

Assimilation of both column- and layer-integrated dust opacity observations in the Martian atmosphere

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Key Points:

- Assimilation of atmospheric measurements of Mars into a global circulation model is extended to include limb profiles of dust opacity.
- Combining nadir and limb profiles of dust opacity enables more accurate recovery of dust vertical structure, including elevated layers.
- The climate reanalysis is significantly improved, as shown in comparisons with independent observations and the Mars Climate Database.

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Abstract

A new dust data assimilation scheme has been developed for the UK version of the Laboratoire de Météorologie Dynamique (LMD) Martian General Circulation Model. The Analysis Correction scheme (adapted from the UK Met Office) is applied with active dust lifting and transport to analyze measurements of temperature, and both column-integrated dust optical depth (CIDO), τ_{ref} (rescaled to a reference level), and layer-integrated dust opacity (LIDO). The results are shown to converge to the assimilated observations, but assimilating either of the dust observation types separately does not produce the best analysis. The most effective dust assimilation is found to require both CIDO and LIDO observations, especially for Mars Climate Sounder (MCS) data that does not access levels close to the surface. The resulting full reanalysis improves the agreement with both in-sample assimilated CIDO and LIDO data and independent observations from outside the assimilated dataset. It is thus able to capture previously elusive details of the dust vertical distribution, including elevated detached dust layers that have not been captured in previous reanalyses. Verification of this reanalysis has been carried out under both clear and dusty atmospheric conditions during Mars Years 28 and 29, using both in-sample and out of sample observations from orbital remote sensing and contemporaneous surface measurements of dust opacity from the Spirit and Opportunity landers. The reanalysis was also compared with a recent version of the Mars Climate Database (MCD v5), demonstrating generally good agreement though with some systematic differences in both time mean fields and day-to-day variability.

Plain Language Summary

Data assimilation is a method of combining atmospheric observations, which are inevitably uncertain and incomplete in their coverage, with a global numerical model. It is commonly used for the Earth to reconstruct the best possible record of the changing global climate. This has also been used for the Martian atmosphere in recent years, using measurements of temperature, dust and ice from satellites in orbit around Mars. But most previous efforts have only used measurements of the total amount of dust in a vertical column from instruments that “look” vertically downwards to the Martian surface. In new work presented here, however, we also use detailed measurements of the vertical structure of the dust distribution from an instrument that “looks” towards the edge of the planet. This is much more effective when atmospheric dust is not mainly concentrated near the ground. Such events are reasonably common on Mars, when elevated layers of dust are formed, which can strongly affect how the atmosphere is heated by the Sun. We present examples of situations when previous methods failed to recover the correct dust distribution, as verified against independent measurements e.g. from the Spirit and Opportunity Rovers, and compare with the ESA Mars Climate Database.

1 Introduction

The dust cycle is a key component of the Martian climate, and is extremely important for understanding the interannual, seasonal and synoptic evolution of the Martian environment. (e.g. Newman et al., 2002a; Kahre et al., 2017, and references therein). Intensive measurements of atmospheric temperature and dust extending over more than ten Mars years (MY) now exist with unprecedented spatial coverage, thanks to various orbital spacecraft. Such observations have already helped to improve our understanding of Mars’ weather and climate. However, the incomplete coverage of these measurements across the planet constrains our ability to study the general circulation in full detail, particularly those aspects related to dust opacity. For instance, the Thermal Emission Imaging System (THEMIS) carried by the Mars Odyssey (MO) spacecraft can provide multi-annual measurements of Column Integrated Dust Opacity (CIDO), but its coverage in space and time is quite limited.

68 On the other hand, numerical models provide four-dimensional simulated data with
69 moderate-high temporal and spatial resolution and complete coverage in space and time,
70 but often fail to reproduce the dust cycle’s full range of variability. Various authors, start-
71 ing with Newman et al. (2002b) (see also Kahre et al. (2017) for a recent review) showed
72 that a global circulation model (GCM) could capture the onset and growth of regional
73 dust events, but did not realistically capture the observed interannual variability. In par-
74 ticular, they could not reproduce the relatively “quiet” year of dust activity that occurs
75 immediately after a simulated global dust storm (GDS) year. Others have sought to take
76 additional factors into account, such as the finite extent of the surface dust reservoir (e.g.
77 Pankine & Ingersoll, 2004; Szwasz et al., 2006) or nonlinear effects associated with the
78 “shadowing” of pockets of dust behind rocks and boulders (Mulholland et al., 2013). But
79 even the most sophisticated free-running GCMs still struggle to capture realistic inter-
80 annual variability associated with dust lifting and transport.

81 To aid this task, data assimilation has become an optimal approach to provide a
82 solution that is consistent with both observations and modelled physical constraints. Data
83 assimilation corrects model-predicted variables towards observations such that the re-
84 sulting solution can represent the full observed variability of the climate. This approach
85 has been widely used as an effective tool in operational weather forecasting systems or
86 climate models for analyzing meteorological variables for the Earth (e.g. Lorenc et al.,
87 1991; Kalnay, 2003). This approach has already been used for a number of years to in-
88 vestigate tracer/chemical evolution in the Earth’s atmosphere (Collins et al., 2001; J. Wang
89 et al., 2004; Schutgens et al., 2010; Benedetti et al., 2018). Collins et al. (2001), for ex-
90 ample, used an optimal interpolation approach to assimilate satellite retrievals of total
91 column aerosol optical depth (AOD) over the Indian Ocean, which reproduced the daily
92 variations of AOD at a single model grid point. J. Wang et al. (2004) used nudging to
93 assimilate AOD into a nonhydrostatic atmospheric model, which captured the observed
94 evolution of a dust event near Puerto Rico. More recently, Schutgens et al. (2010) ap-
95 plied the Local Ensemble Transform Kalman filter to assimilate AOD from the AERONET
96 global surface observation network, which captured the evolution of AOD and also re-
97 duced uncertainties in model estimates of the evolving aerosol distribution. At the time
98 of writing, around five major operational centres around the world use a variety of as-
99 similation techniques, including optimal analysis (similar to the Analysis Correction scheme
100 presented here), variational methods or ensemble Kalman filters to analyse observations
101 of dust and aerosols from various sources (e.g. Benedetti et al., 2018, and references therein),
102 such as AOD derived from orbiting or surface-based platforms.

103 Few centres have assimilated dust profile observations, however, preferring instead
104 to focus on achieving high horizontal resolution utilizing the much more abundant AOD
105 measurements. Limited publications to date include the work of Yumimoto et al. (2008),
106 who assimilated vertical profiles of the dust extinction coefficients in a regional dust trans-
107 port model. In their study, the data from a ground-based lidar network were interpo-
108 lated to the vertical model levels for analyzing the model prognostic dust variables. More
109 recently, Sekiyama et al. (2010) directly assimilated the total attenuated backscatter-
110 ing coefficient from the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observa-
111 tions (CALIPSO) mission into a global chemistry-transport model. The measurements
112 were averaged approximately to the model horizontal and vertical resolution before the
113 assimilation. This approach has recently been extended by Cheng et al. (2019), who also
114 assimilated CALIPSO profiles of aerosol optical depth using a 4D Ensemble Kalman Fil-
115 ter approach.

116 Data assimilation has also been applied to the Martian atmosphere with success.
117 S. R. Lewis & Read (1995) implemented the analysis correction (AC) scheme (Lorenc
118 et al., 1991) in a simple version of a Mars General Circulation Model (MGCM) in order
119 to assimilate temperature profiles from the Pressure Modulator Infrared Radiome-
120 ter (PMIRR) instrument on-board the short-lived Mars Observer spacecraft (1993). Their

121 results showed that assimilation of such observations was feasible and that it improved
 122 the agreement between model and observations. S. R. Lewis et al. (2007) extended this
 123 approach to include dust tracer assimilation, which was combined with a full MGCM to
 124 assimilate thermal profiles and CIDO rescaled to a reference level (hereafter τ_{ref}) using
 125 retrievals from the Thermal Emission Spectrometer (TES) on-board MGS (M. D. Smith
 126 et al., 2003). The performance of the data assimilation system was validated against in-
 127 dependent radio occultation measurements by Montabone, Lewis, Read, & Hinson (2006).
 128 This showed that combined temperature and τ_{ref} assimilation was able to reduce discrep-
 129 ancies between the model and radio occultation data below 20 km, especially when dust
 130 amounts were large and changing rapidly, although some large discrepancies remained
 131 due to known inconsistencies between TES temperature profiles and radio occultation
 132 data. This approach was further extended by Holmes et al. (2020) to include assimi-
 133 lation of column dust optical depth measurements derived from Mars Reconnaissance Or-
 134 biter (MRO) MCS retrievals, together with measurements of water ice and ozone, into
 135 a version of the LMD/UKMGCM that also advects dust and other tracers with the anal-
 136 ysed winds.

137 An alternative approach to assimilation of Mars observations was developed by Hoff-
 138 man et al. (2010) based on a complementary method using the ensemble Kalman filter
 139 (EnKF, Evensen, 2003) to assimilate TES temperature retrievals. They found generally
 140 improved agreement with TES temperature observations over a free-running model, and
 141 in Greybush et al. (2012) the joint assimilation of TES temperatures with forcing using
 142 a 2D CIDO dust field from TES was shown to improve the agreement of model and TES
 143 temperatures further. This method has also been extended to include assimilation of col-
 144 umn dust optical depths from both MGS/TES and MRO/MCS observations (Greybush
 145 et al., 2019) and a dataset is publicly available known as EMARS.

146 The approach used by S. R. Lewis et al. (2007) first assimilated TES temperature
 147 profiles and τ_{ref} without explicitly advecting the dust tracer field, using an empirical re-
 148 lation (Conrath, 1975) to prescribe the vertical distribution of dust. This system has sub-
 149 sequently been applied in several studies of Martian weather and climate (Montabone
 150 et al., 2005; S. Lewis & Barker, 2005; Montabone, Lewis, Read, & Withers, 2006; S. R. Lewis
 151 et al., 2007; Wilson et al., 2008; Rogberg et al., 2010; S. R. Lewis et al., 2016), and both
 152 a three-year reanalysis covering MY 24-27 using this assimilation system and another
 153 covering much of the MCS period (MY 28-32) have been published (Montabone et al.,
 154 2014; Holmes et al., 2020). Navarro et al. (2014) also assimilated Mars Climate Sounder
 155 (MCS) temperature profiles and modified the dust vertical distribution using its corre-
 156 lation with temperature. However, it is essentially different from the work presented here,
 157 in which dust observations are directly assimilated. The more recent study of data as-
 158 similation issues on Mars by Navarro et al. (2017) is more similar to the present work
 159 in including some cases that assimilated MCS dust opacity profiles using the EnKF method.
 160 Their study indicated some promise for this approach, although they only analysed a part
 161 of MY 29 and noted some difficulties in capturing the diurnal variation in dust vertical
 162 distributions.

163 The data assimilation system developed here is based on the scheme described by
 164 S. R. Lewis et al. (2007). However, that scheme does not assimilate a vertically-resolved
 165 dust distribution, only TES nadir retrievals of CIDO, and the model does not transport
 166 the dust actively. The newly-available dataset from MCS on board MRO (Kleinböhl et
 167 al., 2009) does provide vertically resolved, global measurements of the atmospheric dust
 168 distribution. With this new dataset to hand, here we update the existing data assim-
 169 ilation system to better represent the Martian dust cycle. In later work we will use this
 170 to study the formation and life cycles of regional and global dust storms in detail.

171 Section 2 describes the Mars GCM and current data assimilation scheme, and the
 172 observations of Martian dust are in Section 3. We describe how the assimilation was adapted
 173 in Section 4. Sections 5 and 6 describe verification against in-sample and out-of-sample

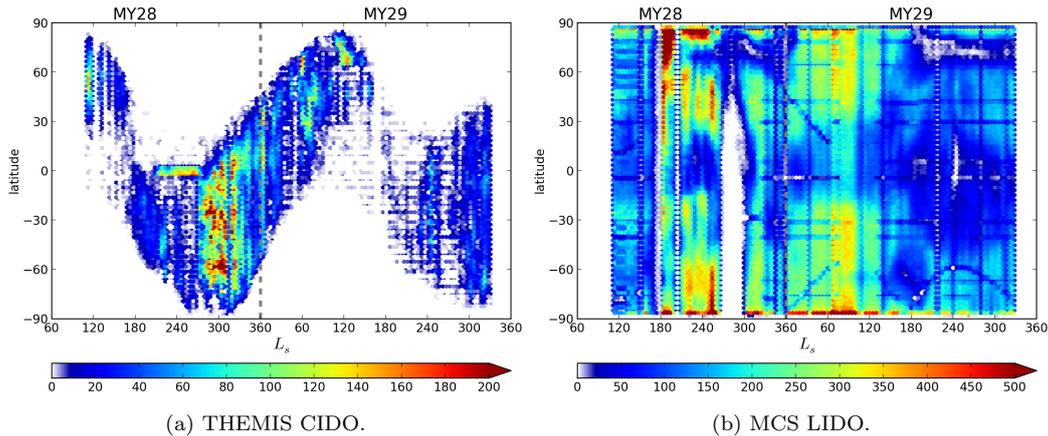


Figure 1: Spatial and temporal distribution of available dust opacity data from THEMIS and MCS during the study period. The colour scales show the number of measurements in $5^\circ L_s$ and 3° latitude bins.

