

# An eigenvalue problem for nonlinear Schrödinger-Poisson system with steep potential well

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## Abstract

In this paper, we study an eigenvalue problem for Schrödinger-Poisson system with indefinite nonlinearity and potential well as follows:

$$\begin{cases} -\Delta u + \mu V(x)u + K(x)\phi u = \lambda f(x)u + g(x)|u|^{p-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases}$$

where  $4 \leq p < 6$ , the parameters  $\mu, \lambda > 0$ ,  $V \in C(\mathbb{R}^3)$  is a potential well, and the functions  $f \in L^{\frac{3}{2}}(\mathbb{R}^3)$  and  $g \in L^\infty(\mathbb{R}^3)$  are allowed to be sign-changing. It is well known that such a system with the potential being positive constant has two positive solutions when  $\lim_{|x| \rightarrow \infty} g(x) = g_\infty < 0$ ,  $K = 0$  in the set  $\{x \in \mathbb{R}^3 : g(x) = 0\}$  and  $\lambda > \lambda_1(f)$  with near  $\lambda_1(f)$ , where  $\lambda_1(f)$  is the first eigenvalue of  $-\Delta + id$  in  $H^1(\mathbb{R}^3)$  (see e.g. Huang et al., J. Differential Equations **255**, 2463 (2013)). The main purpose is to obtain the existence and multiplicity of positive solutions without the above assumptions for  $g$  and  $K$ . The results are obtained via variational method and steep potential. Furthermore, we also consider the concentration of solutions as  $\mu \rightarrow \infty$ .

**Keywords.** Schrödinger-Poisson system, eigenvalue problem, steep potential well, mountain pass theory.

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## 1 Introduction

In the paper, we are concerned the following Schrödinger-Poisson system

$$\begin{cases} -\Delta u + V(x)u + K(x)\phi u = h(x, u) & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3. \end{cases} \quad (1.1)$$

This system has been first derived from Benci and Fortunato [1] in order to study the stationary solutions of Schrödinger equations coupled with Maxwell equations. It describes the interaction of a charged particle with its own electrostatic field. The unknown functions  $u$  and  $\phi$  represent the wave functions associated to the particle and electric potential, respectively. The functions  $V$  and  $K$  denote an external potential and nonnegative density charge, respectively. The presence of the nonlinear term  $h(x, u)$  simulates the interaction effect among many particles. We refer the reader to [1] and [2] for more details.

In recent years, system (1.1) has been widely studied under variant assumptions on  $V, K$  and  $h$ . See, for example, [1–7] for the autonomous case and [4, 8–20] for the non-autonomous case. In

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particular, many papers have been devoted to the nonlinear term  $h(x, u) = g(x)|u|^{p-2}u$ ,  $2 < p < 6$ . For which some assumptions on  $V, K$  and  $g$  are considered in order to overcome the lack of compactness of the embedding of  $H^1(\mathbb{R}^3)$  into  $L^r(\mathbb{R}^3)$ ,  $2 \leq r < 6$ , since system (1.1) is on the whole space  $\mathbb{R}^3$ .

In the work [9], Cerami and Vaira studied the following Schrödinger-Poisson system

$$\begin{cases} -\Delta u + u + \lambda K(x)\phi u = g(x)|u|^{p-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (1.2)$$

where  $\lambda \equiv 1$ ,  $4 < p < 6$ ,  $g$  and  $K$  are nonnegative real functions,

$$\lim_{|x| \rightarrow \infty} g(x) = g_\infty > 0, \quad \lim_{|x| \rightarrow \infty} K(x) = K_\infty = 0 \quad \text{and} \quad K \in L^2(\mathbb{R}^3).$$

They obtained the existence of bound and ground state positive solutions by using the Nehari manifold method and establishing a global compactness lemma to overcome the lack of compactness. Later, Vaira [10] considered system (1.2) with  $\lambda \in \mathbb{R}$ ,  $g$  and  $K$  being nonnegative functions,

$$\lim_{|x| \rightarrow \infty} g(x) = g_\infty > 0, \quad \lim_{|x| \rightarrow \infty} K(x) = K_\infty > 0 \quad \text{and} \quad K - K_\infty \in L^2(\mathbb{R}^3).$$

For which the positive ground state solutions were obtained in the cases where  $2 < p < 6$  if  $\lambda < 0$  and  $4 < p < 6$  if  $\lambda > 0$  by using the Nehari manifold method.

Recently, Huang et al. [11] studied Schrödinger-Poisson system with indefinite nonlinearity, namely

$$\begin{cases} -\Delta u + u + K(x)\phi u = \lambda f(x)u + g(x)|u|^{p-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (1.3)$$

where  $4 < p < 6$ ,  $K$  and  $f$  are nonnegative functions,  $K \in L^2(\mathbb{R}^3)$ ,  $f \in L^{\frac{3}{2}}(\mathbb{R}^3)$ ,  $g \in C(\mathbb{R}^3)$  which changes sign in  $\mathbb{R}^3$ ,

$$\lim_{|x| \rightarrow \infty} g(x) = g_\infty < 0 \quad \text{and} \quad K = 0 \text{ a.e. in the set } \{x \in \mathbb{R}^3 : g(x) = 0\}. \quad (1.4)$$

In which they proved the existence of two positive solutions in  $\lambda > \lambda_1(f)$  and near  $\lambda_1(f)$ , where  $\lambda_1(f)$  is the first eigenvalue of  $-\Delta u + u = \lambda f(x)u$  in  $H^1(\mathbb{R}^3)$ , whose corresponding eigenfunction is denoted by  $\phi_1$ . Besides the condition about  $g$  in (1.4) is different from [9, 10], it is worth noting that they do not need the sign condition  $\int_{\mathbb{R}^3} g(x)\phi_1^p dx < 0$ , which has been shown to be a necessary condition to the existence of positive solutions for semilinear elliptic equations with indefinite nonlinearity, see [21–23]. Later, Chen [12] studied the case of  $K \in L^\infty(\mathbb{R}^3)$  (more precisely,  $K(x) \equiv 1$ ) for system (1.3). Since the condition about  $K$  in (1.4) is not established now, the author assume the additional condition

$$|\{x \in \mathbb{R}^3 : g(x) = 0\}| = 0 \quad (1.5)$$

to ensure that the existence of two positive solutions can be obtained in  $\lambda > \lambda_1(f)$  and near  $\lambda_1(f)$ .

Very recently, Shen and Han [16] studied system (1.3) with  $p = 4$ ,  $f$  and  $K$  being nonnegative functions,  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$ ,  $f \in L^{\frac{3}{2}}(\mathbb{R}^3)$ ,  $g \in C(\mathbb{R}^3)$  to be sign-changing and  $\lim_{|x| \rightarrow \infty} g(x) = g_\infty < 0$ . The authors do not need the condition about  $K$  in (1.4) that  $K = 0$  a.e. in  $\{x \in \mathbb{R}^3 : g(x) = 0\}$ , and assume the following condition:

$$\int_{\mathbb{R}^3} g(x)\phi_1^4(x)dx - \frac{1}{4\pi} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \frac{K(x)K(y)\phi_1^2(x)\phi_1^2(y)}{|x-y|} dydx < 0. \quad (1.6)$$

For which the existence of two positive solutions in  $\lambda > \lambda_1(f)$  and near  $\lambda_1(f)$  was proved via the Nehari manifold method.

In recent year, many papers study Schrödinger-Poisson system with steep potential well, namely

$$\begin{cases} -\Delta u + \mu V(x)u + K(x)\phi u = h(x, u) & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (1.7)$$

where  $\mu > 0$  is a parameter and the potential  $V$  satisfies the following conditions:

- (V1)  $V$  is a non-negative continuous function in  $\mathbb{R}^3$ .
- (V2) There exists  $c_0 > 0$  such that the set  $\{V < c_0\} := \{x \in \mathbb{R}^3 : V(x) < c_0\}$  is nonempty and has finite positive measure.
- (V3)  $\Omega = \text{int}\{x \in \mathbb{R}^3 : V(x) = 0\}$  is nonempty bounded domain and has a smooth boundary with  $\bar{\Omega} = \{x \in \mathbb{R}^3 : V(x) = 0\}$ .

The conditions (V1)-(V3) were first introduced by Bartsch and Wang [24] in the study of nonlinear Schrödinger equations. Here  $\mu V$  represents a potential well whose depth is controlled by the parameter  $\mu$ , and it is called a steep potential well if  $\mu$  large. Later, steep potential well was first applied into Schrödinger-Poisson system in [25], and then has been widely applied to the study of Schrödinger-Poisson system. We refer the reader to [25–32] and the references therein.

In the paper [26], Zhao et al. studied system (1.7) with  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$ ,  $K \geq 0$  and  $h(x, u) = |u|^{p-2}u$ . In which the existence of nontrivial solution and concentration results are obtained via variational methods when  $3 < p < 6$ . In particular, the potential  $V$  is allowed to be sign-changing for the case  $4 < p < 6$ . Later, Du et al. [27] considered system (1.7) with  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$ ,  $K \geq 0$  and  $h(x, u) = g(x)F(u)$ , where  $g$  is a positive bounded function and  $F$  is either asymptotically linear or asymptotically 3-linear at infinity. The existence and asymptotic behavior of solutions are proved via variational methods. As far as I know, system (1.7) with indefinite nonlinearity has not been studied.

Motivated by the above works [11, 12, 16, 26, 27], we consider Schrödinger-Poisson system in the following form:

$$\begin{cases} -\Delta u + \mu V(x)u + K(x)\phi u = \lambda f(x)u + g(x)|u|^{p-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (SP_{\mu, \lambda})$$

where  $4 \leq p < 6$ , the parameters  $\mu, \lambda > 0$ , and the potential  $V$  satisfies conditions (V1)-(V3). The aim of this paper is to answer the following questions:

- (I) Does system  $(SP_{\mu, \lambda})$  with  $4 \leq p < 6$  still admit multiple positive solutions without the conditions (1.4) and (1.5)?
- (II) Can we consider system  $(SP_{\mu, \lambda})$  with the weight functions  $f$  and  $g$  to be sign-changing?
- (III) Can we deal with the case where  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$  as in [16, 26, 27]?

In order to give the definite answer to questions (I)-(III), we assume that the functions  $K$ ,  $f$  and  $g$  satisfy the following conditions:

- (F)  $f \in L^{\frac{3}{2}}(\mathbb{R}^3) \cap L^\infty(\mathbb{R}^3)$  and  $|\{x \in \Omega : f(x) > 0\}| > 0$ ;
- (G)  $g \in L^\infty(\mathbb{R}^3)$  which changes sign in  $\Omega$ ;
- (K1)  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$ ,  $K \geq 0$  and  $K \not\equiv 0$  in  $\Omega$ ;
- (K2)  $K > 0$  a.e. in the set  $\{x \in \Omega : g(x) = 0\}$ .

