\ \dot{U} &\neq B(\dot{\Omega})U\end{aligned}\tag{26}$$

In words it means that \dot{U} cannot be written as a product of function of $B(\dot{\Omega})$ and U . We know that for time independent Hamiltonians $U = e^{-iHt}$. So it is fair to assume that Ω may depend on the Hamiltonian H in some way. Since Ω depends on the Hamiltonian H , $\dot{\Omega}$ too is a function of H .

\therefore From the previous Equation 26

$$\dot{U} \neq B'(H)U\tag{27}$$

where B' is another function only of H such that $B'(H) = B(\dot{\Omega})$.

Taking things a step further $\dot{U} \neq -iHU$ for time dependent Hamiltonians.

Time independent Hamiltonians

The rule $\dot{U} = -iHU$ still holds good for time independent Hamiltonians, but here too things are not the same as before. From special case 1 from subsection we have

$$[H, \dot{H}] = 0 \quad \text{since } H \text{ is time independent}\tag{28}$$

From the discussion in sub-section we can say that

$$[H, e^{-iHt}] = 0\tag{29}$$

$$[H, U] = 0\tag{30}$$

Thus $\dot{U} = -iHU$ can be transformed to $\dot{U} = -iUH$. So The order of U, H does not really matter on the right hand side the equation $\dot{U} = -iHU$ for time independent Hamiltonians.

Conclusion

In this paper, we have shown that \dot{U} is not always equal to $-iHU$ for time dependant Hamiltonians. This does not mean that it is not possible. One of the ways it may be possible is that the functions a, b, c, d of A align themselves in such a way that $[A, \dot{A}] = 0$ (other than those special cases considered in section). Under *more severely* restrictive conditions than those considered here, [12], [13] have shown that $(\dot{e}^A) = \dot{A}e^A$, even if $[A, \dot{A}] \neq 0$. But these restrictions coupled with the Hermiticity requirements of the Hamiltonian make this very unlikely to happen.

References

- [1] Nouredine Zettili. *Quantum mechanics : concepts and applications*. Wiley, Chichester, U.K, 2009.

- [2] Paola Cappellaro. Quantum Theory of Radiation Interactions — Nuclear Science and Engineering — MIT OpenCourseWare. <https://ocw.mit.edu/courses/nuclear-engineering/>, fall 2012. (Accessed on 06/25/2020).
- [3] Raj Chakrabarti and Herschel Rabitz. Quantum control landscapes. *International Reviews in Physical Chemistry*, 26(4):671–735, 2007.
- [4] Navin Khaneja, Timo Reiss, Cindie Kehlet, Thomas Schulte-Herbrüggen, and Steffen J Glaser. Optimal control of coupled spin dynamics: design of NMR pulse sequences by gradient ascent algorithms. *Journal of magnetic resonance*, 172(2):296–305, 2005.
- [5] Shai Machnes, Elie Assémat, David J Tannor, and Frank K Wilhelm. Gradient optimization of analytic controls: the route to high accuracy quantum optimal control. *arXiv preprint arXiv:1507.04261*, 2015.
- [6] T Schulte-Herbrüggen, A Spörl, N Khaneja, and SJ Glaser. Optimal control for generating quantum gates in open dissipative systems. *Journal of Physics B: Atomic, Molecular and Optical Physics*, 44(15):154013, 2011.
- [7] Peter Lax. *Linear algebra and its applications*. Wiley-Interscience, Hoboken, N.J., 2007.
- [8] Carl Wunsch. *Discrete inverse and state estimation problems: with geophysical fluid applications*. Cambridge University Press, 2006.
- [9] Thomas Mehen. BCH.pdf. <http://webhome.phy.duke.edu/~mehen/760/ProblemSets/BCH.pdf>, Jan 2009. (Accessed on 06/24/2020).
- [10] Ralph M Wilcox. Exponential operators and parameter differentiation in quantum physics. *Journal of Mathematical Physics*, 8(4):962–982, 1967.
- [11] quantum mechanics - Does a time-dependent Hamiltonian commute with its self at different times? <https://physics.stackexchange.com>, jan 2020. (Accessed on 1/15/2020).
- [12] Wen-Xiu Ma, Xiang Gu, and Liang Gao. A note on exact solutions to linear differential equations by the matrix exponential. *Adv. Appl. Math. Mech*, 1(4):573–580, 2009.
- [13] Wen-Xiu Ma and Boris Shekhtman. Do the chain rules for matrix functions hold without commutativity? *Linear and Multilinear Algebra*, 58(1):79–87, 2010.
- [14] James H Liu. A remark on the chain rule for exponential matrix functions. *The College Mathematics Journal*, 34(2):141–143, 2003.
- [15] operators - The formal solution of the time-dependent Schrödinger equation. <https://physics.stackexchange.com>, jan 2020. (Accessed on 1/14/2020).
- [16] Jay A Wood. The chain rule for matrix exponential functions. *The College Mathematics Journal*, 35(3):220–222, 2004.
- [17] quantum field theory - Does the time ordering operator have a rigorous definition? <https://physics.stackexchange.com>, jan 2020. (Accessed on 1/15/2020).
- [18] Jearl Walker and Resnick Halliday. *Fundamentals of physics*. John Wiley & Sons, Inc, Hoboken, NJ, 2014.

- [19] quantum mechanics - Time ordering operator and derivative with respect to time. <https://physics.stackexchange.com>, jan 2020. (Accessed on 1/15/2020).
- [20] quantum mechanics - Utility of the time-ordered exponential. <https://physics.stackexchange.com>, jan 2020. (Accessed on 1/15/2020).