174 observations respectively, while Section 7 describes a systematic comparison of the Mars
 175 Climate Database against the reanalysis. We conclude in Section 8.

176 2 Overview of Mars GCM and data assimilation scheme

177 In this work the model used is based on the UK version of a three-dimensional Martian
 178 Global Climate Model (UK-LMD MGCM, v5.1.3) (Forget et al., 1999; Mulholland et al.,
 179 2013). The model combines a spectral dynamical solver at triangular truncation T31,
 180 corresponding to a 96×48 longitude-latitude grid in real space, a tracer transport scheme
 181 and dust lifting and deposition routines, along with a full range of physical parameter-
 182 izations.

183 The equations for a hydrostatic, adiabatic and inviscid gas surrounding a rotating
 184 spherical planet are cast in vorticity-divergence form. In the vertical, levels are defined
 185 in terms of the terrain-following σ coordinate system using a standard finite difference
 186 approach. There are 25 levels with the first three at 4, 19, and 44 m above the surface,
 187 to resolve detailed surface processes represented in the model. The model top varies in
 188 altitude over time but is typically at around 100 km, with a sponge layer (applying a lin-
 189 ear drag on eddy vorticity and divergence) in the uppermost three levels to reduce spu-
 190 rious reflections of vertically propagating waves. There are typically 480 dynamical and
 191 96 physics timesteps per sol (where a sol is a mean solar day on Mars).

192 The radiative transfer scheme calculates atmospheric absorption and emission due
 193 to carbon dioxide and airborne dust; the radiative effects of water vapour and ice are not
 194 included since our focus here is on the dust cycle. We rely, therefore, on the tempera-
 195 ture assimilation to account for the radiative effects of clouds. The balance between in-
 196 coming radiative flux and thermal conduction in the soil contributes to changes in sur-
 197 face temperature, using a surface thermal inertia field derived from TES and Viking (For-
 198 get et al., 1999) and topography from the Mars Orbiter Laser Altimeter on MGS (D. E. Smith
 199 et al., 2001). The surface roughness length z_0 is based on a global map compiled by Hébrard
 200 et al. (2012), and implemented in the UK-MGCM by Mulholland et al. (2015).

201 The dust transport scheme includes dust lifting parameterizations, tracer advec-
 202 tion, gravitational sedimentation and dry deposition. We assume a $1.5 \mu\text{m}$ particle size

for simplicity, based on Mars Exploration Rover (MER) observations (Lemmon et al., 2004). The two most important distinct mechanisms responsible for the injection of dust into the atmosphere are thought to be dust lifting by near-surface wind stress, and dust lifting by dust devils (Newman et al., 2002a).

The data assimilation scheme is based on the analysis correction sequential estimation (AC) scheme (Lorenç et al., 1991) but with modifications specific to Mars (S. R. Lewis et al., 2007). The assimilation step is computationally inexpensive compared with the rest of the model, and so is performed at each dynamical timestep. S. R. Lewis et al. (2007) describe the scheme in full detail. Temperature assimilations are the same as in that work, except for the observational dataset used. They assimilated dust CIDO observations without advecting the dust tracer, instead setting the vertical distribution of dust opacity using an empirical relation following Conrath (1975). In this work we extend the dust assimilation to incorporate advective transport of radiatively active dust in the simulation model as well as to assimilate both CIDO and LIDO (Layer integrated dust opacity) observations; this is described in Sect. 4.

The ratio of observational error to first guess error used in the normalization factor \tilde{Q}_i (Lorenç et al., 1991, Eq. 3.20) is set to 1 for assimilation of TES temperature observations (as previously done by S. R. Lewis et al., 2007), implying that the model and observation errors are comparable. Following the study of ice opacity assimilation by Steele et al. (2014), we also set this ratio to 1 for the dust observations.

3 Observations of Martian dust

Thanks to various spacecraft in orbit around Mars since 1997, measurements of atmospheric temperature and dust exist covering the Martian atmosphere over more than ten MYs. The instruments on board these spacecraft for determining temperature and dust in the Martian atmosphere that have been used for assimilation, such as the present study, include (amongst others) TES on MGS (M. D. Smith, 2004), THEMIS on MO (M. D. Smith, 2009) and MCS on MRO (Kleinböhl et al., 2009).

TES and THEMIS dust retrievals contain CIDO data only, while MCS data contain more recent satellite observations (MCS v3 was used here for this initial proof of concept) with vertically resolved, asynchronously-sampled global retrievals of atmospheric profiles of temperature, and LIDO (McCleese et al., 2010). This contains information on the day-to-day variability of Martian weather from the near surface to the top of the middle atmosphere around 80 km altitude (Kleinböhl et al., 2009). The spacecraft have different operational periods and orbits, so their retrievals have different temporal and spatial coverage. Only THEMIS has overlapping observational periods with the other two datasets. Further details of the THEMIS dataset, including the retrieval algorithm, can be found in M. D. Smith et al. (2000, 2003) and M. D. Smith (2009).

The analysis in this paper focuses on part of the MCS mapping period from MY28 $L_s = 110^\circ$ (solar longitude) to MY29 $L_s = 330^\circ$. During this period, THEMIS CIDO data and MCS LIDO data are both available. The spatial coverage of THEMIS retrievals varies significantly with L_s , while MCS retrievals are more consistent and uniform except during the global-scale dust storm (GDS) season (roughly from MY28 $L_s = 270^\circ - 305^\circ$; see Fig. 1(b)). Coverage was restricted to the very early stage of the storm and poleward of 45°N throughout. The spatial coverage of these two datasets within the study period is shown in Fig. 1.

Because the dust is assumed by the THEMIS dust opacity retrieval algorithm to be well mixed, the THEMIS CIDO data is commonly reported rescaled (i.e. as τ_{ref}) to a reference pressure of 610 Pa (M. D. Smith, 2009), to remove the effects of variable topography. However, this assumption may introduce uncertainty when the dust is not well

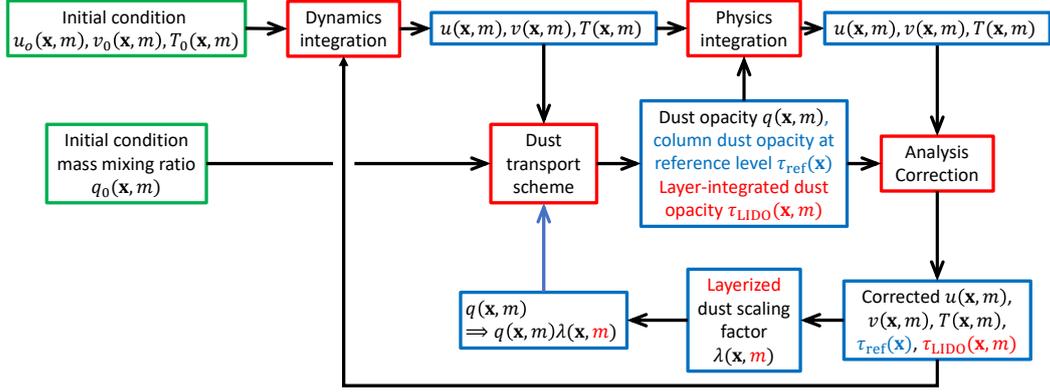


Figure 2: Sequence of operations in the new data assimilation scheme with active dust transport. Green boxes show initial conditions, blue boxes show data generated by MGCM integration, and red boxes show individual MGCM modules. Text in black applies to both CIDO and LIDO assimilation, text in blue applies to CIDO assimilation only, and text in red applies to LIDO assimilation only. Only variables related to the data assimilation scheme are included.

252 mixed. When an intense detached dust layer exists (Heavens et al., 2011), for example,
253 rescaling under the well-mixed assumption could lead to an overestimate in τ_{ref} .

254 It is also important to note that THEMIS dust observations are provided as an in-
255 frared absorption optical depth, while the modeled τ_{ref} is the visible extinction optical
256 depth. Using numerical experiments, M. D. Smith (2009) determined the conversion be-
257 tween IR absorption and visible extinction optical depth to be $\gamma \sim 1.3$, and Clancy et
258 al. (2003) used a scaling factor $\varepsilon = 2$ to convert from IR to visible for dust of size 1.5–
259 2.0 μm . Lemmon et al. (2004) compared visible optical depths with MER measurements
260 at 9 μm wavelength, and found agreement with Clancy et al. (2003). For simplicity, in
261 this work the dust particle size is approximated by a constant 1.5 μm , which is reason-
262 ably consistent with various observational studies (Pollack et al., 1995; Clancy et al., 2003;
263 Lemmon et al., 2004). Hence the conversion factor from THEMIS IR absorption opti-
264 cal depth to a model-compatible visible extinction optical depth is 2.6, and in this work
265 we imply the visible extinction optical depth when referring to CIDO and τ_{ref} , unless
266 otherwise stated. For MCS LIDO retrievals, the infrared opacities (at a wavelength cen-
267 tred on 21.6 μm) were multiplied by a factor of 7.3 to convert them to a visible equiv-
268 alent (Montabone et al., 2015).

269 For a fully-independent validation of the analysis, upward-looking surface obser-
270 vations provide a bottom-up view of CIDO that is independent of satellite-based datasets,
271 although only over particular locations. The MER missions, Spirit (14.57°S, 175.48°E)
272 and Opportunity (1.95°S, 5.53°W) provide almost continuous data coverage during MY28
273 and MY29, concurrent with the study period in this work. Both rovers carried a Pan-
274 cam camera, which included solar filters at 440 nm and 880 nm wavelengths. The effec-
275 tive dust particle radii based on the rovers’ observations were $1.47 \pm 0.21 \mu\text{m}$ for Spirit
276 and $1.52 \pm 0.18 \mu\text{m}$ for Opportunity (Lemmon et al., 2004). However, the 440 nm filter
277 is significantly affected by a red leak (Lemmon et al., 2015). As there should not be a
278 significant difference between the measured CIDO at 880 nm and CIDO at 700 nm (used
279 as the model visible wavelength), we therefore rescale the CIDO measurements at 880 nm
280 to the reference pressure 610 Pa, and make the comparison directly with the modelled
281 sol-averaged τ_{ref} (also rescaled to 610 Pa).

282 About 10% of the THEMIS and MCS data are excluded from the assimilation and
 283 used for out-of-sample validation. The withheld data were taken as every 10th THEMIS
 284 data point and every 10th MCS vertical profile. Withholding this small fraction of the
 285 dataset should not greatly affect the assimilated results. One should not be surprised
 286 that these validation data may be correlated with the assimilated data, which does weaken
 287 their usefulness in validating the model to some extent. However, these datasets were
 288 the best available for the assimilation itself.

289 **4 Dust data assimilation with activated transport**

290 Initial conditions for the prognostic variables and dust tracers are taken from a free-running
 291 spin-up run, which is run for two years prior to the start of the assimilation. Temper-
 292 ature assimilation using MCS data was always included, using the method described by
 293 S. R. Lewis et al. (2007), and such that temperatures were always assimilated before the
 294 dust assimilation. An identical free-running simulation without any assimilation, but with
 295 a fully active dust lifting and transport cycle tuned to reproduce plausible seasonal vari-
 296 ations of dust loading, was run in parallel.

297 Previous assimilation studies using the UK-LMD MGCM excluded active dust trans-
 298 port, instead just correcting the temperature profiles and τ_{ref} . The dust distribution re-
 299 mained static in the absence of dust observations and the vertical distribution was pre-
 300 scribed using the Conrath (1975) distribution.

301 In the new scheme presented here, the data assimilation system is updated to in-
 302 clude full dust transport, lifting and sedimentation while correcting τ_{ref} and/or the model's
 303 vertical dust distribution.

304 **4.1 CIDO assimilation only**

In this configuration only the CIDO retrievals are assimilated. The sequence of opera-
 tions is shown in Fig. 2. First the dynamics timestep is integrated, and then the dust
 is advected to obtain the three dimensional (3D) distribution of dust mass mixing ra-
 tio $q(\mathbf{x}, m)$. The CIDO at position \mathbf{x} , predicted from this distribution (τ_C) is obtained
 by linearly summing up the layer-integrated dust opacity (LIDO) within each model layer
 (τ_{LIDO}). In the model the LIDO at horizontal position \mathbf{x} in model level m is given by

$$\tau_{\text{LIDO}}(\mathbf{x}, m) = q(\mathbf{x}, m) q_{\text{ext}} \frac{p_s(\mathbf{x}) \Delta\sigma(m)}{g} \quad (1)$$

where $q_{\text{ext}} = (3Q_{\text{ext}})/(4\rho r)$, $q(\mathbf{x}, m)$ is the dust mass mixing ratio, $g = 3.72 \text{ m s}^{-2}$ is
 the gravitational acceleration, Q_{ext} is the extinction coefficient, $\rho = 2500 \text{ kg m}^{-3}$ and
 $r = 1.5 \mu\text{m}$ are the density and radius of dust particles, respectively. $p_s(\mathbf{x})$ is the sur-
 face pressure, and $\Delta\sigma(m)$ is the layer thickness. The reference dust opacity τ_{ref} at a ref-
 erence pressure p_{ref} is then determined from the model by

$$\tau_{\text{ref}} = \left(\sum_m \tau_{\text{LIDO}}(\mathbf{x}, m) \right) \frac{p_{\text{ref}}}{p_s(\mathbf{x})} \quad (2)$$

305 The reference pressure p_{ref} is arbitrary; to compare the results with observations, mod-
 306 elled CIDO values are rescaled to 610 Pa.