**Remark 1.1.** If  $f$  is bounded in  $\overline{\Omega}$  and  $|\{x \in \Omega : f(x) > 0\}| > 0$ , then there exists a sequence of eigenvalues  $\{\lambda_n(f_{\overline{\Omega}})\}$  of the problem

$$-\Delta u = \lambda f_{\overline{\Omega}}(x)u, \quad u \in H_0^1(\Omega), \quad (1.8)$$

with  $0 < \lambda_1(f_{\overline{\Omega}}) < \lambda_2(f_{\overline{\Omega}}) \leq \dots$  and each eigenvalue being of finite multiplicity, where  $f_{\overline{\Omega}}$  is a restriction of  $f$  on  $\overline{\Omega}$ . Denote by  $e_1$  the positive principal eigenfunction with

$$\lambda_1(f_{\overline{\Omega}}) = \int_{\Omega} |\nabla e_1|^2 dx = \inf \left\{ \int_{\Omega} |\nabla u|^2 dx : u \in H_0^1(\Omega), \int_{\Omega} f_{\overline{\Omega}} u^2 dx = 1 \right\}.$$

Moreover, we have

$$\lambda_2(f_{\overline{\Omega}}) = \inf \left\{ \int_{\Omega} |\nabla u|^2 dx : u \in H_0^1(\Omega), \int_{\Omega} f_{\overline{\Omega}} u^2 dx = 1, \int_{\Omega} \nabla u \nabla e_1 dx = 0 \right\}.$$

In the present paper, we shall prove the existence and multiplicity of positive solutions for system  $(SP_{\mu,\lambda})$  via the mountain pass theory. In order to show the mountain pass geometry of the related functional, we follow the argument in [23] to decompose each  $u \in X$  (defined later) as  $u = te_{1,\mu} + w$ , where  $w \in \{\text{span}\{e_{1,\mu}\}\}^\perp$  and  $e_{1,\mu}$  is the positive principal eigenfunction corresponding to the positive principal eigenvalue  $\lambda_{1,\mu}(f)$  of the problem

$$-\Delta u + \mu V(x)u = \lambda f(x)u.$$

Then by the approximation estimate

$$\lambda_{1,\mu}(f) \rightarrow \lambda_1^-(f_{\overline{\Omega}}) \text{ as } \mu \rightarrow \infty,$$

we can deduce the mountain pass geometry of the related functional in  $\lambda > \lambda_1(f_{\overline{\Omega}})$  and near  $\lambda_1(f_{\overline{\Omega}})$ . Moreover, we apply the concentration compactness principle [33] and a variant version of the steep well method [24] to overcome the difficulty that the lack of compactness. Furthermore, we also consider the concentration of solutions.

We state the main results.

**Theorem 1.2.** Suppose that  $4 \leq p < 6$  and conditions (V1)-(V3), (F), (G), (K1) and (K2) hold. In addition, for  $p = 4$ , we assume the following:

(K3)  $K = 0$  a.e. in the set  $\{x \in \Omega : g(x) > 0\}$ ;

$$(D1) \quad \|g\|_{L^\infty} |\{V < c_0\}|^{\frac{1}{3}} \leq \frac{S^4}{64\pi\lambda_1(f_{\overline{\Omega}})^2} \int_{\Omega} \int_{\Omega} \frac{K(x)K(y)e_1^2(x)e_1^2(y)}{|x-y|} dy dx.$$

Then we have the following results.

(i) For every  $0 < \lambda \leq \lambda_1(f_{\overline{\Omega}})$ , system  $(SP_{\mu,\lambda})$  has a positive solution for  $\mu > 0$  large enough.

(ii) There exists  $\delta_0 > 0$  such that for every  $\lambda_1(f_{\overline{\Omega}}) < \lambda < \lambda_1(f_{\overline{\Omega}}) + \delta_0$ , system  $(SP_{\mu,\lambda})$  has two positive solutions for  $\mu > 0$  large enough.

**Theorem 1.3.** Let  $(u_\mu, \phi_{u_\mu})$  be the solutions obtained in Theorem 1.2. Then we have  $u_\mu \rightarrow u_\infty$  in  $X$  as  $\mu \rightarrow \infty$ , where  $u_\infty \in H_0^1(\Omega)$  is a positive solution of equation:

$$\begin{cases} -\Delta u + K(x) \left( \int_{\Omega} \frac{K(y)u^2(y)}{4\pi|x-y|} dy \right) u = \lambda f(x)u + g(x)|u|^{p-2}u & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases} \quad (S_{\infty,\lambda})$$

The outline of the paper is as follows. In next section, we introduce the variational setting and some basic estimates. In section 3, we consider the eigenvalue problem  $-\Delta u + \mu V(x)u = \lambda f(x)u$ . In section 4, we prove the corresponding energy functional having the mountain pass geometry and satisfying the Palais-Smale condition. In section 5, we prove Theorem 1.2. In last section, we give the proof of Theorem 1.3.

## 2 Notations and variational setting

We denote the  $L^q(\mathbb{R}^3)$ -norm for  $1 \leq q \leq \infty$  by  $\|\cdot\|_{L^q}$ . As we take a subsequence for a sequence  $\{u_n\}$ , we shall still use  $\{u_n\}$  to denote it. We use  $o(1)$  to denote a quantity that depends on  $n \in \mathbb{N}$  and goes to zero as  $n \rightarrow \infty$ . Let  $S$  be the best Sobolev constant for the embedding of  $D^{1,2}(\mathbb{R}^3)$  in  $L^6(\mathbb{R}^3)$ , where  $D^{1,2}(\mathbb{R}^3)$  is the completion of  $C_0^\infty(\mathbb{R}^3)$  with respect to the norm  $\|u\|_{D^{1,2}}^2 = \|\nabla u\|_{L^2}^2$ .

Next, we consider the variational setting for system  $(SP_{\mu,\lambda})$ . It is well known that for any  $K \in L^2(\mathbb{R}^3) \cup L^\infty(\mathbb{R}^3)$  and  $u \in H^1(\mathbb{R}^3)$ , the Lax-Milgram theorem implies that there exists a unique  $\phi_u \in D^{1,2}(\mathbb{R}^3)$  such that

$$\int_{\mathbb{R}^3} \nabla \phi_u \nabla v dx = \int_{\mathbb{R}^3} K(x) u^2 v dx \quad \text{for all } v \in D^{1,2}(\mathbb{R}^3). \quad (2.1)$$

That is,  $\phi_u$  is the weak solution of  $-\Delta \phi = K(x) u^2$ . Moreover,  $\phi_u$  can be represented by

$$\phi_u(x) = \frac{1}{4\pi} \int_{\mathbb{R}^3} \frac{K(y) u^2(y)}{|x-y|} dy. \quad (2.2)$$

Therefore, we can transform system  $(SP_{\mu,\lambda})$  into a nonlinear Schrödinger equation with a non-local term as follows:

$$-\Delta u + \mu V(x) u + K(x) \phi_u u = \lambda f(x) u + g(x) |u|^{p-2} u \quad (S_{\mu,\lambda})$$

with  $\phi_u$  as in (2.2). Next, we set the space

$$X = \left\{ u \in H^1(\mathbb{R}^3) : \int_{\mathbb{R}^3} V u^2 dx < \infty \right\}$$

with the following inner product and norm

$$\langle u, v \rangle_\mu = \int_{\mathbb{R}^3} (\nabla u \nabla v + \mu V u v) dx, \quad \|u\|_\mu = \langle u, u \rangle_\mu^{1/2},$$

for  $\mu > 0$ . By condition (V2), the Hölder and Sobolev inequalities, we deduce that

$$\int_{\mathbb{R}^3} u^2 dx = \int_{\{V \geq c_0\}} u^2 dx + \int_{\{V < c_0\}} u^2 dx \leq \frac{1}{\mu c_0} \int_{\mathbb{R}^3} \mu V u^2 dx + \frac{|\{V < c_0\}|^{\frac{2}{3}}}{S^2} \|u\|_{D^{1,2}}^2,$$

which implies that the embedding  $X \hookrightarrow H^1(\mathbb{R}^3)$  is continuous. Furthermore, for  $2 \leq r \leq 6$  and  $\mu \geq \mu_0 := S^2 \left( c_0 |\{V < c_0\}|^{\frac{2}{3}} \right)^{-1}$ , we have

$$\begin{aligned} \int_{\mathbb{R}^3} |u|^r dx &\leq \left( \int_{\mathbb{R}^3} u^2 dx \right)^{\frac{6-r}{4}} \left( \int_{\mathbb{R}^3} u^6 dx \right)^{\frac{r-2}{4}} \\ &\leq \left( \frac{1}{\mu c_0} \int_{\mathbb{R}^3} \mu V u^2 dx + \frac{|\{V < c_0\}|^{\frac{2}{3}}}{S^2} \|u\|_{D^{1,2}}^2 \right)^{\frac{6-r}{4}} \left( \frac{\|u\|_{D^{1,2}}^6}{S^6} \right)^{\frac{r-2}{4}} \\ &\leq |\{V < c_0\}|^{\frac{6-r}{6}} S^{-r} \|u\|_\mu^r. \end{aligned} \quad (2.3)$$

Now, for Eq.  $(S_{\mu,\lambda})$ , we set the energy functional  $J_{\mu,\lambda} : X \rightarrow \mathbb{R}$  which is defined by

$$J_{\mu,\lambda}(u) = \frac{1}{2} \|u\|_\mu^2 + \frac{1}{4} \int_{\mathbb{R}^3} K \phi_u u^2 dx - \frac{\lambda}{2} \int_{\mathbb{R}^3} f u^2 dx - \frac{1}{p} \int_{\mathbb{R}^3} g |u|^p dx.$$

The functional  $J_{\mu,\lambda}$  is a  $C^1$  functional with the derivative given by

$$\langle J'_{\mu,\lambda}(u), \varphi \rangle = \int_{\mathbb{R}^3} (\nabla u \nabla \varphi + \mu V u \varphi) dx + \int_{\mathbb{R}^3} K \phi_u u \varphi dx - \lambda \int_{\mathbb{R}^3} f u \varphi dx - \int_{\mathbb{R}^3} g |u|^{p-2} u \varphi dx$$

for all  $\varphi \in X$ , where  $J'_{\mu,\lambda}$  denotes the Fréchet derivative of  $J_{\mu,\lambda}$ . One can see that the critical points of  $J_{\mu,\lambda}$  are corresponding to the solutions of Eq.  $(S_{\mu,\lambda})$ . Therefore, we conclude that  $u$  is a critical point of  $J_{\mu,\lambda}$  if and only if  $(u, \phi_u)$  solves system  $(SP_{\mu,\lambda})$ .

Let us define the operator  $\Phi : H^1(\mathbb{R}^3) \rightarrow D^{1,2}(\mathbb{R}^3)$  and the functional  $N : H^1(\mathbb{R}^3) \rightarrow \mathbb{R}$  by  $\Phi(u) := \phi_u$  and

$$N(u) := \int_{\mathbb{R}^3} K \phi_u u^2 dx = \frac{1}{4\pi} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \frac{K(x)K(y)u^2(x)u^2(y)}{|x-y|} dy dx \quad (2.4)$$

respectively. By (2.1) and (2.3), we have the following estimates for  $\mu \geq \mu_0$ :

$$\|\phi_u\|_{D^{1,2}}^2 = N(u) \leq \begin{cases} S^{-2} \|K\|_{L^2}^2 \|u\|_{L^6}^4 \leq S^{-6} \|K\|_{L^2}^2 \|u\|_{\mu}^4, & \text{if } K \in L^2(\mathbb{R}^3), \\ S^{-2} \|K\|_{L^\infty}^2 \|u\|_{L^{\frac{12}{5}}}^4 \leq S^{-6} \{V < c_0\} \|K\|_{L^\infty}^2 \|u\|_{\mu}^4, & \text{if } K \in L^\infty(\mathbb{R}^3). \end{cases} \quad (2.5)$$

Next, we state some useful properties.

**Proposition 2.1** ([9], Lemma 2.1). (i)  $\Phi$  is continuous and  $\Phi(tu) = t^2 \Phi(u)$  for all  $t \in \mathbb{R}$ .  
(ii)  $\Phi$  maps bounded sets into bounded sets.

**Proposition 2.2.** Assume that the sequence  $\{u_n\} \subset H^1(\mathbb{R}^3)$  is bounded. Then there exist a subsequence  $\{u_n\}$  (still denote by  $\{u_n\}$ ) and  $u \in H^1(\mathbb{R}^3)$  such that

(i) for  $K \in L^2(\mathbb{R}^3)$ , we have  $N(u_n) = N(u) + o(1)$  and  $N'(u_n) = N'(u) + o(1)$  in  $H^{-1}(\mathbb{R}^3)$ ;  
(ii) for  $K \in L^\infty(\mathbb{R}^3)$ , we have  $N(u_n - u) = N(u_n) - N(u) + o(1)$  and  $N'(u_n) = N'(u) + o(1)$  in  $H^{-1}(\mathbb{R}^3)$ .

**Proposition 2.3** ([34], Lemma 2.13). If  $f \in L^{\frac{3}{2}}(\mathbb{R}^3)$ , the functional  $u \mapsto \int_{\mathbb{R}^3} f u^2 dx$  is weakly continuous on  $H^1(\mathbb{R}^3)$ .

