

Multiple solutions for a class of quasilinear Choquard equations ^{*}

Xian Wu [†]

Department of Mathematics, Kunming University, Kunming, Yunnan 650214, P.R.China

Department of Mathematics, Yunnan Normal University, Kunming, Yunnan 650092, P.R.China

email: wuxian2042@163.com

ABSTRACT. In this paper, we study the following quasilinear Choquard equations of the form

$$-\Delta u + V(x)u - \Delta(|u|^{2\alpha})|u|^{2\alpha-2}u = (|x|^{-\mu} * G(u))g(u), \quad x \in \mathbb{R}^N,$$

where $1 \geq \alpha > \frac{1}{2}$, $V \in C(\mathbb{R}^N, \mathbb{R})$, $g \in C(\mathbb{R}^N, \mathbb{R})$. Distinguished from two situations $\lim_{|x| \rightarrow \infty} V(x) = +\infty$ or $\lim_{|x| \rightarrow \infty} V(x) < +\infty$, we research the existence of nontrivial solutions and a sequence of high energy solutions.

1991 Mathematics Subject Classifications: 35J20, 35J70, 35P05, 35P30, 34B15, 58E05, 47H04.

Key Words: quasilinear Choquard equation; dual approach; high energy solution.

1. Introduction and Preliminaries

Consider the following quasilinear Choquard equations of the form

$$-\Delta u + V(x)u - \Delta(|u|^{2\alpha})|u|^{2\alpha-2}u = (|x|^{-\mu} * G(u))g(u), \quad x \in \mathbb{R}^N, \quad (1.1)$$

^{*}This work is supported in partly by the National Natural Science Foundation of China (11771385)

[†]Corresponding Author: wuxian2042@163.com

where $1 \geq \alpha > \frac{1}{2}$, $V \in C(\mathbb{R}^N, \mathbb{R})$, $g \in C(\mathbb{R}^N, \mathbb{R})$.

As $\alpha = 1$ and $\mu = 0$, the equation (1.1) degenerate into the form

$$-\Delta u + V(x)u - \Delta(|u|^2)u = g(u), \quad x \in \mathbb{R}^N. \quad (1.2)$$

Solutions of the equation (1.2) are standing waves the following quasilinear Schrödinger equation of the form

$$i\psi_t + \Delta\psi - V(x)\psi + k\Delta(h(|\psi|^2))h'(|\psi|^2)\psi + g(\psi) = 0, \quad x \in \mathbb{R}^N. \quad (1.3)$$

The quasilinear Schrödinger equations (1.3) are derived as models of several physical phenomena, such as see [9, 12, 13, 20, 23]. It begins with [22] for the studies on Mathematics. In the resent years, greater important attention has been paid to the equation (1.2), for example, see [5, 6, 7, 11, 14, 15, 16, 19, 24, 26]. Especially, in [14], the ground state solutions for the following problems

$$-\Delta u + V(x)u - \Delta(|u|^{2\alpha})|u|^{2\alpha-2}u = \lambda|u|^{p-1}u, \quad x \in \mathbb{R}^N$$

was studied via the Lagrange multiplier method; in [3], the uniqueness of the ground state solutions for the following the problems

$$-\Delta u + \lambda u - \kappa\Delta(|u|^{2\alpha})|u|^{2\alpha-2}u = |u|^{p-1}u, \quad x \in \mathbb{R}^N$$

was studied via a dual approach.

As $\alpha = \frac{1}{2}$, the equation (1.1) degenerate into the form

$$-\Delta u + V(x)u = (|x|^{-\mu} * G(u))g(u), \quad x \in \mathbb{R}^N. \quad (1.4)$$

(1.4) first appeared in [21], it was used to describe the quantum mechanics of a polaron. Next, Choquard used (1.4) to describe an electron trapped in its own hole (see [10]). In [17], (1.4) was used as a model of self-gravitating matter. Recently, greater important attention has been paid to the equation (1.4), for example, see [1, 2, 4, 18, 27, 30, 31].

In [29], the author had studied the existence of positive solutions, negative solutions and sequence of high energy solutions for the equation (1.1) with $\frac{1}{2} < \alpha \leq 1$, $\mu = 0$ and

$\lim_{|x| \rightarrow \infty} V(x) = +\infty$. In the present paper, we study the equation (1.1) with $\frac{1}{2} < \alpha \leq 1$, $0 < \mu < 2 < N$. Distinguished from two situations $\lim_{|x| \rightarrow \infty} V(x) = +\infty$ or $\lim_{|x| \rightarrow \infty} V(x) < +\infty$, we research the existence of nontrivial solutions and a sequence of high energy solutions.

In order to reduce the statements for main results, we list the assumptions as follows:

$$(V) \quad V \in C(\mathbb{R}^N, \mathbb{R}), \quad 0 < V_0 := \inf_{x \in \mathbb{R}^N} V(x) \text{ and } V_\infty := \lim_{|x| \rightarrow \infty} V(x) \leq +\infty.$$

$$(g_1) \quad g \in C(\mathbb{R}, \mathbb{R}), \quad tg(t) > 0 \text{ for all } t \neq 0, \text{ and there exist } C_1 > 0, 2\alpha(2 - \frac{\mu}{N}) \leq q_1 \leq q_2 < 2\alpha(\frac{2N-\mu}{N-2}) \text{ such that}$$

$$|g(t)| \leq C_1(|t|^{q_1-1} + |t|^{q_2-1})$$

for all $t \in \mathbb{R}$.

$$(g_2)$$

$$tg(t) - 4\alpha G(t) \geq 0, \quad \forall t \in \mathbb{R},$$

where $G(t) = \int_0^t g(s)ds$.