The advected dust opacity and mass mixing ratio fields are then used to integrate
 the physical parametrizations. Finally, τ_{ref} and T are updated using data assimilated by
 the AC scheme, followed by increments to u and v in thermal wind balance. The dust
 transport scheme transports a 3D dust mass mixing ratio field, but the assimilated CIDO
 dust observations constrain τ_{ref} only, so the dust mass mixing ratio at each model layer
 must be adjusted after the assimilation. This adjustment simply consists of a multiplica-
 tive scale factor $\lambda(\mathbf{x})$, which ensures that the shape of the vertical dust profile at each

horizontal grid point remains the same before and after the assimilation of τ_{ref} :

$$\lambda(\mathbf{x}) = \frac{\tau_{\text{LIDO}}(\mathbf{x}, m)'}{\tau_{\text{LIDO}}(\mathbf{x}, m)} = \frac{\tau_{\text{ref}}'(\mathbf{x})}{\tau_{\text{ref}}(\mathbf{x})} \quad (3)$$

307 where variables without and with primes are before and after assimilation, respectively.
 308 Since the extinction coefficient and layer thickness are constant within a particular time
 309 step, Eq. 1 implies that the dust mass mixing ratio $q(\mathbf{x}, m)$ is proportional to $\tau_{\text{LIDO}}(\mathbf{x}, m)$,
 310 and therefore the adjustment $\lambda(\mathbf{x})$ can be applied directly to $q(\mathbf{x}, m)$. A similar assump-
 311 tion was also used when assimilating AOD on Earth (e.g. Collins et al., 2001; J. Wang
 312 et al., 2004).

313 4.2 LIDO only

314 A more advanced method is required to make proper use of the MCS vertically-resolved
 315 dust profiles. This section describes how retrievals of dust profiles from MCS are assim-
 316 ilated into the model by themselves (i.e., without assimilating CIDO). Figure 2 shows
 317 the procedure for assimilation of LIDO, and it is very similar to CIDO-only, but now LIDO
 318 (τ_{LIDO}) is analyzed directly. When the model layers have a smaller vertical spacing than
 319 the MCS measurements (typical in the lower and middle atmosphere), this approach avoids
 320 the direct interpolation of observational data to the model levels.

321 The assimilation of the vertical dust distribution is not yet widely used in Earth
 322 aerosol modelling or forecasting. Work to date includes Yumimoto et al. (2008), who as-
 323 similated vertical profiles of dust extinction coefficients into a regional dust transport
 324 model. Data from a ground-based LIDAR network were interpolated to the appropri-
 325 ate vertical model levels. Sekiyama et al. (2010) directly assimilated total attenuated backscat-
 326 tering coefficients measured by the CALIPSO spacecraft into a global chemistry-transport
 327 model. In that case they averaged the observations over the model’s horizontal and ver-
 328 tical resolution before assimilation.

329 Since MCS does not (in general) take data at the same levels as those used in the
 330 model, we need to pre-process the observed dust distribution. Our approach differs in
 331 approach from both Yumimoto et al. (2008) and Sekiyama et al. (2010), who both used
 332 the CALIPSO satellite LIDAR measurements of dust opacity in the Earth’s atmosphere.
 333 Here, MCS dust retrievals are reported as dust opacities at atmospheric pressures typ-
 334 ically 1–1.5 km apart, but their intrinsic vertical resolution is about 5 km (Kleinböhl et
 335 al., 2009), so the dataset oversamples the actual MCS measurements. First we integrate
 336 the MCS dust retrievals vertically with a 5 km grid spacing in order to recover the ob-
 337 served LIDO. This ensures that the assimilated data has the same vertical resolution as
 338 the actual measurements. This preserves smaller-scale vertical variability in the mod-
 339 elled dust profile that would be unresolved in the MCS observations.

340 The approach used here resembles the assimilation of thermal profiles into the UK-
 341 LMD MGCM (S. R. Lewis et al., 2007). First, we use the modelled dust opacities $\tau_{\text{LIDO}}(m)$
 342 to predict the dust opacity $\tau_{\text{LIDO}}^{(\text{back})}(i)$ within observation layer i . Model layers that over-
 343 lap more than one observed layer are split linearly in $\ln p$ among the observed layers. Fig-
 344 ure 3 shows an example where three model layers overlap one observed layer. In this in-
 345 stance the modelled dust opacity of the observed layer is

$$\begin{aligned} \tau_{\text{LIDO}}^{(\text{back})}(i) &= \tau_{\text{LIDO}}(m-1) \frac{\ln[p(i)/p(m-1/2)]}{\ln[p(m-3/2)/p(m-1/2)]} \\ &+ \tau_{\text{LIDO}}(m) \\ &+ \tau_{\text{LIDO}}(m+1) \frac{\ln[p(m+1/2)/p(i+1)]}{\ln[p(m+1/2)/p(m+3/2)]} \end{aligned} \quad (4)$$

The LIDO increment within this observation layer is then

$$\Delta\tau_{\text{LIDO}}(i) = \tau_{\text{LIDO}}(i) - \tau_{\text{LIDO}}^{(\text{back})}(i) \quad (5)$$

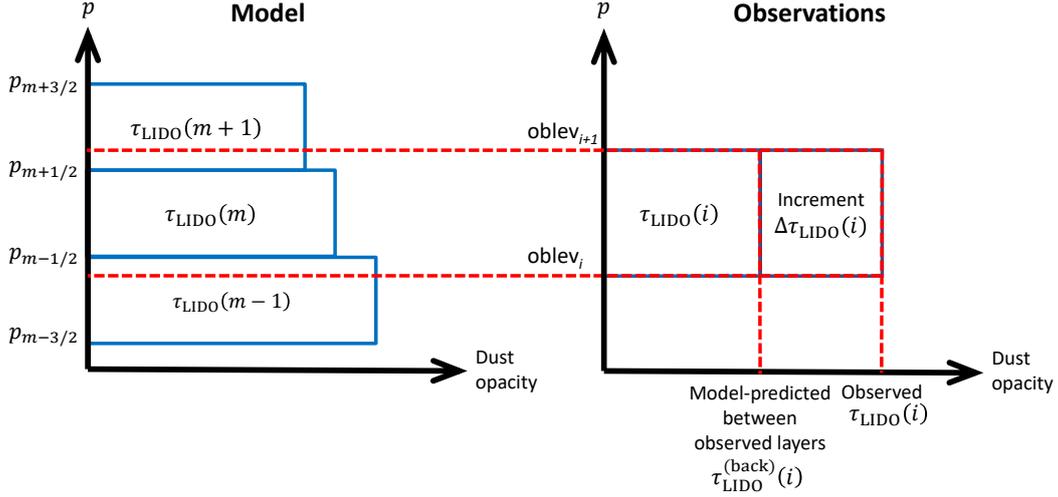


Figure 3: Schematic showing the calculation of the dust opacity increment $\Delta\tau_{\text{LIDO}}(i)$ within observation layer i , given the observed LIDO within that layer $\tau_{\text{LIDO}}(i)$ (right) and the modelled dust opacities $\tau_{\text{LIDO}}(m)$ at the overlapping model levels (left).

346 From which the LIDO increment at each model layer due to observation layer i is

$$\Delta\tau_{\text{LIDO}}(m-1) = \frac{\ln[p(i)/p(m-1/2)]}{\ln[p(i)/p(i+1)]} \Delta\tau_{\text{LIDO}}(i) \quad (6)$$

$$\Delta\tau_{\text{LIDO}}(m) = \frac{\ln[p(m-1/2)/p(m+1/2)]}{\ln[p(i)/p(i+1)]} \Delta\tau_{\text{LIDO}}(i) \quad (7)$$

$$\Delta\tau_{\text{LIDO}}(m+1) = \frac{\ln[p(m+1/2)/p(i+1)]}{\ln[p(i)/p(i+1)]} \Delta\tau_{\text{LIDO}}(i) \quad (8)$$

347 The horizontal assimilation then uses these increments to update the modelled dust field,
 348 following the standard procedure for τ_{ref} .

349 Dust is transported in terms of dust mass mixing ratio, so the assimilation needs
 350 to correct this quantity. As dust mass mixing ratio is proportional to LIDO, it is mul-
 351 tiplied by a factor $\eta = \tau'_{\text{LIDO}}(\mathbf{x}, m)/\tau_{\text{LIDO}}(\mathbf{x}, m)$, where the primed and unprimed quan-
 352 tities are the corrected and uncorrected LIDO values.

353 4.3 Joint CIDO and LIDO

354 To take advantage of both datasets simultaneously, we can assimilate both CIDO and
 355 LIDO together. In principle, one could use the measured dust profiles to correct the dust
 356 in model layers where there are observations, then use the CIDO data to correct the rest
 357 of the column. This avoids unnecessarily adjusting the vertical distribution using CIDO
 358 when part of the distribution has already been corrected using LIDO data. However, THEMIS
 359 and MCS measurements are not normally taken at the same time and place, so it is dif-
 360 ficult to use both simultaneously at one location. Therefore we instead assimilate the
 361 LIDO and CIDO datasets independently. Testing suggested marginally better results when
 362 CIDO was assimilated first, so in the joint assimilation temperature is assimilated first,
 363 followed by CIDO, and finally LIDO.

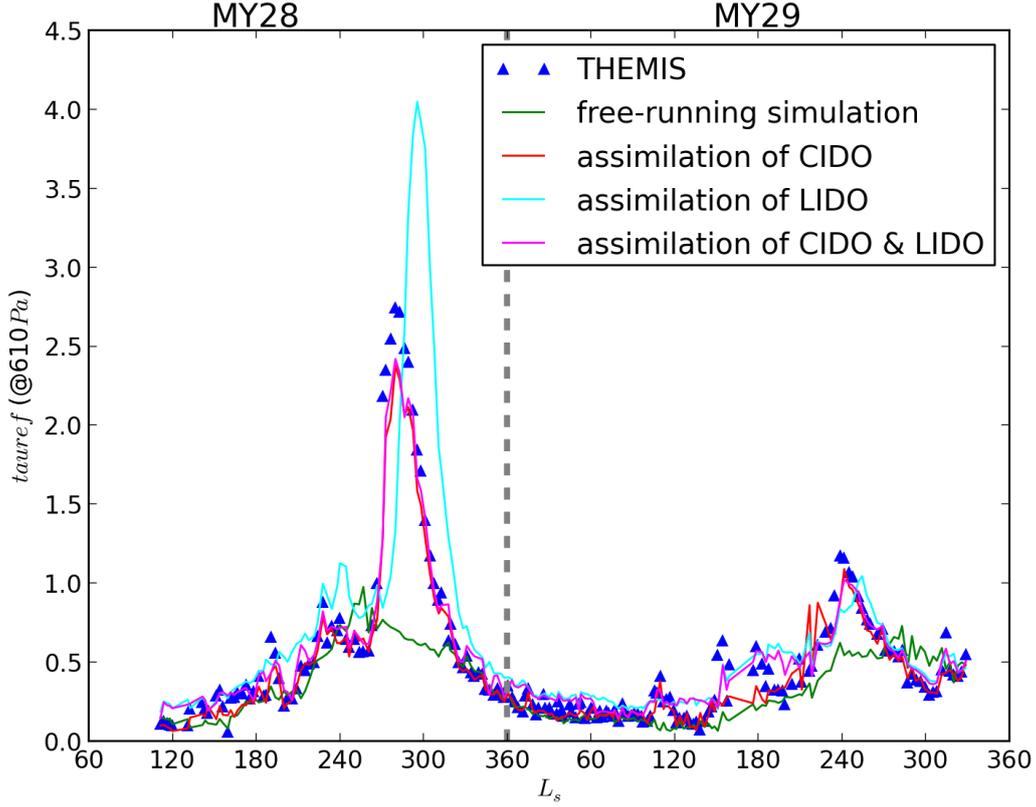


Figure 4: Global mean τ_{ref} (rescaled to 610 Pa) over the study period for the free-running simulation (green), the CIDO-only reanalysis (red), LIDO-only reanalysis (cyan), and joint CIDO/LIDO reanalysis (magenta). THEMIS CIDO observations used in the assimilation are shown as triangles. Each point is an average over five sols.

364 5 Verification I.: in-sample observations and free-running model

365 The methods described above were used to analyse various combinations of THEMIS and
 366 MCS observations obtained during Mars Years 28 and 29, representing a typical pair of
 367 years that include dusty seasons both with and without a planet encircling event. In this
 368 section we present results that compare assimilated analyses with a free-running model
 369 simulation with full dust transport and seasonal variability and evaluate the convergence
 370 of the assimilation towards the input data. Further results and figures can be found in
 371 Section S.1 of the Supplementary Material.

372 The model-predicted and assimilated values of τ_{ref} (rescaled to 610 Pa) from each
 373 variant of the scheme were interpolated to the positions of THEMIS CIDO measurements.
 374 Figure 4 (red line) shows the global mean τ_{ref} of these interpolated data over the course
 375 of the study period. This shows that all three of the reanalyses converge to the assimilated
 376 THEMIS data outside the GDS period in MY28, but in contrast to the other variants,
 377 the LIDO-only assimilation overestimates the peak τ_{ref} during the GDS and mis-
 378 represents the timing of its onset. The free-running simulation, on the other hand, cap-
 379 tures some of the variability, but completely misses the development of the GDS around
 380 MY28 $L_s = 300^\circ$.