**Remark 2.4.** According to the fact that the embedding  $X \hookrightarrow H^1(\mathbb{R}^3)$  is continuous, the above all properties are still valid as the space  $H^1(\mathbb{R}^3)$  replaced by  $X$ .

We omit the proofs of Propositions 2.1-2.3 and the readers are referred to [9, 20, 26, 34].

### 3 Eigenvalue problems

In this section, we study the following eigenvalue problem

$$-\Delta u + \mu V(x)u = \lambda f(x)u \quad \text{in } X. \quad (3.1)$$

In order to find the positive principal eigenvalue of (3.1) we solve the following minimization problem:

$$\min \left\{ \int_{\mathbb{R}^3} (|\nabla u|^2 + \mu V u^2) dx : u \in X, \int_{\mathbb{R}^3} f u^2 dx = 1 \right\}. \quad (3.2)$$

Let

$$\lambda_{1,\mu}(f) := \inf \left\{ \int_{\mathbb{R}^3} (|\nabla u|^2 + \mu V u^2) dx : u \in X, \int_{\mathbb{R}^3} f u^2 dx = 1 \right\}.$$

By condition (F), we deduce that

$$\frac{\int_{\mathbb{R}^3} (|\nabla u|^2 + \mu V u^2) dx}{\int_{\mathbb{R}^3} f u^2 dx} \geq \frac{\|u\|_{D^{1,2}}^2}{\|f\|_{L^{\frac{3}{2}}} S^{-2} \|u\|_{D^{1,2}}^2} = \frac{S^2}{\|f\|_{L^{\frac{3}{2}}}} > 0,$$

which implies that  $\lambda_{1,\mu}(f) \geq S^2 \|f\|_{L^{\frac{3}{2}}}^{-1} > 0$  for all  $\mu > 0$ . Moreover, by condition (V3),

$$\inf_{u \in X \setminus \{0\}} \frac{\int_{\mathbb{R}^3} |\nabla u|^2 + \mu V u^2 dx}{\int_{\mathbb{R}^3} f u^2 dx} \leq \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\int_{\mathbb{R}^3} |\nabla u|^2 + \mu V u^2 dx}{\int_{\mathbb{R}^3} f u^2 dx} = \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^2 dx}{\int_{\Omega} f_{\Omega} u^2 dx},$$

which indicates that  $\lambda_{1,\mu}(f) \leq \lambda_1(f_{\Omega})$  for all  $\mu > 0$ , where  $\lambda_1(f_{\Omega})$  is the positive principal eigenvalue of (1.8). Next, we state the useful lemma to solve problem (3.2).

**Lemma 3.1** ([35], Lemma 2.4). *Let  $\mu_n \rightarrow \infty$  as  $n \rightarrow \infty$  and  $\{v_n\} \subset X$  with  $\|v_n\|_{\mu_n} \leq C_0$  for some  $C_0 > 0$ . Then there exists a subsequence  $\{v_n\}$  and  $v_0 \in H_0^1(\Omega)$  such that  $v_n \rightarrow v_0$  in  $X$  and  $v_n \rightarrow v_0$  in  $L^r(\mathbb{R}^3)$  for all  $2 \leq r < 6$ . Furthermore, we have  $N(v_n) \rightarrow N(v_0)$  as  $n \rightarrow \infty$ .*

**Lemma 3.2.** *For each  $\mu > 0$ , problem (3.2) has a positive solution  $e_{1,\mu} \in X$  such that*

$$\lambda_{1,\mu}(f) = \int_{\mathbb{R}^3} |\nabla e_{1,\mu}|^2 + \mu V e_{1,\mu}^2 dx.$$

Furthermore, we have

- (i)  $\lambda_{1,\mu}(f)$  is simple and principal eigenvalue of Eq. (3.1) and  $e_{1,\mu}$  is a corresponding eigenfunction.
- (ii)  $\lambda_{1,\mu}(f) < \lambda_1(f_{\Omega})$ ,  $\lambda_{1,\mu}(f) \rightarrow \lambda_1^-(f_{\Omega})$  and  $e_{1,\mu} \rightarrow e_1$  in  $X$  as  $\mu \rightarrow \infty$ , where  $e_1$  is the positive principal eigenfunction corresponding to  $\lambda_1(f_{\Omega})$ .

*Proof.* Some results has been proved in [35]. For the reader's convenience, we give the proof in detail here. Let  $\{u_n\} \subset X$  be a minimizing sequence of problem (3.2). Clearly,  $\{u_n\}$  is bounded and thus there exist a subsequence  $\{u_n\}$  and  $e_{1,\mu} \in X$  such that

$$u_n \rightharpoonup e_{1,\mu} \quad \text{in } X; \tag{3.3}$$

$$u_n \rightarrow e_{1,\mu} \quad \text{in } L_{loc}^r(\mathbb{R}^3) \text{ for } 2 \leq r < 6. \tag{3.4}$$

Then by Proposition 2.3 and (3.3), we have

$$\int_{\mathbb{R}^3} f e_{1,\mu}^2 dx = \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f u_n^2 dx = 1. \tag{3.5}$$

We next show that  $u_n \rightarrow e_{1,\mu}$  in  $X$ . If it is false, then we have

$$\int_{\mathbb{R}^3} |\nabla e_{1,\mu}|^2 + \mu V e_{1,\mu}^2 dx < \liminf_{n \rightarrow \infty} \|u_n\|_{\mu}^2 = \lambda_{1,\mu}(f),$$

which is impossible according to the definition of  $\lambda_{1,\mu}(f)$  and (3.5). Hence  $u_n \rightarrow e_{1,\mu}$  in  $X$  and  $\lambda_{1,\mu}(f) = \|e_{1,\mu}\|_{\mu}^2$ . Since  $e_{1,\mu} \in X$  and  $\|e_{1,\mu}\|_{\mu}^2 = \|e_{1,\mu}\|_{\mu}^2 = \lambda_{1,\mu}(f)$ , we may assume that  $e_{1,\mu} \geq 0$ .

Let

$$F(t) := \frac{\|e_{1,\mu} + t\varphi\|_{\mu}^2}{\int_{\mathbb{R}^3} f (e_{1,\mu} + t\varphi)^2 dx}$$

for any  $\varphi \in C_0^\infty(\mathbb{R}^3)$ . Then we have  $F'(0) = 0$  since  $F$  has a minimizer at  $t = 0$ . By  $F'(0) = 0$ , we deduce that

$$\int_{\mathbb{R}^3} (\nabla e_{1,\mu} \nabla \varphi + \mu V e_{1,\mu} \varphi) dx = \lambda_{1,\mu}(f) \int_{\mathbb{R}^3} f e_{1,\mu} \varphi dx,$$

which indicates that  $e_{1,\mu}$  is an eigenfunction of (3.1) corresponding to the eigenvalue  $\lambda_{1,\mu}(f)$ . According to the strong maximum principle, we may assume that  $e_{1,\mu} > 0$ .

Here we show that  $\lambda_{1,\mu}(f)$  is simple. By Sobolev embedding and conditions (V1) and (F), we conclude that any corresponding eigenfunctions belong to  $W_{loc}^{2,q}(\mathbb{R}^3) \cap C_{loc}(\mathbb{R}^3)$  for  $2 \leq q \leq 6$ . Suppose that  $\lambda_{1,\mu}(f)$  is not simple. Then there exist an eigenfunction  $v \in X$  and  $d > 0$  such that  $u = v - de_{1,\mu}$  both takes positive and negative values. Thus, there exists  $x_0 \in \mathbb{R}^3$  such that  $u(x_0) = 0$  since  $v, e_{1,\mu} \in C_{loc}(\mathbb{R}^3)$ . Moreover, we may obtain

$$-\Delta|u| + \mu V|u| = \lambda_{1,\mu}(f)f|u|.$$

By the strong maximum principle, we have  $|u| > 0$ , which contradicts  $u(x_0) = 0$ . Hence  $\lambda_{1,\mu}(f)$  is simple.

Next, we show that  $\lambda_{1,\mu}(f) < \lambda_1(f_{\bar{\Omega}})$ . Arguing by contradiction, we assume that  $\lambda_{1,\mu}(f) = \lambda_1(f_{\bar{\Omega}})$ . Then one can obtain  $e_1$  is also an eigenfunction of (3.1), which is a contradiction due to Harnack inequality. Thus, we conclude that  $\lambda_{1,\mu}(f) < \lambda_1(f_{\bar{\Omega}})$ .

For any sequence  $\mu_n \rightarrow \infty$ , let  $v_n = e_{1,\mu_n}$  be the minimizer of  $\lambda_{1,\mu_n}(f)$ , then we have

$$\int_{\mathbb{R}^3} f v_n^2 dx = 1 \quad \text{and} \quad \lambda_{1,\mu_n}(f) = \|v_n\|_{\mu_n}^2 < \lambda_1(f_{\bar{\Omega}}).$$

By Lemma 3.1, there exist a subsequence  $\{v_n\}$  and  $v_0 \in H_0^1(\Omega)$  such that  $v_n \rightharpoonup v_0$  in  $X$  and  $v_n \rightarrow v_0$  in  $L^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ . Then by Proposition 2.3, we have

$$\int_{\Omega} f v_0^2 dx = \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f v_n^2 dx = 1.$$

Moreover,

$$\int_{\Omega} |\nabla v_0|^2 dx = \int_{\mathbb{R}^3} |\nabla v_0|^2 + V v_0^2 dx \leq \liminf_{n \rightarrow \infty} \|v_n\|_{\mu_n}^2 \leq \lambda_1(f_{\bar{\Omega}}).$$

By Remark 1.1, we conclude that  $v_0 = e_1$  and  $v_n \rightarrow e_1$  in  $X$ . It is easy to deduce that  $\lambda_{1,\mu_1}(f) \leq \lambda_{1,\mu_2}(f)$  for  $\mu_1 < \mu_2$ , and thus we can conclude that  $e_{1,\mu} \rightarrow e_1$  in  $X$ . This completes the proof.  $\square$

In order to find the other positive eigenvalues of (3.1) we solve the following problem

$$\min \left\{ \int_{\mathbb{R}^3} |\nabla u|^2 + \mu V u^2 dx : u \in X, \int_{\mathbb{R}^3} f u^2 dx = 1, \langle u, e_{i,\mu} \rangle_{\mu} = 0 \text{ for all } 1 \leq i \leq n-1 \right\}. \quad (3.6)$$

**Lemma 3.3.** *For each  $\mu > 0$  and  $n \geq 2$ , problem (3.6) has a solution  $e_{n,\mu} \in X$ . Furthermore,  $e_{n,\mu}$  is an eigenfunction of (3.1) corresponding to the eigenvalue  $\lambda_{n,\mu}(f) := \|e_{n,\mu}\|_{\mu}^2$  and*

$$\lambda_{2,\mu}(f) > \frac{\lambda_1(f_{\bar{\Omega}}) + \lambda_2(f_{\bar{\Omega}})}{2} \quad \text{for all } \mu \text{ large enough.}$$

*Proof.* The existence of  $e_{n,\mu}$  is proved as in Lemma 3.2. However, we cannot assume that  $e_{n,\mu} \geq 0$  because of the condition  $\langle e_{n,\mu}, e_{i,\mu} \rangle_{\mu} = 0$ . Let  $v \in X$  satisfying  $\langle v, e_{i,\mu} \rangle_{\mu} = 0$  for all  $1 \leq i \leq n-1$  and

$$F(t) = \frac{\|e_{n,\mu} + tv\|_{\mu}^2}{\int_{\mathbb{R}^3} f(e_{n,\mu} + tv)^2 dx}.$$

Then  $F'(0) = 0$  implies that

$$\int_{\mathbb{R}^3} (\nabla e_{n,\mu} \nabla v + \mu V e_{n,\mu} v) dx = \lambda_{n,\mu}(f) \int_{\mathbb{R}^3} f e_{n,\mu} v dx. \quad (3.7)$$



Let  $\varphi \in C_0^\infty(\mathbb{R}^3)$  and  $w = \varphi - \sum_{i=1}^{n-1} c_i e_{i,\mu}$ , where  $c_i = \frac{\langle \varphi, e_{i,\mu} \rangle_\mu}{\langle e_{i,\mu}, e_{i,\mu} \rangle_\mu}$ . Then  $\langle w, e_{i,\mu} \rangle_\mu = 0$  for all  $1 \leq i \leq n-1$ , and thus by (3.7),

$$\int_{\mathbb{R}^3} \nabla e_{n,\mu} \nabla \varphi + \mu V e_{n,\mu} \varphi dx = \int_{\mathbb{R}^3} \nabla e_{n,\mu} \nabla w + \mu V e_{n,\mu} w dx = \lambda_{n,\mu}(f) \int_{\mathbb{R}^3} f e_{n,\mu} w dx = \lambda_{n,\mu}(f) \int_{\mathbb{R}^3} f e_{n,\mu} \varphi dx.$$

This indicates that  $e_{n,\mu}$  is an eigenfunction of (3.1) corresponding to the eigenvalue  $\lambda_{n,\mu}(f)$ .