Set

$$H^1(\mathbb{R}^N) = \{u \in L^2(\mathbb{R}^N) : \nabla u \in L^2(\mathbb{R}^N)\}$$

with the inner product

$$\langle u, v \rangle_{H^1} = \int_{\mathbb{R}^N} (\nabla u \cdot \nabla v + uv) dx$$

and the norm

$$\|u\|_{H^1} = [\int_{\mathbb{R}^N} (|\nabla u|^2 + u^2) dx]^{1/2}.$$

When $V_\infty = +\infty$, set

$$E = \{u \in H^1(\mathbb{R}^N) : \int_{\mathbb{R}^N} [|\nabla u|^2 + V(x)u^2] dx < +\infty\}$$

with the inner product

$$\langle u, v \rangle_E = \int_{\mathbb{R}^N} [\nabla u \cdot \nabla v + V(x)uv] dx$$

and the norm

$$\|u\|_E = \langle u, u \rangle_E^{1/2}.$$

When $V_\infty < +\infty$, set $E = H^1(\mathbb{R}^N)$. Then E is a Hilbert space. By the continuity of the embedding

$$E \hookrightarrow L^s(\mathbb{R}^N), s \in [2, 2^*],$$

there exist constants $a_s > 0$, $2 \leq s \leq 2^*$, such that

$$\|u\|_s \leq a_s \|u\|_E, \quad \forall u \in E,$$

where we denote by $\|\cdot\|_s$ the norm of $L^s(\mathbb{R}^N)$. Moreover, by Lemma 3.4 in [33] we know that as $V_\infty = +\infty$, the embedding $E \hookrightarrow L^s(\mathbb{R}^N)$ is compact for each $2 \leq s < 2^*$.

We observe that formally equation (1.1) is the Euler-Lagrange equation associated of the natural energy functional $J : E \rightarrow \mathbb{R}$ given by

$$J(u) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} V(x) u^2 dx + \frac{1}{4\alpha} \int_{\mathbb{R}^N} |\nabla(|u|^{2\alpha})|^2 dx - \frac{1}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(u)) G(u) dx.$$

Clearly, $\frac{1}{4\alpha} \int_{\mathbb{R}^N} |\nabla(|u|^{2\alpha})|^2 dx = \alpha \int_{\mathbb{R}^N} |u|^{2(2\alpha-1)} |\nabla u|^2 dx$. Hence

$$J(u) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} V(x) u^2 dx + \alpha \int_{\mathbb{R}^N} |u|^{2(2\alpha-1)} |\nabla u|^2 dx - \frac{1}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(u)) G(u) dx.$$

According to [3], we can define f by

$$f'(t) = \frac{1}{\sqrt{1 + 2\alpha|f(t)|^{2(2\alpha-1)}}} \quad \text{on } t \in [0, +\infty)$$

and

$$f(-t) = -f(t) \quad \text{on } t \in (-\infty, 0].$$

After the change of variables, we obtain the following functional

$$I(v) := J(f(v)) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla v|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} V(x) f^2(v) dx - \frac{1}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v))) G(f(v)) dx,$$

which is well defined in E under the assumptions (V) and (g_1) . Moreover, a critical point v of the functional I corresponds to a weak solutions of the following equation

$$-\Delta v = \frac{1}{\sqrt{1 + 2\alpha|f(v)|^{2(2\alpha-1)}}} [(|x|^{-\mu} * G(f(v))) g(f(v)) - V(x) f(v)] \quad \text{in } \mathbb{R}^N \quad (1.5)$$

and $u = f(v)$ is a weak solution (1.1).

The following lemma appeared in [29].

Lemma 2.1 The function $f(t)$ enjoys the following properties:

(f_1) f is uniquely defined C^∞ function and invertible.

(f_2) $0 < f'(t) \leq 1, \forall t \in \mathbb{R}$.

(f_3) $|f(t)| \leq |t|, \forall t \in \mathbb{R}$.

(f_4) $\lim_{t \rightarrow 0} \frac{f(t)}{t} = 1$.

(f_5) $\lim_{t \rightarrow +\infty} \frac{f^{2\alpha}(t)}{t} = \sqrt{2\alpha}, \quad \lim_{t \rightarrow -\infty} \frac{f^{2\alpha}(t)}{t} = -\sqrt{2\alpha}$.

(f_6) $\frac{f(t)}{2} \leq \alpha t f'(t) \leq \alpha f(t), \quad \forall t \geq 0; \quad \alpha f(t) \leq \alpha t f'(t) \leq \frac{f(t)}{2}, \quad \forall t \leq 0$.

(f_7) $f^{2\alpha}(t) \leq \sqrt{2\alpha}|t|, \quad \forall t \in \mathbb{R}$.

(f_8) The function $f^2(t)$ is strictly convex.

(f_9) There exists a positive constant $\theta > 0$ such that

$$|f(t)| \geq \begin{cases} \theta|t|, & |t| \leq 1, \\ \theta|t|^{\frac{1}{2\alpha}}, & |t| \geq 1. \end{cases}$$

(f_{10}) There exist positive constants C_1 and C_2 such that

$$|t| \leq C_1|f(t)| + C_2|f(t)|^{2\alpha}, \quad \forall t \in \mathbb{R}.$$

(f_{11}) $|f^{2\alpha-1}(t)f'(t)| < \frac{1}{\sqrt{2\alpha}}, \quad \forall t \in \mathbb{R}$.

(f_{12}) $f(t)$ is odd, $f^2(t)$ is even.

(f_{13}) For each $\xi > 0$, there exists a positive constant $C(\xi)$ such that

$$f^{2\alpha}(\xi t) \leq C(\xi)f^{2\alpha}(t).$$

(f_{14}) The function $f(t)f'(t)t^{-1}$ is strictly decreasing for $t > 0$.

(f_{15}) The function $f^p(t)f'(t)t^{-1}$ is strictly increasing for $p \geq 4\alpha - 1$ and $t > 0$.

By the Hardy-Littlewood-Sobolev inequality (For example, see Proposition 1.1 in [4]) and Lemma 2.1 we can prove the following Lemma 2.2.

Lemma 2.2 If assumptions (V) and (g_1) hold, then the functionals I is well defined on E , and $I \in C^1(E, \mathbb{R})$.

Through out the paper, C and C_i are used in various places to denote positive constants.

3. Main results

Theorem 3.1 Assume the conditions (V) and (g_1) -(g_2) hold. Then the equation (1.1) has a nontrivial solution. Furthermore, if g is odd, then the equation (1.1) has a sequence of solutions $\{u_n\} \subset E$ such that $\|u_n\|_E \rightarrow \infty$ and $J(u_n) \rightarrow +\infty$.