5.1 Assimilating CIDO only vs. MCS dust observations

It is also useful to compare the dust reanalysis with the observed time-zonal mean dust distribution. A set of vertical dust distributions (MCS-binned observations hereafter) were produced by sampling MCS dust profiles in 5° horizontal grids during daytime (local time 06:00–18:00) and nighttime (local time 18:00–06:00), after binning the data in $L_s = 5^\circ$ intervals. The model results were interpolated to the same grid and averaged over the same L_s time windows, and restricted to altitudes where MCS-binned observations were available before taking the zonal mean.

A comparison is shown in Fig. 5. With CIDO assimilation, the top of the dust layer when detached dust layers are absent is broadly similar to observations, as is the free-running model (Fig. 5a middle frame). However, elevated detached dust layers (Fig. 5b) cannot be reproduced in either the CIDO-only reanalysis or the free-running model. Such detached dust layers were observed in MCS night-time retrievals (Heavens et al., 2011) and later confirmed by other instruments (M. D. Smith et al., 2013; Guzewich et al., 2013). In the version of the model used in this study, dust tends to be lifted to a lower height than the observed detached dust layer, and it is then well mixed all the way to the ground. Hence a successful reanalysis is likely to require assimilating vertically resolved dust measurements (i.e. LIDO) to reproduce the detached dust layers in a reanalysis.

5.2 Assimilating LIDO only vs. THEMIS dust observations

In the LIDO-only assimilation, the period with a GDS is captured (Fig. 4, cyan line) despite the limited MCS coverage during that period (Fig. 1b). These observations are sufficient for the reanalysis to capture the initial condition and northern boundary condition of the GDS, but the inferred peak in τ_{ref} is higher and later than in the observations.

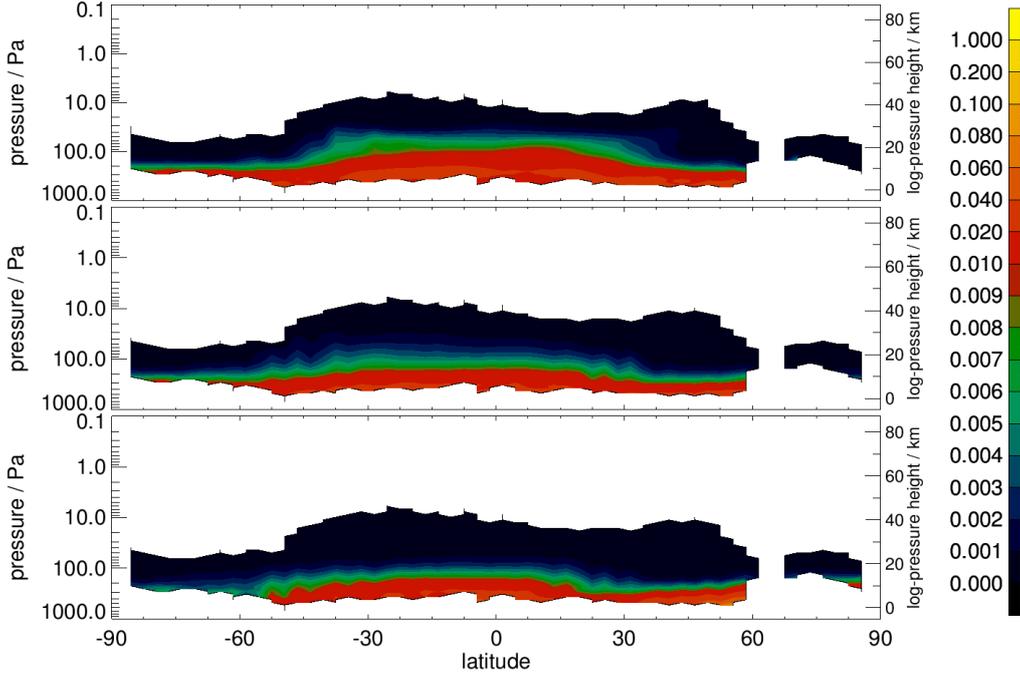
The GDS onset occurs in the southern hemisphere, so the available observations during the dust storm period capture this in the assimilation at a later time than if these observations had been in the southern hemisphere. Once MCS data becomes available in the southern hemisphere again during the “cleanup” of the storm, τ_{ref} returns towards the observed values, but again later than observed in the THEMIS data. The larger peak in the LIDO-only assimilation suggests sedimentation in the model is not efficient enough to remove dust transported southward from the northern boundary of the GDS into the unobserved regions during the peak of the storm.

5.3 Joint CIDO/LIDO assimilation vs. THEMIS and MCS dust observations

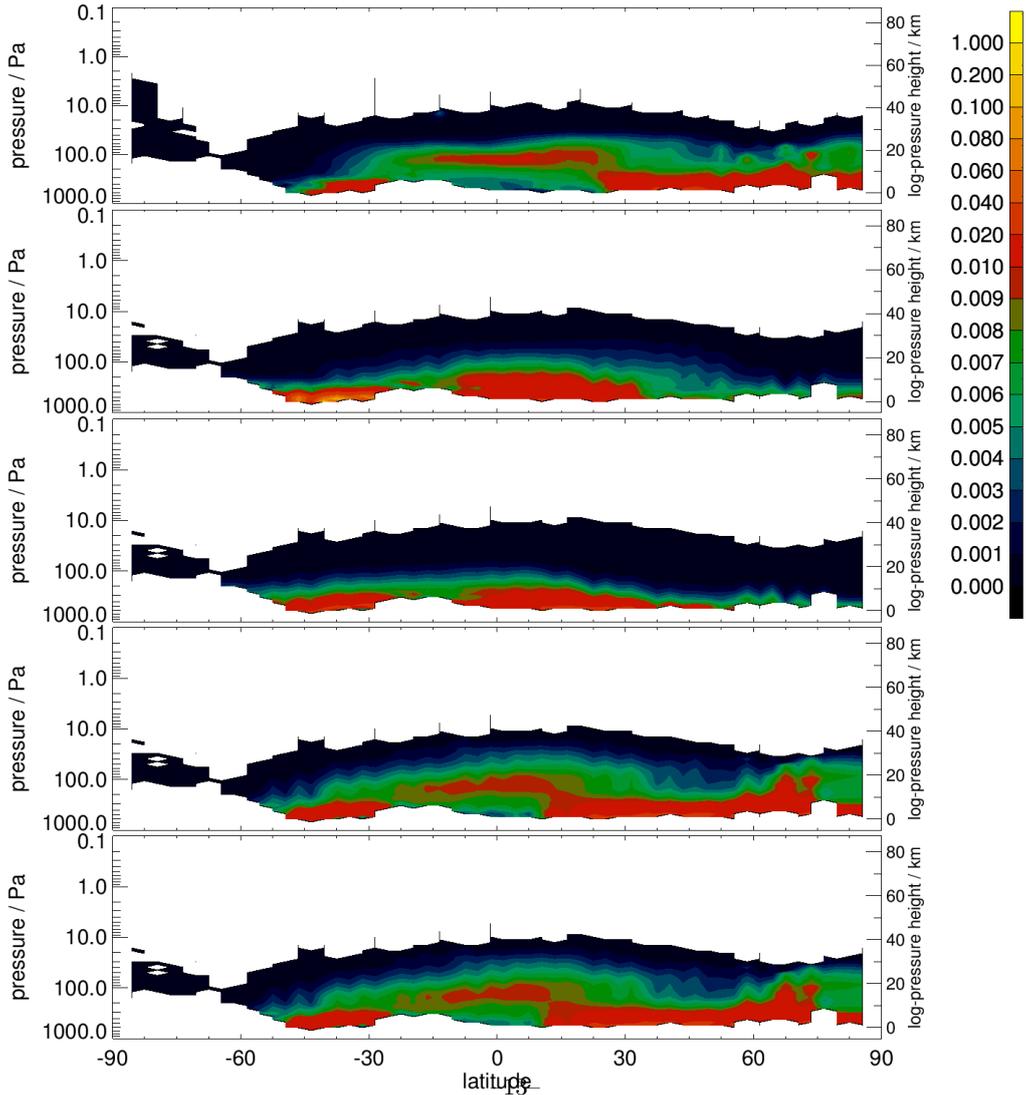
Assimilating CIDO improves the dust horizontal spatial distribution, while assimilating LIDO improves the vertical distribution. By assimilating both we capture features such as the inter-annual variability of the global mean τ_{ref} in the THEMIS observations between GDS and non-GDS years (Fig. 4, magenta line). The global mean τ_{ref} is reproduced during the MY28 GDS, as in the CIDO-only assimilation (red line), and unlike the LIDO-only assimilation (cyan line).

During the “quiet” dust season ($L_s = 0^\circ - 180^\circ$), jointly assimilating CIDO and LIDO gives a reasonable agreement with THEMIS observations (Fig. 4, magenta line). Overestimates in the LIDO-only assimilation during MY29 $L_s = 0^\circ - 90^\circ$ were reduced by assimilating CIDO as well. τ_{ref} is slightly overestimated by the joint assimilation where $\tau_{\text{ref}} < 0.3$ and slightly underestimated where $\tau_{\text{ref}} > 2$.

In the zonal-time mean dust opacity profile compared with MCS binned observations centered at $L_s = 122.5^\circ$ (Fig. 5b, bottom panel), the joint assimilation of CIDO



(a) $L_s = 352.5^\circ$, without detached dust layers. From top: MCS observations, free-running model, CIDO-only reanalysis.



(b) $L_s = 122.5^\circ$, with detached dust layers. From top: MCS observations, free-running model, CIDO-only reanalysis, LIDO-only reanalysis, joint CIDO/LIDO reanalysis.

Figure 5: Night-time (18:00-06:00 local time) zonal-time mean dust opacity (km^{-1}) during MY28 (a) without and (b) with detached dust layers. Time averages are over $5^\circ L_s$.

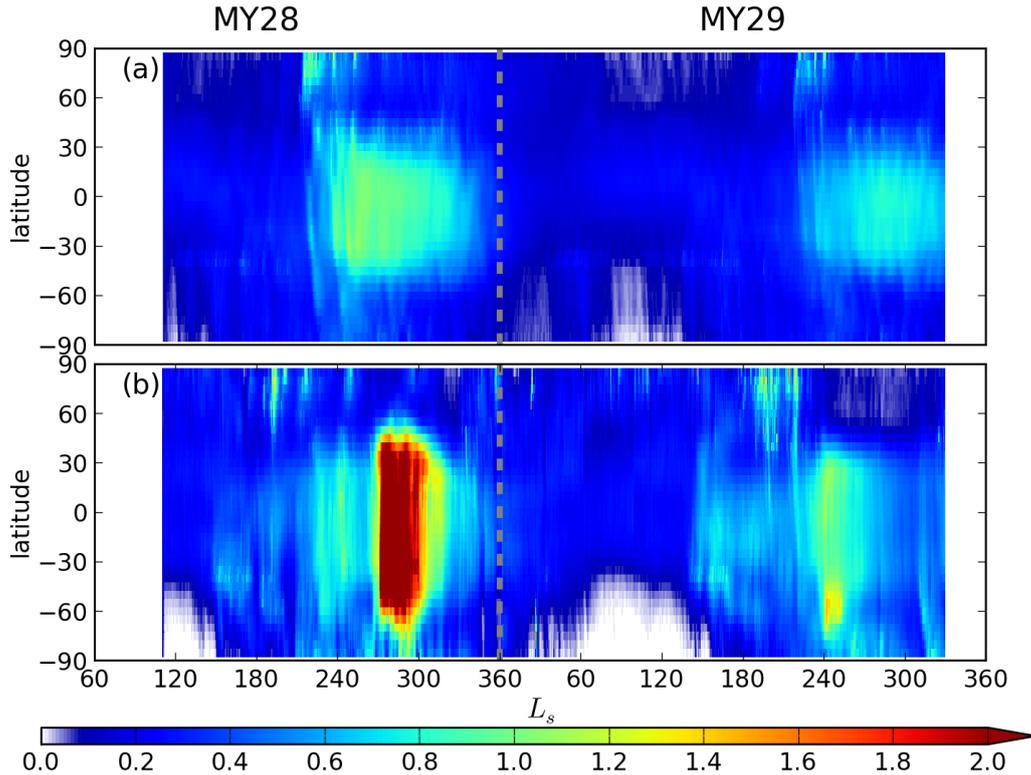


Figure 6: Seasonal evolution of the zonal mean τ_{ref} (CIDO rescaled to 610 Pa), (a) in the free-running simulation, (b) in the joint CIDO/LIDO reanalysis. Because dust is generally well mixed in the lower atmosphere, and hence varies strongly with surface pressure, CIDO is rescaled to the 610 Pa pressure surface to account for Mars' topography.

428 and LIDO produces very similar results to the assimilation of LIDO-only (Fig. 5b, 4th
429 panel).

430 **5.4 Comparison between free-running model and joint CIDO/LIDO re-** 431 **analysis**

432 Figure 6 shows zonal mean column dust opacity τ_{ref} in the free-running simulation
433 and the joint CIDO/LIDO reanalysis over MY28–29. Within each Martian year, the sea-
434 sonal variability of the dust opacity in both free-running simulation and reanalysis ex-
435 hibits at least some features that are generally consistent with spacecraft observations
436 (M. D. Smith, 2008). Global dust opacity is higher during the second half of the year,
437 and relatively quiet during the first half of the year. In the free-running simulation (Fig. 6a),
438 however, the active dust period in each MY lasts longer than in observations (M. D. Smith,
439 2009), while the dust opacity in the reanalysis (Fig. 6b) shows more realistically inter-
440 mittent seasonal variability within each dusty season. The peak in dust opacity is also
441 sharper in L_s in the reanalysis, and tends to shut down prior to the decline in solar forc-
442 ing that occurs towards the end of northern winter.