We now show that  $\lambda_{2,\mu}(f) > \frac{1}{2}(\lambda_1(f_{\bar{\Omega}}) + \lambda_2(f_{\bar{\Omega}}))$  for all  $\mu$  large enough. Suppose on the contrary. Then there exists a sequence  $\{\mu_n\}$  with  $\mu_n \rightarrow \infty$  such that

$$\lambda_{2,\mu_n}(f) \leq \frac{\lambda_1(f_{\bar{\Omega}}) + \lambda_2(f_{\bar{\Omega}})}{2}.$$

Let  $v_n = e_{2,\mu_n}$  be the minimizer of  $\lambda_{2,\mu_n}(f)$ , then  $\int_{\mathbb{R}^3} f v_n^2 dx = 1$ ,  $\langle v_n, e_{1,\mu_n} \rangle_{\mu_n} = 0$  and

$$\lambda_{2,\mu_n}(f) = \|v_n\|_{\mu_n}^2 \leq \frac{\lambda_1(f_{\bar{\Omega}}) + \lambda_2(f_{\bar{\Omega}})}{2}. \quad (3.8)$$

By Lemma 3.1, there exist a subsequence  $\{v_n\}$  and  $v_0 \in H_0^1(\Omega)$  such that  $v_n \rightharpoonup v_0$  in  $X$  and  $v_n \rightarrow v_0$  in  $L^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ . Then by Proposition 2.3, we have

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f v_n^2 dx = \int_{\Omega} f v_0^2 dx = 1. \quad (3.9)$$

Moreover,

$$\int_{\Omega} |\nabla v_0|^2 dx = \int_{\mathbb{R}^3} |\nabla v_0|^2 + V v_0^2 dx \leq \liminf_{n \rightarrow \infty} \|v_n\|_{\mu_n}^2 \leq \frac{\lambda_1(f_{\bar{\Omega}}) + \lambda_2(f_{\bar{\Omega}})}{2}.$$

According to the fact in Lemma 3.2 that  $e_{1,\mu} \rightarrow e_1$  in  $X$  as  $\mu \rightarrow \infty$ , we deduce that

$$\|e_{1,\mu_n} - e_1\|_{\mu_n} \rightarrow 0 \quad \text{as } n \rightarrow \infty. \quad (3.10)$$

It follows from (3.8), (3.10) and  $v_n \rightharpoonup v_0$  in  $X$  that

$$\lim_{n \rightarrow \infty} \langle v_n, e_{1,\mu_n} \rangle_{\mu_n} = \int_{\Omega} \nabla v_0 \nabla e_1 dx = 0. \quad (3.11)$$

From (3.9), (3.11) and Remark 1.1, we conclude that  $\int_{\Omega} |\nabla v_0|^2 dx \geq \lambda_2(f_{\bar{\Omega}})$  which is a contradiction. This completes the proof.  $\square$

## 4 Palais-Smale sequence

In this section, we study the mountain pass geometry and the compactness condition for the functional  $J_{\mu,\lambda}$ . Let us recall the well known mountain pass theorem.

**Theorem 4.1** ([36], Mountain pass theorem). *Let  $E$  be a Banach space,  $J \in C^1(E, \mathbb{R})$ ,  $v \in E$  and  $\rho > 0$  be such that  $\|v\| > \rho$  and*

$$b := \inf_{\|u\|=\rho} J(u) > J(0) \geq J(v).$$

*If  $J$  satisfies the Palais-Smale condition at level  $\alpha$  with*

$$\alpha := \inf_{\gamma \in \Gamma} \max_{t \in [0,1]} J(\gamma(t)) \quad \text{and} \quad \Gamma := \{\gamma \in C([0,1], E) : \gamma(0) = 0, \gamma(1) = v\},$$

*then  $\alpha$  is a critical value of  $J$  and  $\alpha \geq b$ .*

**Definition 4.2.** (i) A sequence  $\{u_n\} \subset X$  is called a Palais-Smale sequence at level  $\alpha \in \mathbb{R}$  ( $(PS)_\alpha$ -sequence for short) for the functional  $J_{\mu,\lambda}$  if  $J_{\mu,\lambda}(u_n) \rightarrow \alpha$  and  $J'_{\mu,\lambda}(u_n) \rightarrow 0$ .  
(ii) We say that the functional  $J_{\mu,\lambda}$  satisfies the Palais-Smale condition at level  $\alpha$  ( $(PS)_\alpha$ -condition) if every  $(PS)_\alpha$ -sequence has a convergent subsequence.

Next, we show that the functional  $J_{\mu,\lambda}$  has the mountain pass geometry in  $\lambda > \lambda_1(f_\Omega)$  and near  $\lambda_1(f_\Omega)$  for  $\mu > 0$  large enough. To prove the mountain pass geometry of the functional  $J_{\mu,\lambda}$ , we decompose each  $u \in X$  as  $u = te_{1,\mu} + w$ , where  $t \in \mathbb{R}$ ,  $w \in X$  and  $\langle w, e_{1,\mu} \rangle_\mu = 0$ . Then

$$\|u\|_\mu^2 = \|e_{1,\mu}\|_\mu^2 t^2 + \|w\|_\mu^2 = \lambda_{1,\mu}(f)t^2 + \|w\|_\mu^2. \quad (4.1)$$

Moreover, by Lemmas 3.2 and 3.3, we have

$$\lambda_{1,\mu}(f) \int_{\mathbb{R}^3} f e_{1,\mu} w dx = \int_{\mathbb{R}^N} \nabla e_{1,\mu} \nabla w + \mu V e_{1,\mu} w dx = 0 \quad (4.2)$$

and

$$\lambda_{2,\mu}(f) \int_{\mathbb{R}^3} f w^2 dx \leq \|w\|_\mu^2. \quad (4.3)$$

Thus, by (4.1) – (4.3), we deduce that

$$\begin{aligned} & \frac{1}{2} \left( \|u\|_\mu^2 - \lambda \int_{\mathbb{R}^3} f u^2 dx \right) \\ &= \frac{1}{2} \left( \lambda_{1,\mu}(f)t^2 + \|w\|_\mu^2 - \lambda \int_{\mathbb{R}^3} (t^2 f e_{1,\mu}^2 + 2t f e_{1,\mu} w + f w^2) dx \right) \\ &\geq \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{1,\mu}(f)} \right) \lambda_{1,\mu}(f)t^2 + \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{2,\mu}(f)} \right) \|w\|_\mu^2 \\ &= \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{1,\mu}(f)} \right) \|u\|_\mu^2 + \frac{\lambda}{2} \left( \frac{1}{\lambda_{1,\mu}(f)} - \frac{1}{\lambda_{2,\mu}(f)} \right) \|w\|_\mu^2. \end{aligned} \quad (4.4)$$

**Lemma 4.3.** Suppose that  $4 \leq p < 6$  and conditions (V1)-(V3), (F), (G) and (K1) hold. In addition, for  $p = 4$ , assume (K3) and (D1). Then there exists  $\delta_0 > 0$  such that for every  $0 < \lambda < \lambda_1(f_\Omega) + \delta_0$ , there exist  $\rho_\lambda, \eta_\lambda > 0$  and  $\varphi_0 \in H_0^1(\Omega)$  such that  $\|\varphi_0\|_\mu > \rho_\lambda$  and

$$\inf_{\|u\|_\mu = \rho_\lambda} J_{\mu,\lambda}(u) \geq \eta_\lambda > 0 > J_{\mu,\lambda}(\varphi_0)$$

for  $\mu > 0$  large enough.

*Proof.* We first show that there exists  $\delta_0 > 0$  such that for every  $0 < \lambda < \lambda_1(f_\Omega) + \delta_0$ , there exist  $\rho_\lambda, \eta_\lambda > 0$  such that

$$\inf_{\|u\|_\mu = \rho_\lambda} J_{\mu,\lambda}(u) \geq \eta_\lambda > 0$$

as  $\mu > 0$  large enough. We separate this part into two cases.

Case (1):  $0 < \lambda < \lambda_1(f_\Omega)$ . By Lemma 3.2(ii), we have

$$\lambda_{1,\mu}(f) \geq \frac{\lambda + \lambda_1(f_\Omega)}{2} \quad \text{for } \mu > 0 \text{ large enough.} \quad (4.5)$$

By (2.3), (4.4), (4.5) and condition (G), we deduce that for  $\mu > 0$  large enough,

$$\begin{aligned} J_{\mu,\lambda}(u) &\geq \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{1,\mu}(f)} \right) \|u\|_\mu^2 + \frac{1}{4} N(u) - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{pS^p} \|u\|_\mu^p \\ &\geq \|u\|_\mu^2 \left( \frac{1}{2} \left( \frac{\lambda_1(f_\Omega) - \lambda}{\lambda_1(f_\Omega) + \lambda} \right) - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{pS^p} \|u\|_\mu^{p-2} \right), \end{aligned}$$

where  $N(u)$  is as in (2.4). Let

$$\rho_\lambda = \left( \frac{1}{4} \left( \frac{\lambda_1(f_\Omega) - \lambda}{\lambda_1(f_\Omega) + \lambda} \right) \frac{pS^p}{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}} \right)^{\frac{1}{p-2}}.$$

Then for  $\mu > 0$  large enough and  $\|u\|_\mu = \rho_\lambda$ , we have

$$J_{\mu,\lambda}(u) \geq \frac{1}{4} \left( \frac{\lambda_1(f_\Omega) - \lambda}{\lambda_1(f_\Omega) + \lambda} \right) \rho_\lambda^2 =: \eta_\lambda > 0.$$

Case (2):  $\lambda \geq \lambda_1(f_\Omega)$ . Using (2.3) and (4.4), for  $\mu > 0$  large enough, we have

$$J_{\mu,\lambda}(u) \geq \Lambda_{1,\mu} \|u\|_\mu^2 + \Lambda_{2,\mu} \|w\|_\mu^2 + \frac{N(te_{1,\mu})}{4} + \frac{N(u) - N(te_{1,\mu})}{4} - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{pS^p} \|u\|_\mu^p, \quad (4.6)$$

where  $u = te_{1,\mu} + w$  with  $t \in \mathbb{R}$  and  $w \in \{v \in X : \langle v, e_{1,\mu} \rangle_\mu = 0\}$ ,

$$\Lambda_{1,\mu} = \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{1,\mu}(f)} \right) \quad \text{and} \quad \Lambda_{2,\mu} = \frac{\lambda}{2} \left( \frac{1}{\lambda_{1,\mu}(f)} - \frac{1}{\lambda_{2,\mu}(f)} \right).$$

By Lemma 3.3, we may deduce that for  $\mu > 0$  large enough and  $\lambda \geq \lambda_1(f_\Omega)$ ,

$$\Lambda_{2,\mu} \geq \frac{1}{2} \left( 1 - \frac{\lambda_1(f_\Omega)}{\lambda_{2,\mu}(f)} \right) \geq \frac{\lambda_2(f_\Omega) - \lambda_1(f_\Omega)}{2(\lambda_2(f_\Omega) + \lambda_1(f_\Omega))} =: \Lambda_0. \quad (4.7)$$