Proof. First, we prove that I satisfies the Cerami condition. Let $\{v_n\} \subset E$ be any Cerami sequence of I , i.e. $\{I(v_n)\}$ is bounded and $(1 + \|v_n\|_E)I'(v_n) \rightarrow 0$ in E^* . Set $A_n^2 := \int_{\mathbb{R}^N} [|\nabla v_n|^2 + V(x)f^2(v_n)]dx$. By Lemma 2.1 (f_{11}) and (g_2) , there exists a constant $C > 0$ such that

$$\begin{aligned} C &\geq I(v_n) - \frac{1}{8\alpha} \langle I'(v_n), \frac{f(v_n)}{f'(v_n)} \rangle \\ &\geq \frac{1}{4} \int_{\mathbb{R}^N} |\nabla v_n|^2 dx + \frac{4\alpha - 1}{8\alpha} \int_{\mathbb{R}^N} V(x)f^2(v_n) dx \\ &\quad + \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v_n))) \left[\frac{1}{8\alpha} f(v_n)g(f(v_n)) - \frac{1}{2} G(f(v_n)) \right] dx \\ &\geq \frac{1}{4} A_n^2. \end{aligned}$$

This shows that the sequence $\{A_n\}$ is bounded. Moreover, from Step 2 of proof of Theorem 3.1 in [29] we know that there exists a constant $C > 0$ such that $A_n^2 \geq C\|v_n\|_E^2$, and hence the sequence $\{v_n\}$ is bounded in E . Note that, for each $t \in \mathbb{R}$, by Lemma 2.1 (f_{11}) one has

$$\frac{d}{dt} [f(t)f'(t)] = (f'(t))^2 + f(t)f''(t) > 2(1 - \alpha)|f'(t)|^2 = \frac{2(1 - \alpha)}{1 + 2\alpha|f(t)|^{2(2\alpha-1)}} \geq 0.$$

Hence, as the proof of Lemma 3.11 in [8], we may prove that there is a constant $C > 0$ such

that

$$\int_{\mathbb{R}^N} |\nabla(v_n - v)|^2 dx + \int_{\mathbb{R}^N} V(x)[f(v_n)f'(v_n) - f(v)f'(v)](v_n - v)dx \geq C\|v_n - v\|_E^2. \quad (3.1)$$

Moreover, by (g_1) and Lemma 2.1 (f_7) , (f_{11}) , we have

$$|g(f(t))f'(t)| \leq C_2(|t|^{\frac{q_1}{2\alpha}-1} + |t|^{\frac{q_2}{2\alpha}-1}), \quad \forall t \in \mathbb{R}, \quad (3.2)$$

and

$$|G(f(t))| \leq C_2(|t|^{\frac{q_1}{2\alpha}} + |t|^{\frac{q_2}{2\alpha}}), \quad \forall t \in \mathbb{R}. \quad (3.3)$$

Set $s = \frac{2N}{2N-\mu}$. Then $s \in (1, \frac{N}{\mu})$ and $G(f(w))$, $g(f(w))f'(w)w \in L^s(\mathbb{R}^N)$ for all $w \in E$.

Notice that

$$2 \leq \frac{sq_1}{2\alpha} \leq \frac{sq_2}{2\alpha} < 2^*.$$

(1⁰) For the case $V_\infty = +\infty$, then by the compactness of embedding $E \hookrightarrow L^s(\mathbb{R}^N)$ ($2 \leq s < 2^*$), up to a subsequence, one has $v_n \rightharpoonup v$ in E , $v_n \rightarrow v$ in $L^s(\mathbb{R}^N)$ for all $2 \leq s < 2^*$ and $v_n(x) \rightarrow v(x)$ a.e. on \mathbb{R}^N . By the Hardy-Littlewood-Sobolev inequality we know that

$$\begin{aligned} & \left| \int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(v_n)))g(f(v_n))f'(v_n) - (|x|^{-\mu} * G(f(v)))g(f(v))f'(v)](v_n - v)dx \right| \\ & \leq C[\|G(f(v_n))\|_s \cdot \|g(f(v_n))f'(v_n)(v_n - v)\|_s + \|G(f(v))\|_s \cdot \|g(f(v))f'(v)(v_n - v)\|_s] \\ & \leq C[\|g(f(v_n))f'(v_n)(v_n - v)\|_s + \|g(f(v))f'(v)(v_n - v)\|_s] \\ & \leq C(\|v_n - v\|_{\frac{sq_1}{2\alpha}} + \|v_n - v\|_{\frac{sq_2}{2\alpha}}) \rightarrow 0. \end{aligned} \quad (3.4)$$

(2⁰) For the case $V_\infty < +\infty$, then by the boundedness of $\{v_n\}$ in E , up to a subsequence, we can assume that $v_n \rightharpoonup v$ in E , $v_n \rightarrow v$ in $L_{loc}^s(\mathbb{R}^N)$ for each $s \in [1, 2^*)$ and $v_n(x) \rightarrow v(x)$ a.e. $x \in \mathbb{R}^N$. By (3.2) and (3.4) we know that

$$\begin{aligned} & \left| \int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(v_n)))g(f(v_n))f'(v_n) - (|x|^{-\mu} * G(f(v)))g(f(v))f'(v)](v_n - v)dx \right| \\ & \leq C[\|g(f(v_n))f'(v_n)(v_n - v)\|_s + \|g(f(v))f'(v)(v_n - v)\|_s] \end{aligned} \quad (3.5)$$

and for any $\varepsilon > 0$ there exist $0 < \delta_0 < \rho_0$ such that

$$|g(f(t))f'(t)| \leq \varepsilon(|t|^{\frac{N-\mu}{N}} + |t|^{\frac{2+N-\mu}{N-2}}) + \chi_{[\delta_0, \rho_0]}(|t|)|g(f(t))f'(t)|, \quad \forall t \in \mathbb{R},$$

where $\chi_{[\delta_0, \rho_0]}$ is the characteristic function on $[\delta_0, \rho_0]$. Hence

$$\begin{aligned} & \|g(f(v_n))f'(v_n)(v_n - v)\|_s^s \\ & \leq C\varepsilon \int_{\mathbb{R}^N} (|v_n|^{\frac{s(N-\mu)}{N}} + |v_n|^{\frac{s(2+N-\mu)}{N-2}})|v_n - v|^s dx \\ & \quad + C \int_{\mathbb{R}^N} \chi_{[\delta_0, \rho_0]}(|v_n|)|g(f(v_n))f'(v_n)|^s |v_n - v|^s dx \end{aligned} \quad (3.6)$$

and

$$\begin{aligned} & C\varepsilon \int_{\mathbb{R}^N} (|v_n|^{\frac{s(N-\mu)}{N}} + |v_n|^{\frac{s(2+N-\mu)}{N-2}})|v_n - v|^s dx \\ & \leq C\varepsilon [\|v_n\|_2^{\frac{2(N-\mu)}{2N-\mu}} \|v_n - v\|_2^s + \|v_n - v\|_{2^*}^s \|v_n\|_{2^*}^{\frac{2^*(2+N-\mu)}{2N-\mu}}] \leq C\varepsilon. \end{aligned} \quad (3.7)$$