443 The interannual variability in the reanalysis (Fig. 6b) is essentially the same as in
444 observations (e.g. M. D. Smith (2008, Fig. 8a), M. D. Smith (2009, Fig. 6) and Montabone
445 et al. (2015)). Global dust storms (GDS) do not happen every MY, but the reanalysis
446 successfully reproduces the observed GDS around MY28 $L_s \approx 265^\circ - 310^\circ$. The initi-

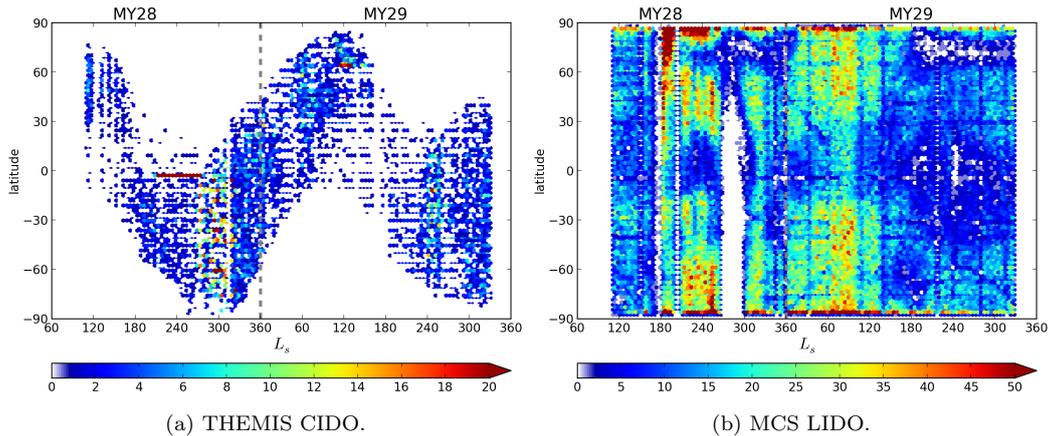


Figure 7: As Fig. 1, but for out-of-sample dust retrievals.

447 ation and duration of the GDS in the reanalysis are also consistent with THEMIS dust
 448 retrievals (M. D. Smith, 2009, Fig. 6 upper panel).

449 The mild dusty season in MY29 following the MY28 GDS also suggests a more re-
 450 realistic interannual variability in the reanalysis. The free-running simulation displays
 451 some variability, with a slightly stronger dusty season in MY28 than in MY29 (Fig. 6a),
 452 but remains some considerable way away from the observations.

453 6 Verification II.: independent, non-assimilated observations

454 In this section, the reanalysis from the new dust assimilation system is validated
 455 against non-assimilated data, including the independent upward-looking measurements
 456 from the Mars Exploration Rovers (MER) “Spirit” and “Opportunity” (Bell III et al.,
 457 2003; Lemmon et al., 2004). In order to have a more comprehensive validation, about
 458 10% of the THEMIS and MCS data were withheld from the assimilation, and they are
 459 also used as an out-of-sample validation. Those withheld data were selected from 1
 460 in every 10 of the data (for THEMIS) and of the profiles (for MCS). It would not be sur-
 461 prising to see that these selected THEMIS and MCS data for validation may have cor-
 462 relation with the data assimilated into the model, and this, to some degree, compromises
 463 their application to validate the reanalysis. The completely independent datasets from
 464 the MER landers, however, provide a complementary way of validation that does not suf-
 465 fer from these correlations. Hereafter, the reanalysis/assimilation used refers to the joint
 466 assimilation of CIDO and LIDO, as described in section 4.3.

467 6.1 Out-of-sample THEMIS dust observations

468 The distribution of non-assimilated (out-of-sample) THEMIS CIDO retrievals is shown
 469 in Fig. 7a. Coverage is similar to the full THEMIS dataset (Fig. 1a), and while it has
 470 only 10% of the data points, its distribution in latitude and L_s reflects the spread in the
 471 full THEMIS dataset. To compare with the out-of-sample THEMIS data, both the ob-
 472 servations and model results were rescaled to the 610 Pa pressure level to account for Mars’
 473 topography, and the model results were interpolated both horizontally and in time to
 474 the out-of-sample data points.

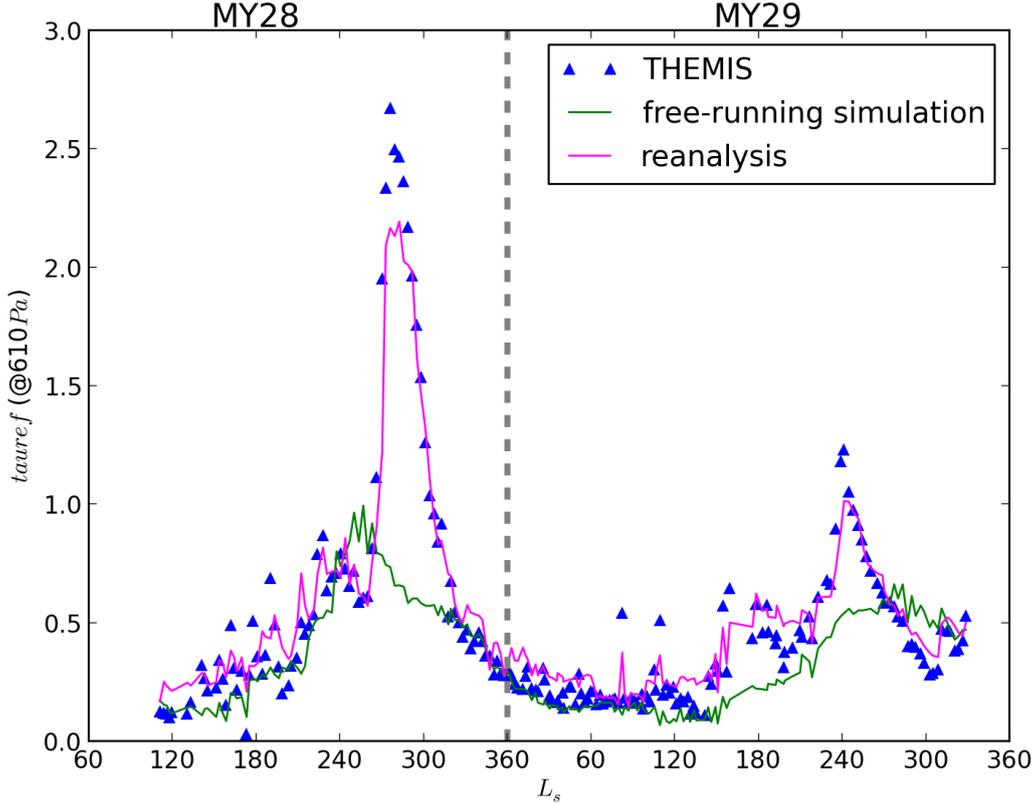


Figure 8: 5-sol global mean τ_{ref} over MY28–29 showing the free-running simulation (green), joint CIDO/LIDO reanalysis (magenta), and out-of-sample THEMIS observations (triangles). Compare Fig. 4 for the in-sample observations.

475 Figure 8 shows the comparison with the global mean τ_{ref} using out-of-sample obser-
 476 vations. The free-running simulation tracks the observations up to $L_s = 240^\circ$ of MY28,
 477 but fails to capture the subsequent GDS. It does predict a marginally milder dust season
 478 in MY29 compared to MY28, as observed, but does not reproduce the observed dust
 479 in either case.

480 The reanalysis performs significantly better, capturing the MY28 GDS as well as
 481 the precursor initiation events and subsequent decay, and the interannual variability dur-
 482 ing MY29’s dusty season. The magnitudes in the reanalysis are also more consistent with
 483 observations than the free-running model, although the maximum τ_{ref} during the MY28
 484 GDS is still lower than observations. Measurement uncertainties in the THEMIS data
 485 may be 20% or higher, however (M. D. Smith, 2004), so the reanalysis could still be broadly
 486 consistent with the observations at the peak of the GDS.

487 During the “quiet” season, the free-running model predictions generally fit the THEMIS
 488 data well, especially during MY29, at least in a global average sense. During this period
 489 both free-running model and observations fall within the minimum observational uncer-
 490 tainty, which is 0.104 for the visible extinction opacity (M. D. Smith, 2009).

491 Correlations between τ_{ref} in the out-of-sample THEMIS observations and the free-
 492 running simulation and reanalysis are shown in Fig. 9. As with the in-sample observa-
 493 tions, the free-running simulation generally underestimates dust loading, mainly in the
 494 dusty season. The reanalysis produces significantly better correlations with out-of-sample

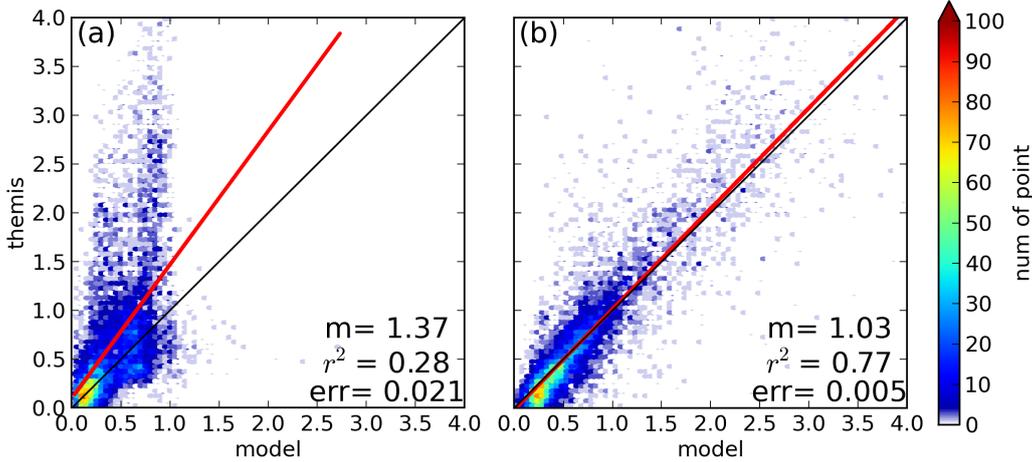


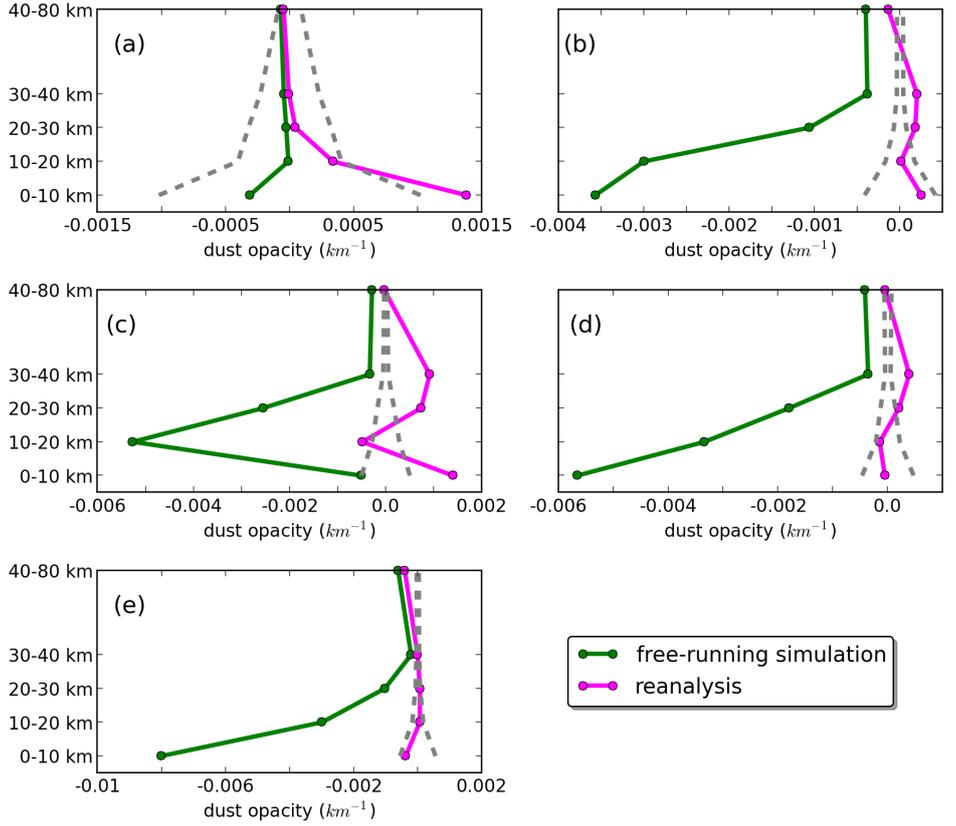
Figure 9: Scatter plots showing individual τ_{ref} points comparing the out of sample THEMIS observations with the free-running model and various reanalyses over the period shown in Fig. 4. Colours show the data density as the number of points per square of side $\tau_{\text{ref}} = 0.05$. Red lines show the linear least square fit, with m the fitting coefficient, r^2 the coefficient of determination, and err the standard error in m . Black lines show $m = 1$. (a) Free-running simulation, (b) joint CIDO/LIDO reanalysis.