By Proposition 2.1(i) and Lemma 3.2(ii), we deduce that  $N(e_{1,\mu}) \rightarrow N(e_1)$  as  $\mu \rightarrow \infty$ , which implies that

$$N(e_{1,\mu}) \geq \frac{1}{2} N(e_1) \quad \text{for } \mu > 0 \text{ large enough.} \quad (4.8)$$

Moreover, by condition (K1), we have  $N(e_1) > 0$ . By the mean value theorem, there exists  $\theta$  with  $0 < \theta < 1$  such that

$$\frac{1}{4} |N(te_{1,\mu} + w) - N(te_{1,\mu})| = \left| \int_{\mathbb{R}^3} K \phi_{te_{1,\mu} + \theta w}(te_{1,\mu} + \theta w) w dx \right|. \quad (4.9)$$

When  $K \in L^2(\mathbb{R}^3)$ , by (2.5), Hölder, Sobolev and Young's inequalities, we deduce that

$$\begin{aligned} &\left| \int_{\mathbb{R}^3} K \phi_{te_{1,\mu} + \theta w}(te_{1,\mu} + \theta w) w dx \right| \\ &\leq S^{-1} \|K\|_{L^2} \|\phi_{te_{1,\mu} + \theta w}\|_{D^{1,2}} \|te_{1,\mu} + \theta w\|_{L^6} \|w\|_{L^6} \\ &\leq S^{-6} \|K\|_{L^2}^2 \|te_{1,\mu} + \theta w\|_\mu^3 \|w\|_\mu \\ &= S^{-6} \|K\|_{L^2}^2 (t^2 \|e_{1,\mu}\|_\mu^2 + \theta^2 \|w\|_\mu^2)^{\frac{3}{2}} \|w\|_\mu \\ &\leq \sqrt{2} S^{-6} \|K\|_{L^2}^2 |t|^3 \|e_{1,\mu}\|_\mu^3 \|w\|_\mu + \sqrt{2} S^{-6} \|K\|_{L^2}^2 \|w\|_\mu^4 \\ &< \frac{N(e_1)}{16} t^4 + \frac{12^3 \|K\|_{L^2}^8 \lambda_1(f_\Omega)^6}{S^{24} N(e_1)^3} \|w\|_\mu^4 + \sqrt{2} S^{-6} \|K\|_{L^2}^2 \|w\|_\mu^4. \end{aligned} \quad (4.10)$$

Analogously, for  $K \in L^\infty(\mathbb{R}^3)$ , by (2.3), (2.5) and Young's inequality, we have

$$\begin{aligned}
& \left| \int_{\mathbb{R}^3} K \phi_{te_{1,\mu} + \theta w} (te_{1,\mu} + \theta w) w dx \right| \\
& \leq \|K\|_{L^\infty} \|\phi_{te_{1,\mu} + \theta w}\|_{L^6} \|te_{1,\mu} + \theta w\|_{L^{\frac{12}{5}}} \|w\|_{L^{\frac{12}{5}}} \\
& \leq \sqrt{2} S^{-6} |\{V < c_0\}| \|K\|_{L^\infty}^2 |t|^3 \|e_{1,\mu}\|_\mu^3 \|w\|_\mu + \sqrt{2} S^{-6} |\{V < c_0\}| \|K\|_{L^\infty}^2 \|w\|_\mu^4 \\
& < \frac{N(e_1)}{16} t^4 + \frac{12^3 |\{V < c_0\}|^4 \|K\|_{L^\infty}^8 \lambda_1(f_\Omega)^6}{S^{24} N(e_1)^3} \|w\|_\mu^4 + \sqrt{2} S^{-6} |\{V < c_0\}| \|K\|_{L^\infty}^2 \|w\|_\mu^4. \quad (4.11)
\end{aligned}$$

Subsequently, combining (4.6) – (4.11), for  $\mu > 0$  large enough, we have

$$\begin{aligned}
& J_{\mu,\lambda}(u) \\
& \geq -|\Lambda_{1,\mu}| \|u\|_\mu^2 + \Lambda_0 \|w\|_\mu^2 + \frac{N(e_1)}{8} t^4 + \frac{N(te_{1,\mu} + w) - N(te_{1,\mu})}{4} - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^p \\
& \geq -|\Lambda_{1,\mu}| \|u\|_\mu^2 + \Lambda_0 \|w\|_\mu^2 + \frac{N(e_1)}{16} t^4 - C_K \|w\|_\mu^4 - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^p \\
& = -|\Lambda_{1,\mu}| \|u\|_\mu^2 + \Lambda_0 \|w\|_\mu^2 + \frac{N(e_1)}{16 \lambda_{1,\mu}(f)^2} (\|u\|_\mu^2 - \|w\|_\mu^2)^2 - C_K \|w\|_\mu^4 - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^p \\
& \geq -|\Lambda_{1,\mu}| \|u\|_\mu^2 + \Lambda_0 \|w\|_\mu^2 + \frac{N(e_1)}{16 \lambda_1(f_\Omega)^2} \left( \frac{\|u\|_\mu^4}{2} - \|w\|_\mu^4 \right) - C_K \|w\|_\mu^4 - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^p \\
& = -|\Lambda_{1,\mu}| \|u\|_\mu^2 + \|u\|_\mu^4 \left( \frac{N(e_1)}{32 \lambda_1(f_\Omega)^2} - \frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^{p-4} \right) \\
& \quad + \|w\|_\mu^2 \left( \Lambda_0 - \left( \frac{N(e_1)}{16 \lambda_1(f_\Omega)^2} + C_K \right) \|w\|_\mu^2 \right), \quad (4.12)
\end{aligned}$$

where

$$C_K = \begin{cases} \frac{12^3 \|K\|_{L^2}^8 \lambda_1(f_\Omega)^6}{S^{24} N(e_1)^3} + \frac{\sqrt{2} \|K\|_{L^2}^2}{S^6}, & \text{for } K \in L^2(\mathbb{R}^3), \\ \frac{12^3 |\{V < c_0\}|^4 \|K\|_{L^\infty}^8 \lambda_1(f_\Omega)^6}{S^{24} N(e_1)^3} + \frac{\sqrt{2} |\{V < c_0\}| \|K\|_{L^\infty}^2}{S^6}, & \text{for } K \in L^\infty(\mathbb{R}^3). \end{cases}$$

Let

$$\rho_\lambda = \begin{cases} \min \left\{ \left( \frac{N(e_1) p S^p}{64 \lambda_1(f_\Omega)^2 \|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}} \right)^{\frac{1}{p-4}}, \left( \Lambda_0 \left( \frac{N(e_1)}{16 \lambda_1(f_\Omega)^2} + C_K \right)^{-1} \right)^{\frac{1}{2}} \right\}, & \text{for } 4 < p < 6, \\ \left( \Lambda_0 \left( \frac{N(e_1)}{16 \lambda_1(f_\Omega)^2} + C_K \right)^{-1} \right)^{\frac{1}{2}}, & \text{for } p = 4. \end{cases} \quad (4.13)$$

Then for  $\|u\|_\mu = \rho_\lambda$ , we have

$$\frac{\|g\|_{L^\infty} |\{V < c_0\}|^{\frac{6-p}{6}}}{p S^p} \|u\|_\mu^{p-4} \leq \frac{N(e_1)}{64 \lambda_1(f_\Omega)^2} \quad \text{and} \quad \Lambda_0 - \left( \frac{N(e_1)}{16 \lambda_1(f_\Omega)^2} + C_K \right) \|w\|_\mu^2 \geq 0. \quad (4.14)$$

Note that under condition (D1), the first inequality in (4.14) holds for  $p = 4$ . Let

$$\delta_0 = \frac{N(e_1)}{256 \lambda_1(f_\Omega)} \rho_\lambda^2 \quad (4.15)$$

and

$$\delta_\mu = \frac{N(e_1)\lambda_{1,\mu}(f)}{64\lambda_1(f_\Omega)^2}\rho_\lambda^2.$$

By Lemma 3.2(ii), we deduce that for  $\mu > 0$  large enough,

$$\lambda_1(f_\Omega) + \delta_0 \leq \lambda_{1,\mu}(f) + 2\delta_0 \leq \lambda_{1,\mu}(f) + \delta_\mu.$$

This implies that for every  $\lambda_1(f_\Omega) \leq \lambda < \lambda_1(f_\Omega) + \delta_0$  and  $\mu > 0$  large enough,

$$|\Lambda_{1,\mu}| = \frac{1}{2} \left( \frac{\lambda}{\lambda_{1,\mu}(f)} - 1 \right) \leq \frac{1}{2} \left( \frac{\lambda_{1,\mu}(f) + \delta_\mu}{\lambda_{1,\mu}(f)} - 1 \right) = \frac{N(e_1)}{128\lambda_1(f_\Omega)^2}\rho_\lambda^2. \quad (4.16)$$

It follows from (4.12), (4.14) and (4.16) that for every  $\lambda_1(f_\Omega) \leq \lambda < \lambda_1(f_\Omega) + \delta_0$ ,  $\|u\|_\mu = \rho_\lambda$  and  $\mu > 0$  large enough,

$$J_{\mu,\lambda}(u) \geq \frac{N(e_1)}{128\lambda_1(f_\Omega)^2}\rho_\lambda^4 =: \eta_\lambda > 0.$$

Consequently, for every  $0 < \lambda < \lambda_1(f_\Omega) + \delta_0$ , we have  $\inf_{\|u\|_\mu=\rho_\lambda} J_{\mu,\lambda}(u) \geq \eta_\lambda > 0$  as  $\mu > 0$  large enough.

Next, we show that there exists  $\varphi_0 \in H_0^1(\Omega)$  such that  $\|\varphi_0\|_\mu > \rho_\lambda$  and  $J_{\mu,\lambda}(\varphi_0) < 0$ . We divide the part into two cases:

Case (I):  $4 < p < 6$ . By condition (G), we may choose a function  $\varphi \in H_0^1(\Omega)$  such that  $\int_{\mathbb{R}^3} g|\varphi|^p dx > 0$ . Let  $s > 0$ , then

$$J_{\mu,\lambda}(s\varphi) = \frac{\|\varphi\|_\mu^2 - \lambda \int_{\mathbb{R}^3} f\varphi^2 dx}{2} s^2 + \frac{N(\varphi)}{4} s^4 - \frac{\int_{\mathbb{R}^3} g|\varphi|^p dx}{p} s^p.$$

This implies that there exists  $s_0 > 0$  such that  $\|s_0\varphi\|_\mu > \rho_\lambda$  and  $J_{\mu,\lambda}(s_0\varphi) < 0$ .

Case (II):  $p = 4$ . By conditions (G) and (K3), we can choose a function  $\varphi \in H_0^1(\Omega)$  such that  $\int_{\mathbb{R}^3} g\varphi^4 dx > 0$  and  $N(\varphi) = 0$ . Let  $s > 0$ , then

$$J_{\mu,\lambda}(s\varphi) = \frac{\|\varphi\|_\mu^2 - \lambda \int_{\mathbb{R}^3} f\varphi^2 dx}{2} s^2 - \frac{\int_{\mathbb{R}^3} g\varphi^4 dx}{4} s^4,$$

Thus, there exists  $s_0 > 0$  such that  $\|s_0\varphi\|_\mu > \rho_\lambda$  and  $J_{\mu,\lambda}(s_0\varphi) < 0$ . We complete the proof.  $\square$

Next, we study the  $(PS)_{\beta_\lambda(\mu)}$ -condition for the functional  $J_{\mu,\lambda}$ , where  $\beta_\lambda(\mu)$  is a real value function defined in  $\mu > 0$  and  $\beta_\lambda(\mu) > 0$  for all  $\lambda, \mu > 0$ . Then we have the following results.