For any $r > 0$ and $n \in \mathbb{N}$, set $B_r := \{x \in \mathbb{R}^N : |x| < r\}$, $B_r^c := \mathbb{R}^N \setminus B_r$ and

$$A_n := \{x \in \mathbb{R}^N : \delta_0 \leq |v_n(x)| \leq \rho_0\}.$$

Then

$$\begin{aligned} & C \int_{\mathbb{R}^N} \chi_{[\delta_0, \rho_0]}(|v_n|)|g(f(v_n))f'(v_n)|^s |v_n - v|^s dx \\ & = C \int_{A_n} |g(f(v_n))f'(v_n)|^s |v_n - v|^s dx \\ & = C \left[\int_{A_n \cap B_r} |g(f(v_n))f'(v_n)|^s |v_n - v|^s dx + \int_{A_n \cap B_r^c} |g(f(v_n))f'(v_n)|^s |v_n - v|^s dx \right] \\ & \leq C \int_{A_n \cap B_r} |g(f(v_n))f'(v_n)|^s |v_n - v|^s dx + C|A_n \cap B_r^c|, \end{aligned} \quad (3.8)$$

where $|A_n \cap B_r^c|$ denotes the Lebesgue measure of $A_n \cap B_r^c$. Similar to (3.4), we have

$$C \int_{A_n \cap B_r} |g(f(v_n))f'(v_n)|^s |v_n - v|^s dx = o_n(1). \quad (3.9)$$

Now, we prove

$$\lim_{r \rightarrow +\infty} |A_n \cap B_r^c| = 0. \quad (3.10)$$

Indeed, if this is not true, then there exist $\delta > 0$ and $r_k \uparrow +\infty$ such that

$$|A_n \cap B_{r_k}^c| \geq \delta, \quad \forall k \in \mathbb{N}.$$

Obviously,

$$|A_n \cap B_{r_k}^c| \leq |A_n| := \beta_n < +\infty, \quad \forall k \in \mathbb{N}.$$

On the other hand, set $\Omega_k := B_{r_k}^c \setminus B_{r_{k+1}}^c$. We have $\Omega_i \cap \Omega_j = \emptyset$ whenever $i \neq j$, and

$$B_{r_k}^c = \bigcup_{i=k}^{\infty} \Omega_i, \quad \forall k \in \mathbb{N}.$$

Hence

$$\delta \leq |A_n \cap B_{r_k}^c| = \sum_{i=k}^{\infty} |A_n \cap \Omega_i|, \quad \forall k \in \mathbb{N},$$

and hence

$$\beta_n \geq |A_n \cap B_{r_1}^c| = \sum_{i=1}^{\infty} |A_n \cap \Omega_i| = +\infty.$$

This is a contradiction. Therefore, (3.10) holds. Now, we prove

$$\lim_{r \rightarrow +\infty} |A_n \cap B_r^c| = 0 \quad \text{uniformly in } n \in \mathbb{N}. \quad (3.11)$$

In fact, for any $\varepsilon > 0$ there exists a $r_0 > 1$ such that

$$\int_{B_r^c} |v|^2 dx < \varepsilon \quad \text{whenever } r \geq r_0.$$

Take $t_1 = r_0, t_j \uparrow +\infty$ be such that $D_j := B_{t_j}^c \setminus B_{t_{j+1}}^c$, $B_{r_0}^c = \bigcup_{j=1}^{\infty} D_j$ and

$$\int_{D_j} |v|^2 dx < \frac{\varepsilon}{2^j}, \quad \forall j \in \mathbb{N}.$$

By Faou Lemma, we get

$$\limsup_{n \rightarrow \infty} \int_{A_n \cap D_j} |v_n|^2 dx \leq \int_{D_j} |v|^2 dx < \frac{\varepsilon}{2^j}, \quad \forall j \in \mathbb{N}.$$

Hence, as $r \geq r_0$, one has

$$\begin{aligned} \delta_0^2 \limsup_{n \rightarrow \infty} |A_n \cap B_r^c| &\leq \delta_0^2 \limsup_{n \rightarrow \infty} |A_n \cap B_{r_0}^c| \\ &\leq \limsup_{n \rightarrow \infty} \int_{A_n \cap B_{r_0}^c} |v_n|^2 dx \\ &= \limsup_{n \rightarrow \infty} \sum_{j=1}^{\infty} \int_{A_n \cap D_j} |v_n|^2 dx \\ &< \sum_{j=1}^{\infty} \frac{\varepsilon}{2^j} = \varepsilon. \end{aligned}$$

Together with (3.10), we know that (3.11) holds.

For any $\varepsilon > 0$, by (3.11) there exists a large $r > 0$ such that $|A_n \cap B_r^c| < \varepsilon$. Using (3.6) – (3.9) we know

$$\|g(f(v_n))f'(v_n)(v_n - v)\|_s^s = o_n(1). \quad (3.12)$$

Similarly, we have

$$\|g(f(v))f'(v)(v_n - v)\|_s. \quad (3.13)$$

Consequently, from (3.5) we get

$$\int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(v_n)))g(f(v_n))f'(v_n) - (|x|^{-\mu} * G(f(v)))g(f(v))f'(v)](v_n - v)dx = o_n(1).$$

Summing up (1⁰) and (2⁰), together with (3.1) we obtain

$$\begin{aligned} o_n(1) &= \langle I'(v_n) - I'(v), v_n - v \rangle \\ &= \int_{\mathbb{R}^N} |\nabla(v_n - v)|^2 dx + \int_{\mathbb{R}^N} V(x)[f(v_n)f'(v_n) - f(v)f'(v)](v_n - v)dx \\ &\quad - \int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(v_n)))g(f(v_n))f'(v_n) - (|x|^{-\mu} * G(f(v)))g(f(v))f'(v)](v_n - v)dx \\ &\geq C\|v_n - v\|_E^2 + o_n(1). \end{aligned}$$

Hence $v_n \rightarrow v$ in E . This shows that I satisfies the Cerami condition.