495 THEMIS observations. τ_{ref} in the reanalysis only slightly overestimates observations where
 496 $\tau_{\text{ref}} < 0.5$, and slightly underestimates them where $\tau_{\text{ref}} > 2$.

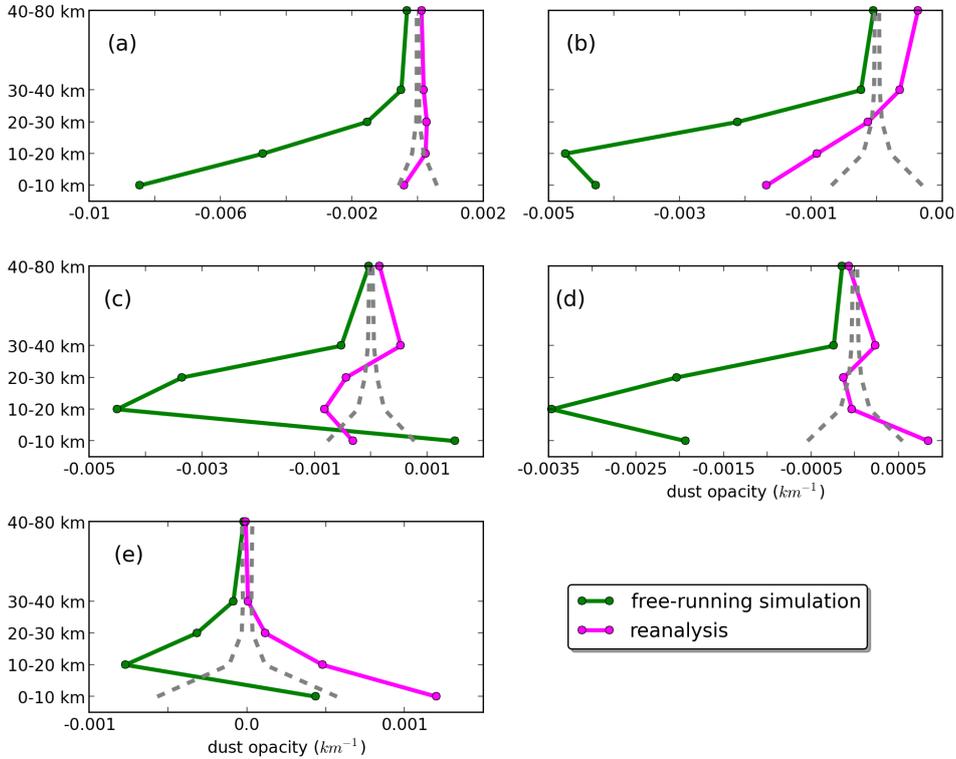
497 6.2 Out-of-sample MCS dust observations

498 Figure 7b shows the distribution of out-of-sample MCS dust profiles. As with THEMIS
 499 there is a similar distribution pattern to the full dataset (Fig. 1b). The model results
 500 were averaged over several pseudo-height ranges (0–10 km, 10–20 km, 20–30 km, 30–40 km,
 501 and 40–80 km), assuming a 10 km scale height and a 610 Pasurface pressure. Within each
 502 pseudo-height range, the mean difference in dust opacity between the out-of-sample data
 503 and the joint CIDO/LIDO reanalysis was calculated for several latitude bands and is shown
 504 in Fig. 10 for the “quiet” and “dusty” seasons.

505 During the “quiet” season (Fig. 10a), in southern high latitudes ($90^\circ\text{S} - 50^\circ\text{S}$)
 506 the free-running model significantly underestimates the dust opacity below 30 km, while
 507 the reanalysis reproduces the observations significantly better. In the midlatitudes of both
 508 hemispheres ($50^\circ\text{S} - 15^\circ\text{S}$ and $15^\circ\text{N} - 50^\circ\text{N}$) the free-running simulation again un-
 509 derestimates the observations, although this underestimate decreases at higher altitude.
 510 The reanalysis generally falls within or close to the observational uncertainty, with the
 511 largest differences at 30–40 km. In the tropics ($15^\circ\text{S} - 15^\circ\text{N}$), the free-running simu-
 512 lation generally underestimates the dust opacity. The reanalysis errors are generally larger
 513 than the MCS observational uncertainties except for 0–10 km. It is an improvement over
 514 the free-running simulation for 10–20 km, 20–30 km, and 40–80 km, but overestimates the
 515 dust opacity for 0–10 km and 30–40 km, with absolute differences larger than those in
 516 the free-running simulation. In northern high latitudes ($50^\circ\text{N} - 90^\circ\text{N}$) the uncertain-
 517 ties in the MCS observations are about 50% larger than elsewhere. The free-running model
 518 falls within the observational uncertainties at all altitudes. The reanalysis also falls within
 519 observational uncertainty except below 10 km.



(a) “Quiet” season ($L_s = 0^\circ - 180^\circ$).



(b) “Dusty” season ($L_s = 180^\circ - 360^\circ$).

Figure 10: Mean difference in dust opacity between the out-of-sample MCS observations and the free-running simulation (green) and joint CIDO/LIDO reanalysis (magenta). In each of the two seasons, the globe is split into latitude bands: (a) $90^\circ\text{S} - 50^\circ\text{S}$, (b) $50^\circ\text{S} - 15^\circ\text{S}$, (c) $15^\circ\text{S} - 15^\circ\text{N}$, (d) $15^\circ\text{N} - 50^\circ\text{N}$, and (e) $50^\circ\text{N} - 90^\circ\text{N}$. Grey dashed lines are the average uncertainties in the MCS observations.

520 Figure 10b shows the same for the dusty season ($L_s = 180^\circ - 360^\circ$). In general,
 521 the reanalysis agrees better with the MCS observations than the free-running simula-
 522 tion. The maximum error in the free-running simulation increases from south to north,
 523 which may be due to the difficulty in predicting frontal dust storms in northern high lat-
 524 itudes during dusty season.

525 In southern high latitudes, the free-running simulation tends to underestimate the
 526 dust opacity above 10 km, and the reanalysis tends to overestimate the dust opacity be-
 527 low 30 km, but fall within the observational uncertainty above 30 km. In southern mid-
 528 latitudes, the reanalysis is closer to the observations than in the tropics and northern
 529 midlatitudes, and at 10–20 km, 20–30 km, and 40–80 km is within or close to observa-
 530 tional uncertainty. In the tropics, the free-running simulation is similar to the northern
 531 middle latitudes, except below 10 km where it slightly overestimates the dust opacity.
 532 The reanalysis tends to underestimate the dust opacity between 10 and 30 km, and over-
 533 estimate the dust opacity above 30 km. In the northern midlatitudes, the reanalysis un-
 534 derestimates the dust opacity below 30 km, but the differences are still smaller than the
 535 free-running simulation. Above 30 km, the reanalysis overestimates the dust opacity with
 536 differences larger than the free-running simulation. In northern high latitudes, the re-
 537 analysis falls either within or close to the MCS observational uncertainties, with lower
 538 differences at higher altitudes.

539 6.3 Independent Pancam observations from Spirit

540 The reanalysis and free-running simulation were compared with Spirit and Opportunity
 541 observations by interpolating τ_{ref} horizontally and in pressure to the rover locations at
 542 Gusev Crater and Meridiani Planum respectively. Figure 11 shows these values at the
 543 locations of the two rovers during MY28–29.

544 Figure 11a shows τ_{ref} at the Spirit rover site. During the relatively “quiet” dust
 545 season Spirit Pancam observations are normally below $\tau_{\text{ref}} = 0.3$. Although the free-
 546 running simulation agreed well with THEMIS τ_{ref} observations globally during this sea-
 547 son (see Fig. 8), at the Spirit landing site it generally underestimates the dust loading.
 548 τ_{ref} only reaches ~ 0.1 during the “quiet” season at this location. During the dusty sea-
 549 son the free-running simulation suggests an increase in τ_{ref} at Spirit’s location, but its
 550 increase does not match the increase in the observations.

551 Conversely, the reanalysis agreed better with the Spirit Pancam data. It captured
 552 the annual and interannual variability in the data well. During the “quiet” season, the
 553 reanalysis reproduces the magnitude and variation in τ_{ref} at the Spirit landing site. Un-
 554 derestimates are mainly during MY29 $L_s = 60^\circ - 120^\circ$. The reanalysis captures the
 555 increase of dust loadings during MY29 $L_s = 140^\circ - 160^\circ$, but not the peak τ_{ref} . The
 556 MY28 GDS is reflected by an increase in τ_{ref} to 3.5 at the Spirit landing site. The re-
 557 analysis reproduces the initiation and decay of the MY28 GDS, and also the variabil-
 558 ity of dust loading during MY29. Nevertheless, the reanalysis dust loading during the
 559 first peak in MY29 ($L_s \approx 160^\circ$) still does not reach the maximum observed by Spirit.
 560 It is worth noting that, although the free-running simulation fails to produce the observed
 561 amount of dust in both dusty seasons, it does exhibit some interannual variability.

562 6.4 Independent Pancam observations from Opportunity

563 Figure 11b shows τ_{ref} at the Opportunity rover site compared with the free-running model
 564 and reanalysis. The observed dust loading at Meridiani Planum has a similar evolution
 565 to that seen at Gusev Crater, but with slightly higher values in general. Both the free-
 566 running simulation and the reanalysis underestimate the peak of the MY28 GDS, though
 567 τ_{ref} in the reanalysis is much closer to the measurements. The reanalysis also better re-
 568 produces the observed variability of dust loading during both dusty seasons, but both

569 the free-running simulation and the reanalysis underestimate the dust loading during the
570 “quiet” season.

571 A similar discrepancy was also noticed in earlier studies when comparing datasets
572 from different instruments. Montabone et al. (2015) found a systematic underestimate
573 of dust opacity over the Opportunity landing site in Meridiani Planum in both TES and
574 THEMIS datasets starting at the spring equinox, up to a factor ~ 2 during northern sum-
575 mer, which may have been linked to the likely presence of clouds. Since the THEMIS
576 data assimilated in this study falls within the same period, it is not surprising to find
577 a similar discrepancy in our reanalysis compared with the Opportunity data. Lemmon
578 et al. (2015) also raised problems with Opportunity 880 nm data around $L_s = 30^\circ -$
579 130° . The source of this discrepancy remains an open problem.

580 7 Validation of the Mars Climate Database against the reanalysis

581 Reanalyses produced by data assimilation are important for verifying models against
582 “reality” or “ground truth”. If the reanalysis is of sufficiently high quality, then it can
583 be used as a surrogate for the real atmosphere when validating and verifying model out-
584 put (e.g. ERA-40, Uppala et al., 2005, for the Earth’s atmosphere). We used the reanal-
585 ysis described in previous sections to verify the Mars Climate Database v5.2 (MCD, S. Lewis
586 et al., 1999; Millour et al., 2015) against our “real” atmosphere.

587 The quantities of interest produced in both the reanalysis and the MCD are sur-
588 face pressure, surface temperature, air temperature, density, and zonal and meridional
589 velocities. We did not compare dust diagnostics, because the reanalysis dust distribu-
590 tion is likely to have more high frequency variability than the MCD, due to the 15-minute
591 timescale over which observations are assimilated; these short time periods are not rep-
592 resented in the MCD, which is forced by dust fields changing over the timescale of a day.
593 There are also significant differences between the way the dust is treated in the MCD
594 and in the GCM used to construct the reanalysis, such as the GCM used here assumes
595 a fixed dust particle radius of $1.5 \mu\text{m}$, while the MCD uses a two-moment scheme which
596 retains information about the full dust particle size distribution (Madeleine et al., 2011).
597 This deficiency of our model will be improved in the future.

598 For each $30^\circ L_s$ period in the reanalysis (MY28 $L_s = 120^\circ$ to the end of MY29)
599 we computed monthly means and day-to-day variability in the same way as the MCD.
600 First we interpolated horizontally from the MCD grid ($5.625^\circ \times 3.75^\circ$) to the reanalysis
601 grid ($5^\circ \times 5^\circ$). Then we interpolated atmospheric quantities linearly in $\log p$ to 30 fixed
602 pressure levels spaced by 2.5 km up to 40 km pseudo-altitude above a reference pressure
603 of 610 Pa (assuming a scale height of 10 km), and spaced by 5 km above that.