**Lemma 4.4.** *Suppose that  $4 \leq p < 6$  and conditions (V1)-(V3), (F), (G), (K1) and (K2) hold. In addition, for  $p = 4$ , assume (K3). If there exists  $d_\lambda > 0$  such that*

$$0 < \beta_\lambda(\mu) < d_\lambda$$

*for all  $\mu > 0$  large enough, then there exists  $D_0 > 0$  such that the  $(PS)_{\beta_\lambda(\mu)}$ -sequence  $\{u_n\}$  for  $J_{\mu,\lambda}$  satisfies  $\|u_n\|_\mu < D_0$  for all  $\mu > 0$  large enough.*

*Proof.* Suppose on the contrary. Then there exist two sequences  $\{\mu_n\}, \{D_n\} \subset \mathbb{R}^+$  with  $\mu_n \rightarrow \infty, D_n \rightarrow \infty$  as  $n \rightarrow \infty$  such that for every  $n \in \mathbb{N}$ , there exists a  $(PS)_{\beta_\lambda(\mu_n)}$ -sequence  $\{u_{n,m}\}_{m \in \mathbb{N}}$  with  $\|u_{n,m}\|_{\mu_n} > D_n$  for all  $m \in \mathbb{N}$ ,

$$\frac{1}{2} \|u_{n,m}\|_{\mu_n}^2 - \frac{\lambda}{2} \int_{\mathbb{R}^3} f u_{n,m}^2 dx - \frac{1}{p} \int_{\mathbb{R}^3} g |u_{n,m}|^p dx + \frac{1}{4} N(u_{n,m}) \rightarrow \beta_\lambda(\mu_n) \quad (4.17)$$

and

$$\int_{\mathbb{R}^3} (\nabla u_{n,m} \nabla \varphi + \mu_n V u_{n,m} \varphi - \lambda f u_{n,m} \varphi) dx - \int_{\mathbb{R}^3} g |u_{n,m}|^{p-2} u_{n,m} \varphi dx + \int_{\mathbb{R}^3} K \phi_{u_{n,m}} u_{n,m} \varphi dx \rightarrow 0 \quad (4.18)$$

as  $m \rightarrow \infty$ .

Let  $w_n = u_{n,n}$ , then  $\|w_n\|_{\mu_n} \rightarrow \infty$  as  $n \rightarrow \infty$ . Let  $v_n = \frac{w_n}{\|w_n\|_{\mu_n}}$ , then  $\|v_n\|_{\mu_n} = 1$ . By Lemma 3.1, there exist a subsequence  $\{v_n\}$  and  $v_0 \in H_0^1(\Omega)$  such that  $v_n \rightharpoonup v_0$  in  $X$ ,  $v_n \rightarrow v_0$  in  $L^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ , and  $\lim_{n \rightarrow \infty} N(v_n) = N(v_0)$ . Then by conditions (F) and (G), we have

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f v_n^2 dx = \int_{\Omega} f v_0^2 dx \quad \text{and} \quad \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} g |v_n|^p dx = \int_{\Omega} g |v_0|^p dx.$$

Dividing (4.17) by  $\|w_n\|_{\mu_n}^2$  and the boundedness of  $\beta_\lambda(\mu_n)$ , we obtain

$$\frac{1}{2} - \frac{\lambda}{2} \int_{\mathbb{R}^3} f v_n^2 dx - \frac{1}{p} \|w_n\|_{\mu_n}^{p-2} \int_{\mathbb{R}^3} g |v_n|^p dx + \frac{1}{4} \|w_n\|_{\mu_n}^2 N(v_n) \rightarrow 0. \quad (4.19)$$

Dividing (4.18) by  $\|w_n\|_{\mu_n}$ , we have

$$\int_{\mathbb{R}^3} (\nabla v_n \nabla \varphi + \mu_n V v_n \varphi - \lambda f v_n \varphi) dx - \|w_n\|_{\mu_n}^{p-2} \int_{\mathbb{R}^3} g |v_n|^{p-2} v_n \varphi dx + \|w_n\|_{\mu_n}^2 \int_{\mathbb{R}^3} K \phi_{v_n} v_n \varphi dx \rightarrow 0. \quad (4.20)$$

If the assumption that  $v_0 = 0$  a.e. in  $\Omega$  holds, then we have

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f v_n^2 dx = \int_{\Omega} f v_0^2 dx = 0. \quad (4.21)$$

Choosing  $\varphi = v_n$  in (4.20), we obtain

$$1 - \lambda \int_{\mathbb{R}^3} f v_n^2 dx - \|w_n\|_{\mu_n}^{p-2} \int_{\mathbb{R}^3} g |v_n|^p dx + \|w_n\|_{\mu_n}^2 N(v_n) \rightarrow 0. \quad (4.22)$$

For  $4 < p < 6$ , combining (4.19), (4.21) and (4.22), we deduce that

$$\lim_{n \rightarrow \infty} \|w_n\|_{\mu_n}^{p-2} \int_{\mathbb{R}^3} g |v_n|^p dx = \frac{p}{4-p} < 0$$

However, by (4.21), (4.22) implies that  $\lim_{n \rightarrow \infty} \|w_n\|_{\mu_n}^{p-2} \int_{\mathbb{R}^3} g |v_n|^p dx > 0$ , which is a contradiction. Similarly, we also have a contradiction for  $p = 4$  by (4.19), (4.21) and (4.22).

Now, we prove the assumption that  $v_0 = 0$  a.e. in  $\Omega$ . For  $4 < p < 6$ , dividing (4.20) by  $\|w_n\|_{\mu_n}^{p-2}$ , we obtain

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} g |v_n|^{p-2} v_n \varphi dx = \int_{\Omega} g |v_0|^{p-2} v_0 \varphi dx = 0. \quad (4.23)$$

Clearly,  $v_0 \chi_{\{x \in \Omega : g(x) > 0\}} \in L^r(\Omega)$  for  $2 \leq r \leq 6$ . Then there exists a sequence  $\{\varphi_n\} \subset C_0^\infty(\Omega)$  such that  $\varphi_n \rightarrow v_0 \chi_{\{x \in \Omega : g(x) > 0\}}$  in  $L^r(\Omega)$  for  $2 \leq r \leq 6$ . Therefore, choosing  $\varphi = \varphi_n$  in (4.23) and taking  $n \rightarrow \infty$ , we have

$$0 = \lim_{n \rightarrow \infty} \int_{\Omega} g |v_0|^{p-2} v_0 \varphi_n dx = \int_{\{x \in \Omega : g(x) > 0\}} g |v_0|^p dx,$$

which implies that  $v_0 = 0$  a.e. in  $\{x \in \Omega : g(x) > 0\}$ . In a same way, we obtain that  $v_0 = 0$  a.e. in  $\{x \in \Omega : g(x) < 0\}$ . For the remaining part, by combining (4.19) and (4.22), we have

$$0 = \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} K \phi_{v_n} v_n^2 dx = \int_{\{x \in \Omega : g(x) = 0\}} K \phi_{v_0} v_0^2 dx.$$

By condition (K2), we obtain that  $v_0 = 0$  a.e. in  $\{x \in \Omega : g(x) = 0\}$ .

For  $p = 4$ , dividing (4.20) by  $\|w_n\|_{\mu_n}^2$ , we obtain

$$\lim_{n \rightarrow \infty} \left( \int_{\mathbb{R}^3} g v_n^3 \varphi dx - \int_{\mathbb{R}^3} K \phi_{v_n} v_n \varphi dx \right) = \int_{\Omega} g v_0^3 \varphi dx - \int_{\Omega} K \phi_{v_0} v_0 \varphi dx = 0. \quad (4.24)$$

In a same way as the argument of the case  $4 < p < 6$  and by condition (K3), we have

$$0 = \int_{\{x \in \Omega : g(x) > 0\}} g v_0^4 dx - \int_{\{x \in \Omega : g(x) > 0\}} K \phi_{v_0} v_0^2 dx = \int_{\{x \in \Omega : g(x) > 0\}} g v_0^4 dx,$$

which implies that  $v_0 = 0$  a.e. in  $\{x \in \Omega : g(x) > 0\}$ . Similarly, we have

$$\int_{\{x \in \Omega : g(x) < 0\}} g v_0^4 dx = \int_{\{x \in \Omega : g(x) < 0\}} K \phi_{v_0} v_0^2 dx$$

and

$$\int_{\{x \in \Omega : g(x) = 0\}} K \phi_{v_0} v_0^2 dx = 0.$$

By condition (K2), we obtain that  $v_0 = 0$  a.e. in  $\{x \in \Omega : g(x) \leq 0\}$ . This completes the proof.  $\square$

**Lemma 4.5.** *Suppose that  $2 < p < 6$  and conditions (V1)-(V2), (F), (G) and (K1) hold. Let  $\{u_n\}$  be a  $(PS)_\beta$ -sequence for  $J_{\mu,\lambda}$ . If there exists  $D_0 > 0$  such that*

$$\|u_n\|_\mu < D_0 \quad (4.25)$$

*for all  $\mu > 0$  large enough, then the sequence  $\{u_n\}$  has a convergent subsequence.*

*Proof.* By (4.25), there exist a subsequence  $\{u_n\}$  and  $u_0 \in X$  such that  $u_n \rightharpoonup u_0$  in  $X$  and  $u_n \rightarrow u_0$  in  $L_{loc}^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ . It follows from Proposition 2.3 and  $X \hookrightarrow H^1(\mathbb{R}^3)$  that

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f u_n^2 dx = \int_{\mathbb{R}^3} f u_0^2 dx. \quad (4.26)$$

Next, we show that  $u_n \rightarrow u_0$  in  $X$ . Let  $v_n = u_n - u_0$ . By (4.25), we deduce that for  $\mu > 0$  large enough,

$$\|v_n\|_\mu^2 \leq \|u_n\|_\mu + \|u_0\|_\mu < D_1. \quad (4.27)$$

From Brézis-Lieb lemma [37] and Proposition 2.2, we have

$$\|v_n\|_\mu^2 = \|u_n\|_\mu^2 - \|u_0\|_\mu^2 + o(1),$$

$$N(v_n) = N(u_n) - N(u_0) + o(1)$$

and

$$\int_{\mathbb{R}^3} g |v_n|^p dx = \int_{\mathbb{R}^3} g |u_n|^p dx - \int_{\mathbb{R}^3} g |u_0|^p dx + o(1).$$

Moreover, we have

$$\int_{\mathbb{R}^3} v_n^2 dx \leq \frac{1}{\mu c_0} \int_{\{V \geq c_0\}} \mu V v_n^2 dx + \int_{\{V < c_0\}} v_n^2 dx \leq \frac{1}{\mu c_0} \int_{\mathbb{R}^3} \mu V v_n^2 dx + o(1). \quad (4.28)$$

By (2.3), (4.27) and (4.28), we deduce

$$\begin{aligned}
\int_{\mathbb{R}^3} g|v_n|^p dx &\leq \|g\|_{L^\infty} \left( \int_{\mathbb{R}^3} |v_n|^p dx \right)^{\frac{p-2}{p}} \left( \int_{\mathbb{R}^3} |v_n|^p dx \right)^{\frac{2}{p}} \\
&\leq \|g\|_{L^\infty} \left( |\{V < c_0\}|^{\frac{6-p}{6}} S^{-p} \|v_n\|_\mu^p \right)^{\frac{p-2}{p}} \left( \|v_n\|_{L^2}^{\frac{6-p}{2}} \|v_n\|_{L^6}^{\frac{3p-6}{2}} \right)^{\frac{2}{p}} \\
&\leq \|g\|_{L^\infty} \left( |\{V < c_0\}|^{\frac{6-p}{6}} S^{-p} D_1^{\frac{p}{2}} \right)^{\frac{p-2}{p}} \left( \left( \frac{1}{\mu c_0} \int_{\mathbb{R}^3} \mu V v_n^2 dx + o(1) \right)^{\frac{6-p}{4}} (S^{-1} \|v_n\|_{D^{1,2}})^{\frac{3p-6}{2}} \right)^{\frac{2}{p}} \\
&\leq \|g\|_{L^\infty} \left( |\{V < c_0\}|^{\frac{6-p}{6}} S^{-p} D_1^{\frac{p}{2}} \right)^{\frac{p-2}{p}} \left( (\mu c_0)^{\frac{p-6}{4}} \|v_n\|_\mu^{\frac{6-p}{2}} S^{\frac{6-3p}{2}} \|v_n\|_\mu^{\frac{3p-6}{2}} \right)^{\frac{2}{p}} + o(1) \\
&\leq D_2 \mu^{\frac{p-6}{2p}} \|v_n\|_\mu^2 + o(1),
\end{aligned} \tag{4.29}$$

where  $D_2 = \|g\|_{L^\infty} |\{V < c_0\}|^{\frac{(6-p)(p-2)}{6p}} S^{-1-p+\frac{6}{p}} D_1^{\frac{p-2}{2}} c_0^{\frac{p-6}{2p}}$ .