Next, we prove that I has a mountain pass geometry. Indeed, set $A(v) := [\int_{\mathbb{R}^N} (|\nabla v|^2 + V(x)f^2(v))dx]^{\frac{1}{2}}$. As paragraph 1 of proof of Lemma 3.3 in [8] we can prove that there exist $C, \rho_1 > 0$ such that

$$A^2(v) \geq C\|v\|_E^2, \quad \text{whenever } \|v\|_E \leq \rho_1. \quad (3.14)$$

Similar to (3.4) we know that

$$|\int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(v)))G(f(v))dx] \leq C\|G(f(v))\|_s^2 \leq C(\|v\|_E^{\frac{q_1}{\alpha}} + \|v\|_E^{\frac{q_2}{\alpha}})$$

Consequently, for small $0 < \rho < \min\{1, \rho_1\}$, we have

$$I(v) \geq C[\|v\|_E^2 - \|v\|_E^{\frac{q_1}{\alpha}} - \|v\|_E^{\frac{q_2}{\alpha}}] = C(\rho^2 - \rho^{\frac{q_1}{\alpha}} - \rho^{\frac{q_2}{\alpha}}) := \delta > 0 \text{ whenever } \|v\|_E = \rho.$$

Take an $e \in E$ with $\|e\|_E = 1$. Set $B(e) := \frac{1}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(e)))G(f(e))dx$ and $\beta(t) := B(te)$ for $t > 0$. Then, by (g₂) and Lemma 2.1 (f_6), one has

$$\frac{\beta'(t)}{\beta(t)} \geq \frac{4}{t}, \quad \forall t > 0. \quad (3.15)$$

For $s > 1$, integrating (3.15) over $[1, s]$ we obtain $B(se) \geq s^4 B(e)$. Hence

$$I(se) \leq \frac{1}{2}s^2 - \frac{1}{2}s^4 \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(e)))G(f(e))dx \rightarrow -\infty \quad (3.16)$$

as $s \rightarrow +\infty$. Consequently, there is a $s_0 > \max\{1, \rho\}$ such that $I(v_0) := I(s_0 e) < 0$. This shows that I has a mountain pass geometry.

Moreover, for any finite-dimensional subspace $\tilde{E} \subset E$, we assert that there exists a constant $R > \rho$ such that $I < 0$ on $\tilde{E} \setminus B_R$. Otherwise, there is a sequence $\{v_n\} \subset \tilde{E}$ such that $s_n := \|v_n\|_E \rightarrow \infty$ and $I(v_n) \geq 0$. Set $e_n := \frac{v_n}{\|v_n\|_E}$. By (3.16), one has

$$0 \leq I(v_n) = I(s_n e_n) \leq \frac{1}{2}s_n^2 - \frac{1}{2}s_n^4 \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(e_n)))G(f(e_n))dx,$$

and hence

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(e_n)))G(f(e_n))dx = 0.$$

By the Fatou Lemma we know that

$$\int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(e)))G(f(e))dx = 0.$$

Hence $e(x) = 0$ a.e. $x \in \mathbb{R}^N$. By the equivalency of all norms in \tilde{E} , there is a constant $C_2 > 0$ such that

$$\|v\|_2^2 \geq C_2 \|v\|_E^2, \quad \forall v \in \tilde{E}.$$

Hence

$$0 = \lim_{n \rightarrow \infty} \|e_n\|_2^2 \geq \lim_{n \rightarrow \infty} C_2 \|e_n\|_E^2 = C_2,$$

a contradiction. This shows that there exists a constant $R > 0$ such that $I < 0$ on $\tilde{E} \setminus B_R$.

Since $E \hookrightarrow L^2(\mathbb{R}^N)$ and $L^2(\mathbb{R}^N)$ is a separable Hilbert space, E has a countable orthogonal basis $\{e_j\}$. Set $E_k := \text{span}\{e_1, \dots, e_k\}$ and $Z_k := E_k^\perp$, then $E = E_k \oplus Z_k$. Hence

$$I|_{S_\rho \cap Z_k} \geq \delta > 0.$$

Notice that the Deformation Theorem still hold under the Cerami condition (see [32]). Hence Theorem 2.2 and Theorem 9.12 in [25] hold under the Cerami condition. Therefore,

Theorem 3.1 follows from Theorem 2.2 and Theorem 9.12 in [25]. This completes the proof.

□

Remark 3.2 From the proof of Theorem 3.1 we know that when $V_\infty = +\infty$, the limit of q_1 can relax as $2\alpha(2 - \frac{\mu}{N}) \leq q_1$.

Theorem 3.3 Assume the conditions (V) with $V_\infty < +\infty$ and (g_1) hold. If $V \in C^1(\mathbb{R}^N, \mathbb{R})$ with $x \cdot \nabla V(x) \in L^\infty(\mathbb{R}^N)$ and $V(rx)$ is non-increasing in $r \in \mathbb{R}$. Then the equation (1.1) has a nontrivial solution. Especially, if $V \equiv 1$, then the equation (1.1) has a ground state solution.