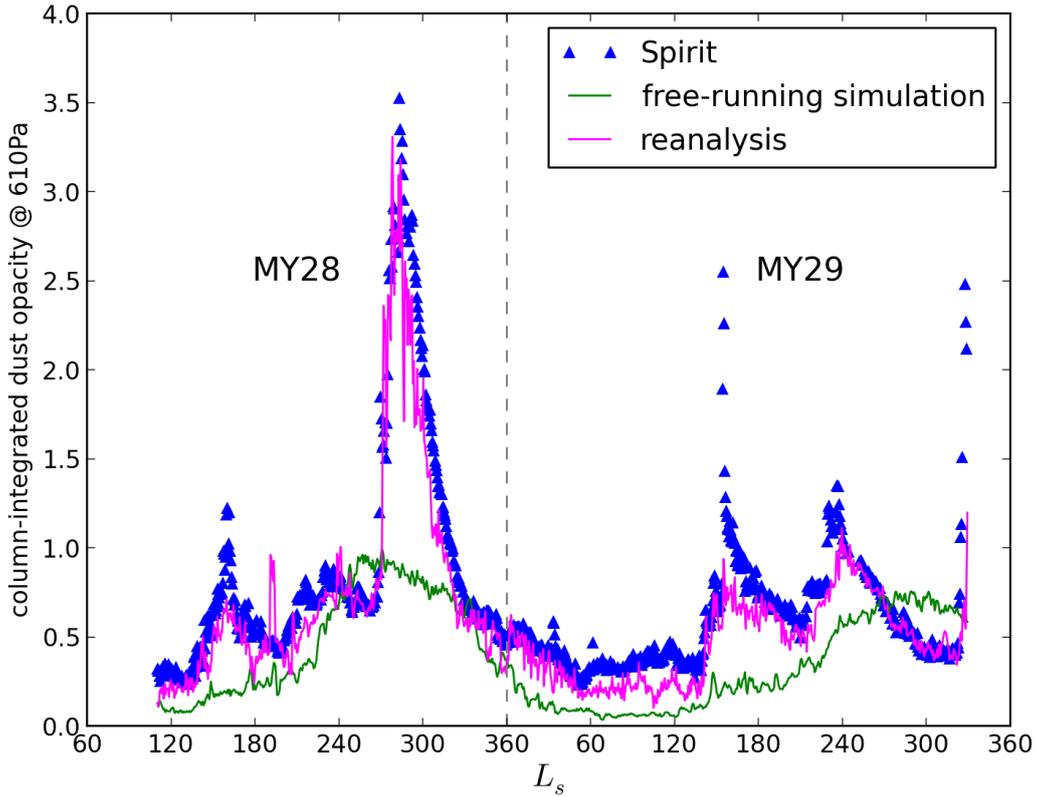
The monthly mean for variable X at (longitude, latitude, pressure) position (i, j, k)
is

$$\overline{X_{ijk}} = \frac{1}{N} \sum_{t=1}^N X_{ijk,t} \quad (9)$$

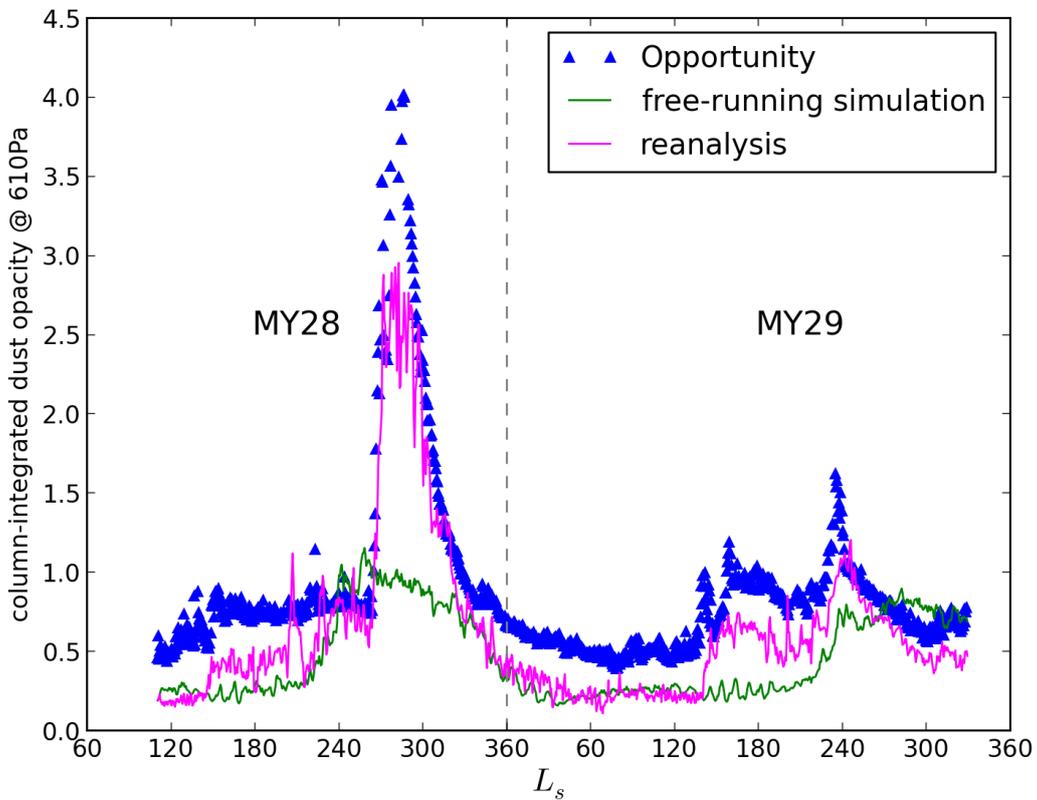
where $t = 1 \dots N$ includes all times within a $30^\circ L_s$ period. Since the orbit is elliptical,
 N varies with season. The day-to-day variability is (Forget et al., 2015, Eq. 5):

$$\text{Var}(X_{ijk}) = \sqrt{\frac{1}{N} \sum_{t=1}^N \left(\overline{X_{ijk,t}^{1 \text{ sol}}} - \overline{X_{ijk,t}^{10 \text{ sols}}} \right)^2} \quad (10)$$

where $\overline{X_{ijk,t}^{1 \text{ sol}}}$ and $\overline{X_{ijk,t}^{10 \text{ sols}}}$ are running means over 1 sol ($t \pm 0.5$ sols) and 10 sols ($t \pm 5$
sols) respectively. The day-by-day variability removes the diurnal cycle, along with any
long-term trend, leaving the variability associated with the day-to-day “weather”. For
atmospheric quantities we calculated zonal-monthly means at (latitude, pressure) points



(a) Spirit, at 14.57°S , 175.48°E (Gusev Crater).



(b) Opportunity, at 1.95°S , 5.53°W (Meridiani Planum).

Figure 11: τ_{ref} from the reanalysis (magenta) and free-running model (green) compared with (a) Spirit and (b) Opportunity Pancam observations. Each point is averaged over one sol.

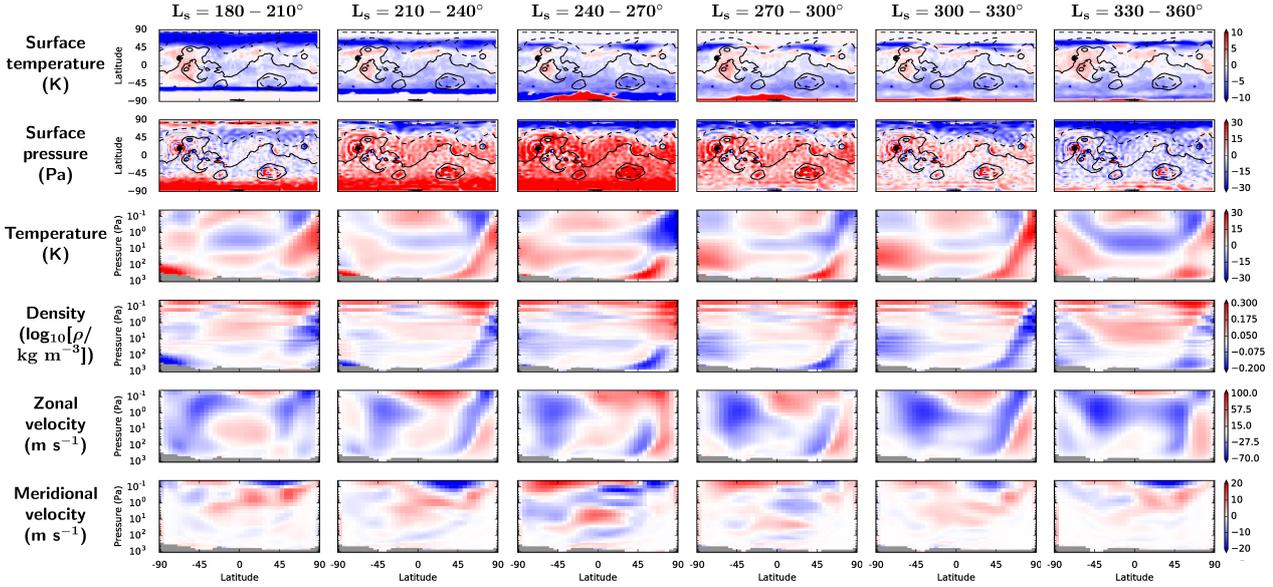


Figure 12: Differences between reanalysis means and MCD means for the period $L_s = 180 - 360^\circ$ in MY29, separated into 30 L_s segments. Positive means the reanalysis value is larger than the MCD value. Surface quantities are monthly means and atmospheric quantities are zonal-monthly means. Grey shows points with missing data (where all points along a latitude circle are below the surface). Black lines show orography at intervals of 4 km between -4 km and +20 km above the geoid (negative values are dashed).

(j, k) :

$$\overline{X}_{jk} = \frac{1}{I} \sum_{i=1}^I \overline{X}_{ijk} \quad (11)$$

604 where I is the number of longitude points above the surface. To compute zonal-monthly
 605 means for the day-to-day variability, we computed the root-mean-square of the day-to-
 606 day variability along each latitude circle. Note this gives us a measure of the day-to-day
 607 variability at points along each latitude circle, rather than the spread of values along the
 608 latitude circle.

609 The MCD contains day-to-day variability as a function of position for each month
 610 and dust scenario, and monthly means as a function of local time of day at zero longi-
 611 tude. To obtain equivalent monthly means for comparison with the reanalysis, we aver-
 612 aged over all local times of day after interpolating each column to the required pres-
 613 sure levels.

614 Figures 12 and 13 summarise the differences between our reanalysis and the MCD
 615 during the second half of the Martian year. Because this paper focuses on dust, we con-
 616 centrate on the dusty season from $L_s = 180 - 360^\circ$. The patterns of mean and vari-
 617 ability differences between the reanalysis and MCD were generally similar in all seasons
 618 (with one exception, discussed below) and for all quantities when comparing MY28 with
 619 MY29, so these summary figures show only MY29. Full sets of figures for all compar-
 620 isons between the reanalysis and MCD for all quantities over all months analysed are in-
 621 cluded as Supplementary Material, Section S.2 and Figs S4–S15.

622 The surface temperature reanalysis was generally cooler than the MCD at most places
 623 and times, particularly between $45-60^\circ$ latitude in each hemisphere. Exceptions were the
 624 polar regions and Tharsis, Arabia, and Elysium, which were persistently warmer in the

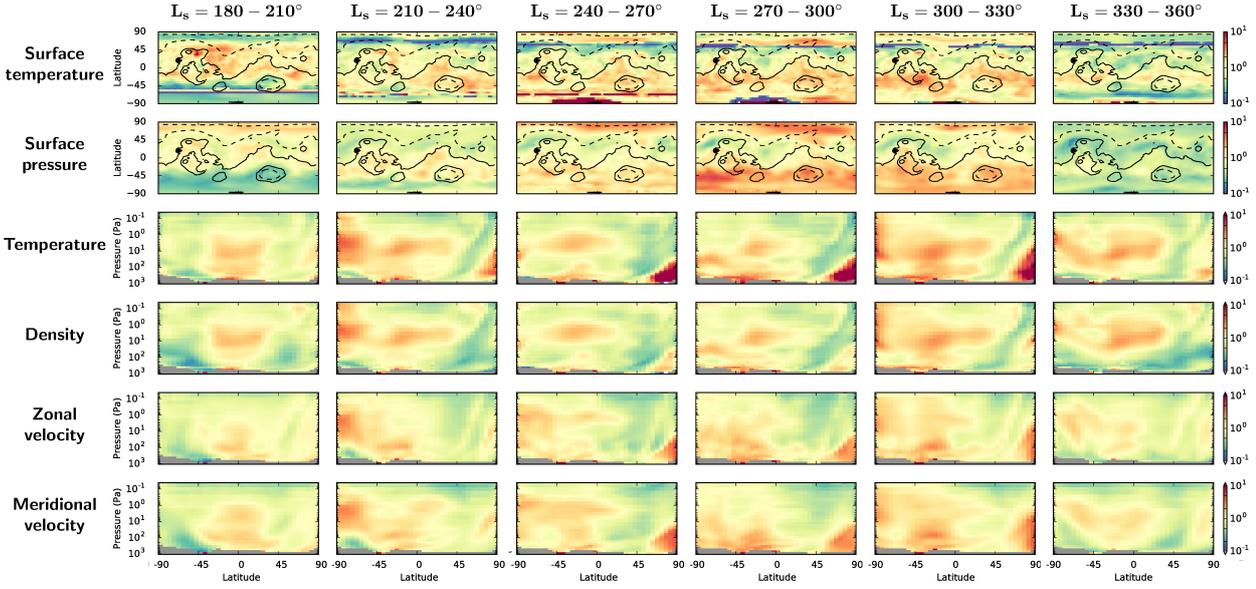


Figure 13: As Fig. 12, but for day-to-day variability. Where the ratio is greater than 1, the reanalysis variability is larger than the MCD variability over that period.

625 reanalysis. The largest differences were near the edge of the polar icecap. The latitude
 626 at the edge of the polar icecap has a large day-to-day variability during a given month,
 627 as the edge of the polar cap moves over time, and so the day-to-day variability ampli-
 628 tude will transition from small (with ice) to large (without ice). This is present in both
 629 the reanalysis and MCD day-to-day variability (Fig. S4). This also means the monthly
 630 mean difference between the MCD and the reanalysis is sensitive to the position of the
 631 edge of the polar cap, and so we see large differences between the reanalysis and the MCD
 632 in that region. There is a strong warm bias at southern polar latitudes between $L_s =$
 633 $240^\circ - 330^\circ$, due to the permanent CO_2 polar ice cap, which is present in the MCD but
 634 is not simulated in the version of the model used in the reanalysis.