Since  $J'_{\mu,\lambda}(u_n) \rightarrow 0$ , we have

$$o(1) = \int_{\mathbb{R}^3} \nabla u_n \nabla \varphi + \mu V u_n \varphi dx - \lambda \int_{\mathbb{R}^3} f u_n \varphi dx - \int_{\mathbb{R}^3} g |u_n|^{p-2} u_n \varphi dx + \int_{\mathbb{R}^3} K \phi_{u_n} u_n \varphi dx, \tag{4.30}$$

which implies that

$$0 = \int_{\mathbb{R}^3} \nabla u_0 \nabla \varphi + \mu V u_0 \varphi dx - \lambda \int_{\mathbb{R}^3} f u_0 \varphi dx - \int_{\mathbb{R}^3} g |u_0|^{p-2} u_0 \varphi dx + \int_{\mathbb{R}^3} K \phi_{u_0} u_0 \varphi dx. \tag{4.31}$$

Choosing  $\varphi = u_n$  in (4.30) and  $\varphi = u_0$  in (4.31), we obtain

$$o(1) = \|u_n\|_\mu^2 - \lambda \int_{\mathbb{R}^3} f u_n^2 dx - \int_{\mathbb{R}^3} g |u_n|^p dx + N(u_n) \tag{4.32}$$

and

$$0 = \|u_0\|_\mu^2 - \lambda \int_{\mathbb{R}^3} f u_0^2 dx - \int_{\mathbb{R}^3} g |u_0|^p dx + N(u_0). \tag{4.33}$$

By (4.26), (4.32) and (4.33), we have

$$o(1) = \|v_n\|_\mu^2 - \int_{\mathbb{R}^3} g |v_n|^p dx + N(v_n). \tag{4.34}$$

It follows from (4.29), (4.34) and  $N(v_n) \geq 0$  that

$$o(1) \geq \|v_n\|_\mu^2 - D_2 \mu^{\frac{p-6}{2p}} \|v_n\|_\mu^2.$$

Therefore, we have  $o(1) \geq \frac{1}{2} \|v_n\|_\mu^2$  for  $\mu$  large enough, which implies that  $v_n \rightarrow 0$  in  $X$ . Hence we have  $u_n \rightarrow u_0$  in  $X$ ,

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} g |u_n|^p dx = \int_{\mathbb{R}^3} g |u_0|^p dx \quad \text{and} \quad \lim_{n \rightarrow \infty} N(u_n) = N(u_0),$$

which implies that  $J_{\mu,\lambda}(u_0) = \beta_\lambda(\mu)$ . This completes the proof.  $\square$



## 5 The proof of Theorems 1.2

In this section, we prove Theorem 1.2. The results will hold by the following two theorems.

**Theorem 5.1.** *Under the assumptions of Theorems 1.2. There exists  $\delta_0 > 0$  such that for every  $0 < \lambda < \lambda_1(f_{\bar{\Omega}}) + \delta_0$ , Eq.  $(S_{\mu,\lambda})$  has a positive solution  $u^+$  with  $J_{\mu,\lambda}(u^+) > 0$  for  $\mu > 0$  large enough.*

*Proof.* By Lemma 4.3, there exists  $\delta_0 > 0$  such that for every  $0 < \lambda < \lambda_1(f_{\bar{\Omega}}) + \delta_0$ , there exist  $\rho_\lambda, \eta_\lambda > 0$  and  $\varphi_0 \in H_0^1(\Omega)$  such that  $\|\varphi_0\|_\mu > \rho_\lambda$  and

$$\inf_{\|u\|_\mu = \rho_\lambda} J_{\mu,\lambda}(u) \geq \eta_\lambda > 0 > J_{\mu,\lambda}(\varphi_0)$$

for all  $\mu > 0$  large enough. Let

$$\alpha_\lambda(\mu) := \inf_{\gamma \in \Gamma} \max_{t \in [0,1]} J_{\mu,\lambda}(\gamma(t)) \quad \text{with} \quad \Gamma := \{\gamma \in C([0,1], X) : \gamma(0) = 0, \gamma(1) = \varphi_0\}.$$

Since  $\varphi_0 \in H_0^1(\Omega)$ , there exists a constant  $d_\lambda > 0$ , independent of  $\mu$ , such that

$$0 < \eta_\lambda \leq \alpha_\lambda(\mu) \leq \max_{0 \leq t \leq 1} J_{\mu,\lambda}(t\varphi_0) \leq d_\lambda \quad \text{for all } \mu > 0 \text{ large enough.} \quad (5.1)$$

In order to prove the positivity of solutions, we follow from the argument in [21]. Since  $J_{\mu,\lambda}(u) = J_{\mu,\lambda}(|u|)$ , we may assume that for every  $n \in \mathbb{N}$ , there exists  $\gamma_n \in \Gamma$  with  $\gamma_n(t) \geq 0$  for all  $t \in [0, 1]$  such that

$$\alpha_\lambda(\mu) \leq \max_{t \in [0,1]} J_{\mu,\lambda}(\gamma_n(t)) < \alpha_\lambda(\mu) + \frac{1}{n}.$$

Denote  $J_{\mu,\lambda}(v_n) = \max_{t \in [0,1]} J_{\mu,\lambda}(\gamma_n(t))$ . By Ekeland's variational principle [38], we have a  $(PS)_{\alpha_\lambda(\mu)}$ -sequence  $\{u_n\}$  for  $J_{\mu,\lambda}$  satisfying

$$\alpha_\lambda(\mu) \leq J_{\mu,\lambda}(u_n) \leq J_{\mu,\lambda}(v_n) < \alpha_\lambda(\mu) + \frac{1}{n}, \quad \|J'_{\mu,\lambda}(u_n)\|_\mu < \frac{1}{\sqrt{n}}$$

and

$$\|v_n - u_n\|_\mu < \frac{1}{\sqrt{n}}. \quad (5.2)$$

By (5.1) and Lemmas 4.4 and 4.5, there exist a subsequence  $\{u_n\}$  and  $u^+ \in X$  such that  $u_n \rightarrow u^+$  in  $X$  as  $n \rightarrow \infty$ ,

$$J_{\mu,\lambda}(u^+) = \alpha_\lambda(\mu) \quad \text{and} \quad J'_{\mu,\lambda}(u^+) = 0.$$

By (5.2) and the fact that  $\gamma_n(t) \geq 0$  for all  $t \in [0, 1]$ , we conclude that  $u^+ \geq 0$  a.e. in  $\mathbb{R}^3$ . It follows from  $J_{\mu,\lambda}(u^+) > 0$  and the strong maximum principle that  $u^+ > 0$  in  $\mathbb{R}^3$ . This proof is complete.  $\square$

**Theorem 5.2.** *Under the assumptions of Theorems 1.2. There exists  $\delta_0 > 0$  such that for every  $\lambda_1(f_{\bar{\Omega}}) < \lambda < \lambda_1(f_{\bar{\Omega}}) + \delta_0$ , Eq.  $(S_{\mu,\lambda})$  has a positive solution  $u^-$  with  $J_{\mu,\lambda}(u^-) < 0$  for  $\mu$  large enough.*

*Proof.* Let  $B_{\rho_\lambda} := \{u \in X : \|u\|_\mu \leq \rho_\lambda\}$  with  $\rho_\lambda$  as in Lemma 4.3. We consider the infimum of  $J_{\mu,\lambda}$  on  $B_{\rho_\lambda}$  and set

$$\bar{\alpha}_\lambda(\mu) := \inf_{\|u\|_\mu \leq \rho_\lambda} J_{\mu,\lambda}(u).$$

By conditions (F) and (G), we deduce

$$J_{\mu,\lambda}(u) \geq -\frac{\|f\|_{L^{\frac{3}{2}}}}{2S^2}\|u\|_{\mu}^2 - \frac{\|g\|_{L^{\infty}}|\{V < c_0\}|^{\frac{6-p}{6}}}{pS^p}\|u\|_{\mu}^p \geq -\frac{\|f\|_{L^{\frac{3}{2}}}}{2S^2}\rho_{\lambda}^2 - \frac{\|g\|_{L^{\infty}}|\{V < c_0\}|^{\frac{6-p}{6}}}{pS^p}\rho_{\lambda}^p.$$

Moreover, for  $t > 0$ , we have

$$J_{\mu,\lambda}(te_1) = -\frac{\lambda - \lambda_1(f_{\overline{\Omega}})}{2}t^2 + \frac{N(e_1)}{4}t^4 + \frac{\int_{\mathbb{R}^3} g e_1^p dx}{p}t^p.$$

This implies there exists  $t_0 > 0$  such that  $\|t_0 e_1\|_{\mu} \leq \rho_{\lambda}$  and  $J_{\mu,\lambda}(t_0 e_1) < 0$ . Hence we conclude that there exist two positive constants  $\overline{\eta}_{\lambda}, \overline{d}_{\lambda} > 0$  such that

$$-\overline{d}_{\lambda} \leq \overline{\alpha}_{\lambda}(\mu) \leq -\overline{\eta}_{\lambda} \quad \text{for all } \mu > 0 \text{ large enough.} \quad (5.3)$$

Since  $J_{\mu,\lambda}(u) = J_{\mu,\lambda}(|u|)$ , we may assume that there exists  $v_n \geq 0$  with  $\|v_n\|_{\mu} \leq \rho_{\lambda}$  such that

$$\overline{\alpha}_{\lambda}(\mu) \leq J_{\mu,\lambda}(v_n) < \overline{\alpha}_{\lambda}(\mu) + \frac{1}{n}.$$

Then by Ekeland's variational principle [38], we have a  $(PS)_{\overline{\alpha}_{\lambda}(\mu)}$ -sequence  $\{u_n\} \subset B_{\rho_{\lambda}}$  satisfying

$$\overline{\alpha}_{\lambda}(\mu) \leq J_{\mu,\lambda}(u_n) \leq J_{\mu,\lambda}(v_n) < \overline{\alpha}_{\lambda}(\mu) + \frac{1}{n}, \quad \|J'_{\mu,\lambda}(u_n)\|_{\mu} < \frac{1}{\sqrt{n}}$$

and

$$\|u_n - v_n\|_{\mu} < \frac{1}{\sqrt{n}}. \quad (5.4)$$

Note that  $\rho_{\lambda}$  is independent of  $\mu$ . Thus, by Lemma 4.5, there exist a subsequence  $\{u_n\}$  and  $u^- \in X$  such that  $u_n \rightarrow u^-$  in  $X$  as  $n \rightarrow \infty$ ,  $J_{\mu,\lambda}(u^-) = \overline{\alpha}_{\lambda}(\mu) < 0$  and  $J'_{\mu,\lambda}(u^-) = 0$ . By (5.4) and the fact that  $v_n \geq 0$ , we conclude that  $u^- \geq 0$  a.e. in  $\mathbb{R}^3$ . It follows from  $J_{\mu,\lambda}(u^-) < 0$  and the strong maximum principle that  $u^- > 0$  in  $\mathbb{R}^3$ . This proof is complete.  $\square$

## 6 Concentration for Solutions

In this section, we follow the argument in [39] to study the asymptotic behavior of positive solutions of system  $(SP_{\mu,\lambda})$ . The results of Theorem 1.3 will hold by the following two theorems.