Proof. Set

$$P(v) := \frac{N-2}{2} \int_{\mathbb{R}^N} |\nabla v|^2 dx + \frac{N}{2} \int_{\mathbb{R}^N} V(x) f^2(v) dx + \frac{1}{2} \int_{\mathbb{R}^N} (x \cdot \nabla V(x)) f^2(v) dx \\ - \frac{2N-\mu}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v))) G(f(v)) dx.$$

Define a mapping $\Phi : \mathbb{R} \times E \rightarrow E$ by

$$\Phi(r, v)(x) := v(e^{-r}x), \quad \forall (r, v) \in \mathbb{R} \times E \rightarrow E.$$

Then

$$I(\Phi(r, v)) = \frac{e^{(N-2)r}}{2} \int_{\mathbb{R}^N} |\nabla v|^2 dx + \frac{e^{Nr}}{2} \int_{\mathbb{R}^N} V(e^r x) f^2(v) dx \\ - \frac{e^{(2N-\mu)r}}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v))) G(f(v)) dx.$$

Set

$$\tilde{\Gamma} := \{\tilde{\gamma} \in C([0, 1], \mathbb{R} \times E) : \tilde{\gamma}(0) = (0, 0), \quad I \circ \Phi(\tilde{\gamma}(1)) < 0\}.$$

Then

$$\Gamma := \{\Phi \circ \tilde{\gamma} : \tilde{\gamma} \in \tilde{\Gamma}\} = \{\gamma \in C([0, 1], E) : \gamma(0) = 0, \quad I(\gamma(1)) < 0\}$$

and

$$\inf_{\tilde{\gamma} \in \tilde{\Gamma}} \sup_{t \in [0, 1]} I \circ \Phi(\tilde{\gamma}(t)) = \inf_{\gamma \in \Gamma} \sup_{t \in [0, 1]} I(\gamma(t)) := c.$$

From the proof of Theorem 3.1 we know that for small $\rho > 0$, we have

$$I(v) \geq \delta > 0 \text{ whenever } \|v\|_E = \rho.$$

Moreover, for any $v \in E \setminus \{0\}$ and $t > 0$, set $v_t(x) = v(\frac{x}{t})$. Then

$$\begin{aligned} I(v_t) &= \frac{t^{N-2}}{2} \int_{\mathbb{R}^N} |\nabla v|^2 dx + \frac{t^N}{2} \int_{\mathbb{R}^N} V(tx) f^2(v) dx - \frac{t^{2N-\mu}}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v))) G(f(v)) dx \\ &\leq \frac{t^{N-2}}{2} \int_{\mathbb{R}^N} |\nabla v|^2 dx + \frac{t^N}{2} C \int_{\mathbb{R}^N} f^2(v) dx - \frac{t^{2N-\mu}}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(v))) G(f(v)) dx \\ &\rightarrow -\infty \end{aligned}$$

as $t \rightarrow +\infty$. Hence, I has a mountain pass geometry, and hence $c > 0$. Using Theorem 2.9 in [28], for $M = [0, 1]$, $M_0 = \{0, 1\}$, $X = \mathbb{R} \times E$ and $\varphi = I \circ \Phi$, we know that there exists a sequence $\{(r_n, v_n)\} \subset \mathbb{R} \times E$ such that

$$\lim_{n \rightarrow \infty} I(\Phi(r_n, v_n)) = c \quad \text{and} \quad (I \circ \Phi)'(r_n, v_n) \rightarrow 0.$$

Notice that

$$\langle (I \circ \Phi)'(r_n, v_n), (r, v) \rangle = \langle I'(\Phi(r_n, v_n)), \Phi(r_n, v) \rangle + rP(\Phi(r_n, v_n)), \quad \forall (r, v) \in \mathbb{R} \times E.$$

We have

$$I'(\Phi(r_n, v_n)) \rightarrow 0, \quad P(\Phi(r_n, v_n)) \rightarrow 0.$$

Set $w_n := \Phi(r_n, v_n)$. Then

$$\begin{aligned} c + o_n(1) &= I(w_n) - \frac{1}{2N - \mu} P(w_n) \\ &= \frac{e^{(N-2)r_n}}{2} \cdot \frac{N + 2 - \mu}{2N - \mu} \int_{\mathbb{R}^N} |\nabla v_n|^2 dx + \frac{e^{Nr_n}}{2} \cdot \frac{N - \mu}{2N - \mu} \int_{\mathbb{R}^N} V(e^{r_n} x) f^2(v_n) dx \\ &\quad - \frac{e^{(N+2)r_n}}{2(2N - \mu)} \int_{\mathbb{R}^N} (x \cdot \nabla V(e^{r_n} x)) f^2(v_n) dx. \end{aligned}$$

For any $x \in \mathbb{R}^N$, set $h(r) := V(rx)$. Then $0 \geq h'(r) = x \cdot \nabla V(rx)$ for all $r \in \mathbb{R}$. Hence

$$\begin{aligned} c + o_n(1) &= I(w_n) - \frac{1}{2N - \mu} P(w_n) \\ &\geq \frac{e^{(N-2)r_n}}{2} \cdot \frac{N + 2 - \mu}{2N - \mu} \int_{\mathbb{R}^N} |\nabla v_n|^2 dx + \frac{e^{Nr_n}}{2} \cdot \frac{N - \mu}{2N - \mu} \int_{\mathbb{R}^N} V(e^{r_n} x) f^2(v_n) dx \\ &= \frac{N + 2 - \mu}{2(2N - \mu)} \int_{\mathbb{R}^N} |\nabla w_n|^2 dx + \frac{N - \mu}{2(2N - \mu)} \int_{\mathbb{R}^N} V(x) f^2(w_n) dx \\ &\geq \frac{N - \mu}{2(2N - \mu)} A^2(w_n) \geq C \|w_n\|_E^2. \end{aligned}$$

This shows $\{w_n\} \subset E$ is bounded.

Next, we prove that there exist $\sigma_0 > 0$, $\delta_0 > 0$ and $\{x_n\} \subset \mathbb{R}^N$ such that

$$\liminf_{n \rightarrow \infty} \int_{B_{\sigma_0}(x_n)} |w_n|^2 \geq \delta_0. \quad (3.17)$$

Indeed, if the conclusion is not true, then by Lemma 1.21 in [28] we know that

$$w_n \rightarrow 0 \text{ in } L^s(\mathbb{R}^N), \quad \forall s \in [2, 2^*).$$

Similar to the case (1⁰) of Theorem 3.1 we can prove

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^N} [(|x|^{-\mu} * G(f(w_n)))g(f(w_n))f(w_n)]dx = 0.$$

Hence

$$\begin{aligned} o_n(1) &= \langle I'(w_n), \frac{f(w_n)}{f'(w_n)} \rangle \\ &= \int_{\mathbb{R}^N} [1 + (2\alpha - 1) \cdot \frac{2\alpha|f(w_n)|^{2(2\alpha-1)}}{1 + 2\alpha|f(w_n)|^{2(2\alpha-1)}}] |\nabla w_n|^2 dx + \int_{\mathbb{R}^N} V(x) f^2(w_n) dx + o_n(1) \\ &\geq C \|w_n\|_E^2 + o_n(1), \end{aligned}$$

and hence $\|w_n\|_E \rightarrow 0$. Consequently, $0 < c = \lim_{n \rightarrow \infty} I(w_n) = 0$, a contradiction. Therefore, (3.17) holds.

By the boundedness of $\{w_n\}$ and (3.17) we know that, up to a subsequence, there exists a $v_0 \in E \setminus \{0\}$ such that $w_n \rightharpoonup v_0$ in E , $w_n \rightarrow v_0$ in $L_{loc}^t(\mathbb{R}^N)$ for each $t \in [1, 2^*)$ and $w_n(x) \rightarrow v_0(x)$ a.e. $x \in \mathbb{R}^N$. From the proof of the case (2⁰) of Theorem 3.1 we know $w_n \rightarrow v_0$ in E . Consequently, $I(v_0) = c$, $I'(v_0) = 0$ and $P(v_0) = 0$. This shows that v_0 is a nontrivial solution of (1.5). Consequently, $u_0 := f(v_0)$ is a nontrivial solution of (1.1).

If $V \equiv 1$, then set $\mathcal{M} := \{v \in E \setminus \{0\} : I'(v) = 0\}$ and $c^* := \inf_{v \in \mathcal{M}} I(v)$. Since $v_0 \in \mathcal{M}$, $c \geq c^*$. On the other hand, for any $w_0 \in M$, one has Pohozaev identity $P(w_0) = 0$. Define $\gamma : [0, +\infty) \rightarrow E$ by

$$\gamma(t)(x) = \begin{cases} w_0(\frac{x}{t}), & t > 0, \\ 0, & t = 0. \end{cases}$$

We claim that $\gamma : [0, +\infty) \rightarrow E$ is continuous. Indeed, if $t_0 \in (0, +\infty)$, $0 < t_n \rightarrow t_0$, then

$$\|\gamma(t_n)\|_E^2 = t_n^{N-2} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + t_n^N \int_{\mathbb{R}^N} |w_0|^2 dx \rightarrow \|\gamma(t_0)\|_E^2.$$

Hence there exists a $v \in E$ such that $\gamma(t_n) \rightharpoonup v$ in E , $\gamma(t_n) \rightarrow v$ in $L_{loc}^p(\mathbb{R}^N)$ for each $p \in [1, 2^*)$, $\gamma(t_n)(x) \rightarrow v(x)$ a.e. $x \in \mathbb{R}^N$. Notice that for any $\varphi \in C_0^\infty(\mathbb{R}^N)$, one has

$$\int_{\mathbb{R}^N} \gamma(t_n)(x) \varphi(x) dx = \int_{\mathbb{R}^N} t_n^N w_0(x) \varphi(t_n x) dx \rightarrow \int_{\mathbb{R}^N} \gamma(t_0)(x) \varphi(x) dx.$$

By the density of $C_0^\infty(\mathbb{R}^N)$ in $L^2(\mathbb{R}^N)$ we know that $\gamma(t_n) \rightharpoonup \gamma(t_0)$ in $L^2(\mathbb{R}^N)$. Consequently, $\gamma(t_0) = v$, and hence $\gamma(t_n) \rightarrow \gamma(t_0)$ in E . This shows that $\gamma : (0, +\infty) \rightarrow E$ is continuous. Moreover, as $0 < t_n \rightarrow 0$, we have $\|\gamma(t_n)\|_E \rightarrow 0$, obviously. Therefore, $\gamma : [0, +\infty) \rightarrow E$ is continuous.

Notice that as $t > 0$,

$$\begin{aligned} I(\gamma(t)) &= \frac{t^{N-2}}{2} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + \frac{t^N}{2} \int_{\mathbb{R}^N} f^2(w_0) dx \\ &\quad - \frac{t^{2N-\mu}}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(w_0))) G(f(w_0)) dx \end{aligned}$$

and

$$\begin{aligned} \frac{d}{dt} I(\gamma(t)) &= \frac{(N-2)t^{N-3}}{2} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + \frac{Nt^{N-1}}{2} \int_{\mathbb{R}^N} f^2(w_0) dx \\ &\quad - \frac{(2N-\mu)t^{2N-\mu-1}}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(w_0))) G(f(w_0)) dx. \end{aligned}$$

Since $0 < \mu < 2 < N$,

$$\begin{aligned} \frac{d}{dt} I(\gamma(t))|_{t=1} &= \frac{N-2}{2} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + \frac{N}{2} \int_{\mathbb{R}^N} f^2(w_0) dx \\ &\quad - \frac{2N-\mu}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(w_0))) G(f(w_0)) dx \\ &= P(w_0) = 0, \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} I(\gamma(t))|_{t<1} &= t^{2N-\mu-1} \left[\frac{N-2}{2t^{N+2-\mu}} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + \frac{N}{t^{N-\mu} 2} \int_{\mathbb{R}^N} f^2(w_0) dx \right. \\ &\quad \left. - \frac{(2N-\mu)}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(w_0))) G(f(w_0)) dx \right] \\ &> t^{2N-\mu-1} P(w_0) = 0 \end{aligned}$$

and

$$\begin{aligned} \frac{d}{dt} I(\gamma(t))|_{t>1} &= t^{2N-\mu-1} \left[\frac{N-2}{2t^{N+2-\mu}} \int_{\mathbb{R}^N} |\nabla w_0|^2 dx + \frac{N}{2t^{N-\mu}} \int_{\mathbb{R}^N} f^2(w_0) dx \right. \\ &\quad \left. - \frac{2N-\mu}{2} \int_{\mathbb{R}^N} (|x|^{-\mu} * G(f(w_0))) G(f(w_0)) dx \right] \\ &< t^{2N-\mu-1} P(w_0) = 0. \end{aligned}$$

Hence $I(\gamma(1)) = \max_{t \geq 0} I(\gamma(t))$. Moreover, obviously, there exists a $t^* > 0$ such that $I(\gamma(t^*)) < 0$. Define $\gamma_1 : [0, 1] \rightarrow E$ by $\gamma_1(t) := \gamma(tt^*)$. Then $\gamma_1 \in \Gamma$. Hence

$$c \leq \sup_{t \in [0,1]} I(\gamma_1(t)) = \sup_{t \in [0,1]} I(\gamma(tt^*)) \leq I(\gamma(1)) = I(w_0),$$

and hence $c \leq c^*$. Therefore, $c = c^*$. This shows v_0 is a ground state solution (1.5), and hence $u_0 = f(v_0)$ is a ground state solution (1.1).