**Theorem 6.1.** *Let  $u_{\mu}$  be the solutions obtained in Theorem 5.1. Then there exists  $u_{\infty} \in H_0^1(\Omega)$  such that  $u_{\mu} \rightarrow u_{\infty}^+$  in  $X$  as  $\mu \rightarrow \infty$  and it is a positive solution of Eq.  $(S_{\infty,\lambda})$  with  $J_{\infty,\lambda}(u_{\infty}^+) > 0$ , where*

$$J_{\infty,\lambda}(u) := \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx - \frac{\lambda}{2} \int_{\Omega} f u^2 dx - \frac{1}{p} \int_{\Omega} g |u|^p dx + \frac{1}{4} \int_{\Omega} K \phi_u u^2 dx.$$

*Proof.* For any sequence  $\mu_n \rightarrow \infty$ , let  $u_n := u_{\mu_n}$  be the positive solutions of Eq.  $(S_{\mu_n,\lambda})$  obtained in Theorem 5.1 for  $\lambda$ . From the proof of Theorem 5.1, there exists a constant  $D_0 > 0$  such that  $\|u_n\|_{\mu_n} < D_0$  for all  $n$ . Then by Lemma 3.1, there exist a subsequence  $\{u_n\}$  and  $u_{\infty}^+ \in H_0^1(\Omega)$  such that  $u_n \rightarrow u_{\infty}^+$  in  $X$  and  $u_n \rightarrow u_{\infty}^+$  in  $L^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ . Since  $\langle J'_{\mu_n,\lambda}(u_n), \varphi \rangle = 0$  for any  $\varphi \in C_0^{\infty}(\Omega)$ , we have

$$\int_{\Omega} \nabla u_n \nabla \varphi dx - \lambda \int_{\Omega} f u_n \varphi dx - \int_{\Omega} g |u_n|^{p-2} u_n \varphi dx + \int_{\Omega} K \phi_{u_n} u_n \varphi dx = 0 \quad (6.1)$$

which implies that

$$\int_{\Omega} \nabla u_{\infty}^+ \nabla \varphi dx - \lambda \int_{\Omega} f u_{\infty}^+ \varphi dx - \int_{\Omega} g |u_{\infty}^+|^{p-2} u_{\infty}^+ \varphi dx + \int_{\Omega} K \phi_{u_{\infty}^+} u_{\infty}^+ \varphi dx = 0. \quad (6.2)$$

That is,  $u_{\infty}^+$  is a weak solution of Eq.  $(S_{\infty, \lambda})$ .

Next, we show that  $u_n \rightarrow u_{\infty}^+$  in  $X$ . Since  $\langle J'_{\mu_n, \lambda}(u_n), u_n \rangle = \langle J'_{\mu_n, \lambda}(u_n), u_{\infty}^+ \rangle = 0$ , we have

$$\|u_n\|_{\mu_n}^2 - \lambda \int_{\mathbb{R}^3} f u_n^2 dx - \int_{\mathbb{R}^3} g |u_n|^p dx + \int_{\mathbb{R}^3} K \phi_{u_n} u_n^2 dx = 0 \quad (6.3)$$

and

$$\int_{\Omega} \nabla u_n \nabla u_{\infty}^+ dx - \lambda \int_{\Omega} f u_n u_{\infty}^+ dx - \int_{\Omega} g |u_n|^{p-2} u_n u_{\infty}^+ dx + \int_{\Omega} K \phi_{u_n} u_n u_{\infty}^+ dx = 0. \quad (6.4)$$

By the fact that  $u_n \rightarrow u_{\infty}^+$  in  $L^r(\mathbb{R}^3)$  for  $2 \leq r < 6$ , we have

$$\lim_{n \rightarrow \infty} \left( \int_{\mathbb{R}^3} f u_n^2 dx - \int_{\Omega} f u_n u_{\infty}^+ dx \right) = 0, \quad \lim_{n \rightarrow \infty} \left( \int_{\mathbb{R}^3} g |u_n|^p dx - \int_{\Omega} g |u_n|^{p-2} u_n u_{\infty}^+ dx \right) = 0$$

and

$$\lim_{n \rightarrow \infty} \left( \int_{\mathbb{R}^3} K \phi_{u_n} u_n^2 dx - \int_{\Omega} K \phi_{u_n} u_n u_{\infty}^+ dx \right) = 0.$$

Thus,

$$\lim_{n \rightarrow \infty} \|u_n\|_{\mu_n}^2 = \lim_{n \rightarrow \infty} \int_{\Omega} \nabla u_n \nabla u_{\infty}^+ dx = \int_{\Omega} |\nabla u_{\infty}^+|^2 dx \geq \lim_{n \rightarrow \infty} \|u_n\|^2$$

which implies  $u_n \rightarrow u_{\infty}^+$  in  $X$ , and thus we have  $J_{\mu_n, \lambda}(u_n) \rightarrow J_{\infty, \lambda}(u_{\infty}^+)$ . Moreover, by (5.1), we can deduce that  $\eta_{\lambda} \leq J_{\mu_n, \lambda}(u_n) \leq d_{\lambda}$  for all  $n$ , which implies that  $\eta_{\lambda} \leq J_{\infty, \lambda}(u_{\infty}^+) \leq d_{\lambda}$ . Hence we conclude that  $u_{\infty}^+ > 0$ . This proof is complete.  $\square$

**Theorem 6.2.** *Let  $u_{\mu}$  be the solutions obtained in Theorem 5.2. Then there exists  $u_{\infty}^- \in H_0^1(\Omega)$  such that  $u_{\mu} \rightarrow u_{\infty}^-$  in  $X$  as  $\mu \rightarrow \infty$  and it is a positive solution of Eq.  $(S_{\infty, \lambda})$  with  $J_{\infty, \lambda}(u_{\infty}^-) < 0$ .*

*Proof.* This proof is essential same as that of Theorem 6.1. Hence we omit it here.  $\square$

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## References

- [1] Benci V, Fortunato D. An eigenvalue problem for the Schrödinger-Maxwell equations. *Topol Methods Nonlinear Anal.* 1998; **11**:283-293.
- [2] D'Aprile T, Mugnai D. Solitary waves for nonlinear Klein-Gordon-Maxwell and Schrödinger-Maxwell equations. *Proc Roy Soc Edinburgh Sect A.* 2004; **134**:893-906.
- [3] Ambrosetti A, Ruiz D. Multiple bound states for the Schrödinger-Poisson problem. *Commun Contemp Math.* 2008; **10**:391-404.

- [4] Azzollini A, Pomponio A. Ground state solutions for the nonlinear Schrödinger-Maxwell equations. *J Math Anal Appl.* 2008; **345**:90-108.
- [5] Ianni I, Vaira G. Non-radial sign-changing solutions for the Schrödinger-Poisson problem in the semiclassical limit. *NoDEA Nonlinear Differ Equ Appl.* 2015; **22**:741-776.
- [6] Kikuchi H. On the existence of a solution for elliptic system related to the Maxwell-Schrödinger equations. *Nonlinear Anal.* 2007; **67**:1445-1456.
- [7] Ruiz D. The Schrödinger-Poisson equation under the effect of a nonlinear local term. *J Funct Anal.* 2006; **237**:655-674.
- [8] Ambrosetti A. On Schrödinger-Poisson Systems. *Milan J Math.* 2008; **76**:257-274.
- [9] Cerami G, Vaira G. Positive solution for some non-autonomous Schrödinger-Poisson systems. *J Differ Equ.* 2010; **248**:521-543.
- [10] Vaira G. Ground states for Schrödinger-Poisson type systems. *Ricerche Math.* 2011; **60**:263-297.
- [11] Huang L, Rocha EM, Chen J. Two positive solutions of a class of Schrödinger-Poisson system with indefinite nonlinearity. *J Differ Equ.* 2013; **255**:2463-2483.
- [12] Chen J. Multiple positive solutions of a class of nonautonomous Schrödinger-Poisson systems. *Nonlinear Anal.* 2015; **21**:13-26.
- [13] Huang L, Rocha EM, Chen J. On the Schrödinger-Poisson system with a indefinite nonlinearity. *Nonlinear Anal.* 2016; **28**:1-19.
- [14] Mercuri C. Positive solutions of nonlinear Schrödinger-Poisson systems with radial potentials vanishing at infinity. *Rend Lincei-Matematicae Applicazioni* 2008; **19**:211-227.
- [15] Mugnai D. The Schrödinger-Poisson system with positive potential. *Commun Partial Differ Equ.* 2011; **36**:1099-1117.
- [16] Shen Z, Han Z. Multiple solutions for a class of Schrödinger-Poisson system with indefinite nonlinearity. *J Math Anal Appl.* 2015; **426**:839-854.
- [17] Sun J, Chen H, Nieto JJ. On ground state solutions for some non-autonomous Schrödinger-Poisson systems. *J Differ Equ.* 2012; **252**:3365-3380.
- [18] Sun J, Wu TF, Feng Z. Multiplicity of positive solutions for a nonlinear Schrödinger-Poisson system. *J Differ Equ.* 2016; **260**:586-627.
- [19] Wang Z, Zhou H. Positive solution for a nonlinear stationary Schrödinger-Poisson system in  $\mathbb{R}^3$ . *Discrete Contin Dyn Syst.* 2007; **18**:809-816.
- [20] Zhao L, Zhao F. On the existence of solutions for the Schrödinger-Poisson equations. *J Math Anal Appl.* 2008; **346**:155-169.
- [21] Alama S, Tarantello G. On semilinear elliptic equations with indefinite nonlinearities. *Calc Var Partial Differ Equ.* 1993; **1**:439-475.
- [22] Chabrowski J, Costa DG. On a class of Schrödinger-Type equations with indefinite weight functions. *Commun Partial Differ Equ.* 2008; **33**:1368-1394.

- [23] Costa DG, Tehrani H. Existence of positive solutions for a class of indefinite elliptic problems in  $\mathbb{R}^3$ . *Calc Var Partial Differ Equ.* 2001; **13**:159-189.
- [24] Bartsch T, Wang ZQ. Existence and multiplicity results for superlinear elliptic problems on  $\mathbb{R}^3$ . *Commun Partial Differ Equ.* 1995; **20**:1725-1741.
- [25] Jiang Y, Zhou H. Schrödinger-Poisson system with steep potential well. *J Differ Equ.* 2011; **251**:582-608.
- [26] Zhao L, Liu H, Zhao F. Existence and concentration of solutions for the Schrödinger-Poisson equations with steep well potential. *J Differ Equ.* 2013; **255**:1-23.
- [27] Du M, Tian L, Wang J, Zhang F. Existence and asymptotic behavior of solutions for nonlinear Schrödinger-Poisson systems with steep potential well. *J Math Phys.* 2016; **57**:031502.
- [28] Sun J, Wu TF. On the nonlinear Schrödinger-Poisson systems with sign-changing potential. *Z Angew Math Phys.* 2015; **66**:16489-1669.
- [29] Sun J, Wu TF, Wu Y. Existence of nontrivial solution for Schrödinger-Poisson systems with indefinite steep potential well. *Z Angew Math Phys.* 2017; **68**:1-22.
- [30] Ye Y, Tang C. Existence and multiplicity of solutions for Schrödinger-Poisson equations with sign-changing potential. *Calc Var Partial Differ Equ.* 2015; **53**:383-411.
- [31] Zhang X, Ma S. Multi-bump solutions of Schrödinger-Poisson equations with steep potential well. *Z Angew Math Phys.* 2015 **66**:1615-1631.
- [32] Zhang W, Tang X, Zhang J. Existence and concentration of solutions for Schrödinger-Poisson system with steep potential well. *Math Meth Appl Sci.* 2016; **39**:2549-2557.
- [33] Lions PL. The concentration-compactness principle in the calculus of variations. The locally compact case part I. *Ann Inst Henri Poincaré, Anal Non Linéaire* 1984; **1**:109-145.
- [34] Willem M. *Minimax Theorems*. Boston: Birkhäuser; 1996.
- [35] Wu TF. On a class of nonlocal nonlinear Schrödinger equations with potential well. *Advances in Nonlinear Analysis* 2020; **9**:665-689.
- [36] Ambrosetti A, Rabinowitz PH. Dual variational methods in critical point theory and applications. *J Funct Anal.* 1973; **14**:349-381.
- [37] Brézis H, Lieb E. A relation between point convergence of functions and convergence of functionals. *Proc Amer Math Soc.* 1983; **88**:486-490.
- [38] Ekeland I. *Convexity methods in Hamiltonian mechanics*. Springer, 1990.
- [39] Bartsch T, Pankov A, Wang ZQ. Nonlinear Schrödinger equations with steep potential well. *Commun Contemp Math.* 2001; **3**:549-569.