References

- [1] C. O. Alves, D. Cassani, C. Tarsi, M. Yang, Existence and concentration of ground state solutions for a critical nonlocal Schrödinger equation in \mathbb{R}^2 , J. Diff. Equ. 261(2016), 1933-1972.
- [2] C. O. Alves, A. B. Nobrega, M. Yang, Multi-bump solutions for Choquard equation with deepening potential well, Calc. Var. 55:48(2016), 1-28.
- [3] S. Adachi, T. Watanabe, Uniqueness of the ground state solutions of quasilinear Schrödinger equations, Nonlinear Anal. 75(2012), 819-833.
- [4] C. O. Alves, M. Yang, Multiplicity and concentration of solutions for quasilinear Choquard equation, J. Math. Phys. 55(2014), 061502, 1-21.
- [5] J. M. Bezerra do Ó, O. H. Miyagaki, S. H. M. Soares, Soliton solutions for quasilinear Schrödinger equations: the critical exponential case, Nonlinear Anal., 67(2007), 3357-3372.
- [6] J. M. Bezerra do Ó, O. H. Miyagaki, S. H. M. Soares, Soliton solutions for quasilinear Schrödinger equations with critical growth, J. Diff. Equ., 248(2010) 722-744.
- [7] M. Colin, L. Jean, Solutions for a quasilinear Schrödinger equation: a dual approach, Nonlinear Anal. 56(2004)213-226.
- [8] X. D. Fang, A. Szulkin, Multiple solutions for a quasilinear Schrödinger equation, J. Diff. Equ., 254(2013), 2015-2032.

- [9] S. Kurihura, Large-amplitude quasi-solitons in superfluid films, J. Phys. Soc. Jpn. 50(1981), 3262-3267.
- [10] E. H. Lieb, Existence and uniqueness of the minimizing solution of Choquard's non-linear equation, Stud. Appl. Math. 57(2)(1976/1977),93-105.
- [11] X. Liu, J. Liu, Z. Q. Wang, Quasilinear elliptic equations via perturbation method, Proc. Amer. Math. Soc., 141(2013), 253-263.
- [12] A. G. Litvak, A. M. Sergeev, One dimensional collapse of plasma waves, JETP Lett., 27(1978), 517-520.
- [13] E. W. Laedke, K. H. Spatschek, L. Stenflo, Evolution theorem for a class of perturbed envelope soliton solutions, J. Math. Phys. 24(1983), 2764-2769.
- [14] J. Liu, Z. Q. Wang, Soliton solutions for quasilinear Schrödinger Equations I, Proc. Amer. Math. Soc. 131(2003) 441-448.
- [15] J. Liu, Y. Wang, Z. Q. Wang, Soliton solutions for quasilinear Schrödinger equations, II, J. Diff. Equ. 187(2003), 473-493.
- [16] J. Liu, Y. Wang, Z. Q. Wang, Solutions for quasilinear Schrödinger Equations via the Nehari Method, Comm. Part. Diff. Equ. 29(2004)879-892.
- [17] L. M. Moroz, R. Penrose, P. Tod, Spherically-symmetric solutions of the Schrödinger-Newton equations, Classical Quantum Gravity, 15(9)(1998), 2733-2742.
- [18] V. Moroz, J. Van Schaftingen, Ground states of nonlinear Choquard equations: Existence, qualitative properties and decay asymptotics, J. Funct. Anal., 265(2013), 153-184.
- [19] J.M. do Ó, U. Severo, Solitary waves for a class of quasilinear Schrödinger equations in dimension two, Calc. Var. 38(2010), 275-315.
- [20] A. Nakamura, Damping and modification of exciton solitary waves, J. Phys. Soc. Jpn 42(1977), 1824-1835.
- [21] S. Pekar, Untersuchung über die Elektronentheorie der Kristalle, Verlag, Berlin, 1954.

- [22] M. Poppenberg, On the local well posedness of quasi-linear Schrödinger equations in arbitrary space dimension, *J. Diff. Equ.* 172(2001), 83-115.
- [23] M. Porkolab, M.V. Goldman, Upper hybrid solitons and oscillating two-stream instabilities, *Phys. Fluids*, 19(1976), 872-881.
- [24] M. Poppenberg, K. Schmitt, Z. Q. Wang, On the existence of soliton solutions to quasilinear Schrödinger equations, *Calc. Var.* 14(2002), 329-344.
- [25] P. H. Rabinowitz, *Minimax Methods in Critical Point Theory with Application to Differential Equations*, in: CBMS Regional Conf. Ser. in Math, vol.65, American Mathematical Society, Providence, RI, 1986.
- [26] David Ruiz, Gaetano Siciliano, Existence of ground states for a modified nonlinear Schrödinger equation, *Nonlinearity*, 23(2010), 1221-1233.
- [27] D. Ruiz, J. V. Schaftingen, Odd symmetry of least energy nodal solutions for the Choquard equation, *J. Diff. Equ.* 264(2018), 1231-1262.
- [28] M. Willem, *Minimax theorems*, in: *Progr. Nonlinear Differential Equations Appl.*, vol. 24, Birkhauser Boston, Inc. Boston, MA, 1996.
- [29] X. Wu, Multiple solutions for quasilinear Schrödinger equations with a parameter, *J. Diff. Equ.* 256(2014), 2619-2632.
- [30] W. Zhang, X. Wu, Nodal solutions for a fractional Choquard equation, *J. Math. Anal. Appl.* 464(2018), 1167-1183.
- [31] W. Zhang, X. Wu, Existence, multiplicity, and concentration of positive solutions for a quasilinear Choquard equation with critical exponent, *J. Math. Phys.* 60(2019), 051501, 1-19.
- [32] Zhong Chen-Kui, Fan Xian-Ling, Chen Wen-yuan, *Introduction of Non-linear Functional Analysis*, Lanzhou University Publishing House, (1998).
- [33] W. M. Zou, M. Schechter, *Critical Point Theory and its Applications*, Springer, New York, (2006).