635 Differences between the MCD and reanalysis surface pressure vary considerably with
 636 season. The surface pressure is persistently higher in the reanalysis in the summer hemi-
 637 sphere, and lower in the reanalysis in the winter hemisphere. There are also rings in the
 638 difference maps around regions with the most extreme elevation changes on the planet,
 639 particularly Olympus Mons, Elysium Mons, and Hellas (Fig. 12). The most likely rea-
 640 son is that the UK version of the Mars GCM (reanalysis) uses a spectral dynamical core,
 641 while the LMD Mars GCM (MCD) uses a grid point dynamical core. Quantities sensi-
 642 tive to surface elevation, such as surface pressure, will have large differences purely as
 643 a result of the topography being represented differently in the two models. The rings them-
 644 selves are characteristic of the Gibbs phenomena that occur when a step function is spec-
 645 trally decomposed, and are therefore likely to be spurious.

646 In the equatorial region the atmospheric temperature reanalysis is typically cooler
 647 than the MCD close to the surface, warmer around 100 Pa, cooler between 1 and 10 Pa,
 648 and warmer above 1 Pa. This pattern is repeated in most months. Both poles are typi-
 649 cally warmer in the reanalysis than in the MCD, at least in the lower atmosphere, with
 650 a warm “tongue” in the difference maps extending into the stratospheric polar region
 651 in the winter hemisphere. There is significantly more day-to-day variability in the re-
 652 analysis near the surface at the winter pole.

Density differences are generally small (up to 0.2 in $\log_{10} \rho$), but with some patterns. At low latitudes the density is generally lower in the reanalysis below 10 Pa, and higher in the reanalysis above this level. In the polar regions the density is nearly always lower in the reanalysis in the winter hemisphere. The day-to-day variability is lower in the reanalysis where the “warm tongue” appears in the air temperature maps. The horizontal striations in Fig. 12 are likely to be artifacts: there are different vertical grids in the reanalysis and MCD, which are interpolated to a common pressure grid for comparison, and density covers several orders of magnitude, so differences in interpolation will be magnified.

Zonal velocities in the reanalysis are, in general, more westward than in the MCD. This means that the eastward mid-latitude jets, the most prominent features of the monthly means (Fig. S12), are weaker in the reanalysis. These differences can be quite large – up to 50 m s^{-1} in magnitude. At low altitudes near the equator, the zonal flow is more eastward in the reanalysis than in the MCD. The day-to-day variability is generally lower in the reanalysis in the winter hemisphere, and larger in the reanalysis in the summer hemisphere (Fig. 13), with the exception of the winter pole near the surface, which has a high day-to-day variability in the reanalysis.

In general the meridional velocity reanalysis has a stronger upper-level equatorial and midlatitude meridional circulation than does the MCD: flow away from the sub-solar point is strengthened in the reanalysis between 0.1–1 Pa during the dusty season, compared with the MCD. Like the zonal velocity, typically the day-to-day variability is generally lower in the reanalysis in the winter hemisphere, and larger in the reanalysis in the summer hemisphere.

The only exception to the similar results for MY28 and MY29 was $L_s = 270 - 300^\circ$, which contains the build up to and peak of the MY28 global dust storm. Figure 14 shows the differences between the reanalysis and MCD during this period for both years. The main difference between the two years is that during MY28 the day-to-day variability is significantly larger in the reanalysis than in the MCD, when compared with the corresponding period during MY29. This applies to the surface temperature (particularly near the equator), and atmospheric temperatures, density, and both zonal and meridional velocities. Throughout the reanalysis sequence the day-to-day variability in the reanalysis is typically 1–2 times that in the MCD. This is likely due to the greater importance of shorter timescales in the reanalysis than in the simulations used to generate the MCD, such as a 15 minute timescale in the temperature field, as that is the interval between successive calls to the data assimilation procedure, and this may be expected to stimulate variability on timescales shorter than a day. However, during the MY28 global dust storm period this difference is amplified. During MY28 there is a clear warm anomaly (and corresponding low-density anomaly) between 1 and 100 Pa in the reanalysis compared with the MCD, compared with the corresponding period during MY29. A second major difference is that the monthly mean meridional velocity (i.e. the cross-equatorial flow) between 0.1–1.0 Pa is stronger by almost $10\text{--}20 \text{ m s}^{-1}$ in the reanalysis compared with the MCD during the MY28 global dust storm, while during the corresponding period in MY29 the difference is $\pm 5 \text{ m s}^{-1}$.

8 Conclusions

The data assimilation system integral to the UK-LMD Mars GCM described by S. R. Lewis et al. (2007) has been updated. That work assimilated temperature and CIDO, prescribing a vertical dust distribution using an empirical function of height (Conrath, 1975). The new scheme adds activated dust lifting, transport, and deposition schemes, assuming a single dust particle size. It also assimilates vertically resolved dust profiles via LIDO, either instead of CIDO or in addition to it. This update has been prompted by the acquisition of vertically-resolved dust profiles by MCS on board MRO (McCleese et al., 2010).

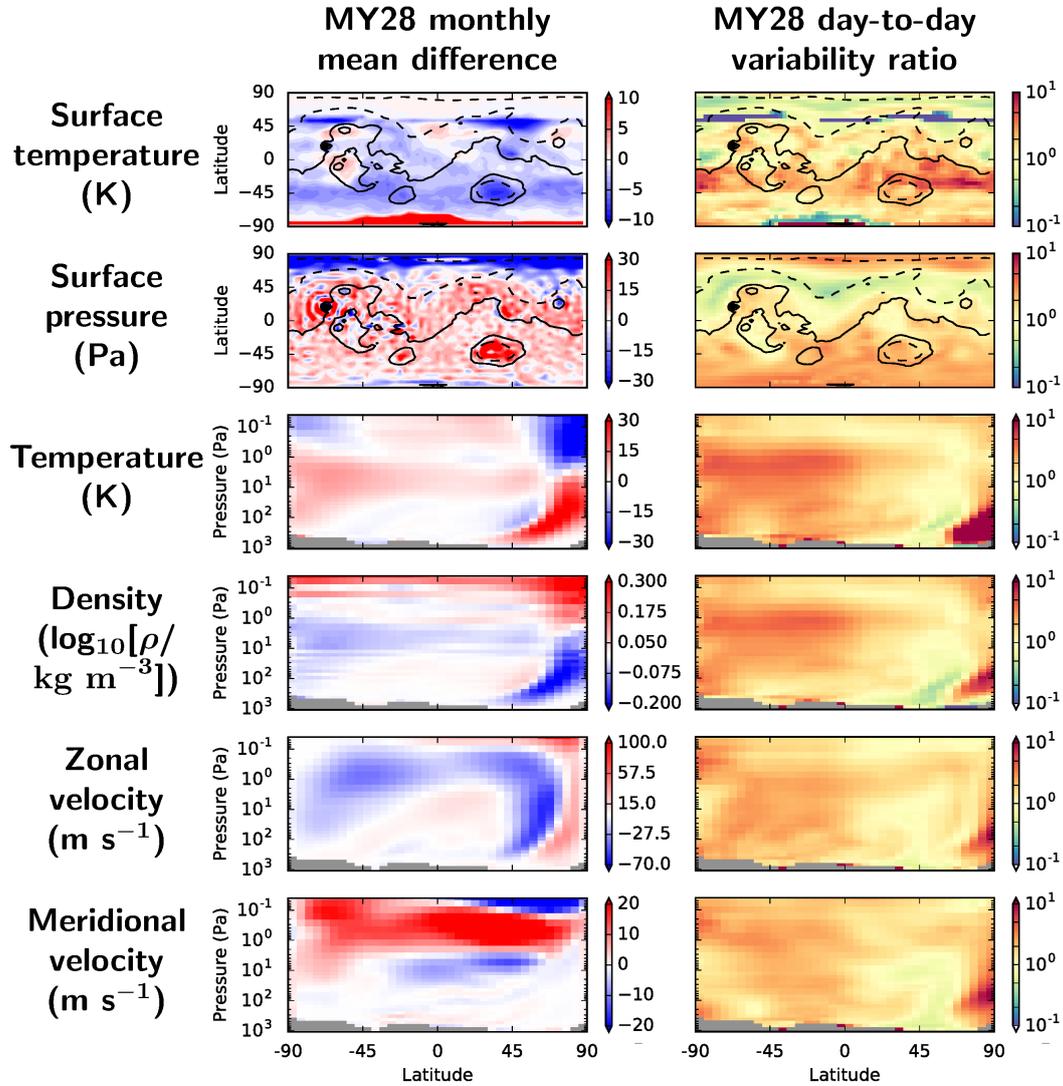


Figure 14: Reanalysis vs. MCD for the MY28 global dust storm period $L_s = 270 - 300^\circ$. Both the difference in monthly means and the ratio between day-to-day variability are shown. The corresponding period in MY29 is shown (in the same column format) in Figs 12 and 13. The units only apply to the mean differences.

704 When CIDO is assimilated by itself, the assimilation can reproduce the observed
 705 interannual variability of the dust horizontal spatial distribution, including the gener-
 706 ation and dissipation of the MY28 GDS (Fig. 4, red line). This inter-annual variability
 707 cannot be reproduced by our free-running model, although some degree of tuning can
 708 be done to reproduce the “quiet” ($L_s = 0^\circ - 180^\circ$) and “dusty” ($L_s = 180^\circ - 360^\circ$)
 709 periods during a single year. Even without assimilating any vertical information, the CIDO
 710 assimilation reduces systematic errors in the model’s estimate of the dust vertical dis-
 711 tribution. However, it misses detached dust layers that form during northern spring and
 712 summer (Fig. 5b), which have been a challenge for Mars GCMs to reproduce.

713 Conversely, when LIDO is assimilated by itself the model can reproduce some fea-
 714 tures of the detached dust layers (Fig. 5b, 4th panel), with reasonable interannual vari-
 715 ability, which the free-standing model is unable to reproduce. Rafkin (2012) discussed
 716 the difficulty of producing this detached dust layer in model simulations, especially in
 717 a relatively coarse resolution GCM. Similar detached dust layers were reproduced in mesoscale
 718 model simulations of “rocket dust storms” by Spiga et al. (2013) and have since been pa-
 719 rameterised in a coarse resolution GCM (C. Wang et al., 2018) with some success. Nev-
 720 ertheless, being able to reproduce them in a reanalysis provides a valuable alternative
 721 means of investigating their observed characteristics and impact on Mars’ atmospheric
 722 circulation. However, the limb-viewing MCS does not continuously observe the lowest
 723 part of the atmosphere, where the dust concentration is generally highest except where
 724 there are detached layers. Consequently assimilating LIDO only does not reproduce the
 725 observed global average τ_{ref} as well as when CIDO is assimilated by itself (see Fig. 4).

726 Once combined together, the joint assimilation of CIDO and LIDO benefits from
 727 information about the total column dust opacity and horizontal distribution from the
 728 THEMIS dataset, and the vertical distribution from the MCS dataset. The evolution of
 729 the MY28 GDS is tracked well, and some features of the detached dust layers are also
 730 reproduced well.

731 The joint assimilation of CIDO and LIDO is a powerful tool that helps us to re-
 732 construct the Martian climate as well as individual dust events. For example, it is dif-
 733 ficult to retrieve a complete vertical dust distribution from MCS measurements during
 734 the MY28 GDS (Fig. 1b). Using the joint assimilation, CIDO provides information to
 735 constrain the model where vertical profiles are sparse, so assimilation can map the four-
 736 dimensional dust distribution during the MY28 GDS in the absence of complete obser-
 737 vations.

738 The reanalysis was validated against out-of-sample THEMIS and MCS dust ob-
 739 servations, as well as upward-looking MER Pancam observations at 880 nm. The reanal-
 740 ysis successfully reproduced the observed interannual and intraseasonal variability in the
 741 original THEMIS data, and generally improved the representation of the dust vertical
 742 distribution compared to free-running simulations. In general, the free-running model
 743 tends to underestimate τ_{ref} , particularly during the dusty season, which can be signifi-
 744 cant during major dust storm events. Hence assimilating dust observations serves to greatly
 745 reduce the model uncertainty during such events. Although free-running simulations were
 746 able to simulate the pattern of dust loading during the quiet season, they failed to sim-
 747 ulate the vertical dust distribution. This is consistent with the general observation that
 748 dust accumulates close to the ground, below the base of typical MCS dust profiles.

749 At the Spirit rover location the reanalysis captured the dust variability and inten-
 750 sity during both quiet and dusty seasons (Fig. 11a). At the Opportunity rover location,
 751 the reanalysis captured the dust variability (Fig. 11b) and improved the pattern of dust
 752 loading over the free-running simulation during the dusty season. However, there were
 753 persistent underestimates of τ_{ref} by the reanalysis, particularly during northern spring
 754 and summer. This is likely to be due to a systematic disagreement between THEMIS
 755 and Opportunity data (Montabone et al., 2015).

756 A systematic comparison between the reanalysis and the MCD during MY28 and
 757 MY29 exhibited a number of trends, although agreement between the reanalysis and the
 758 MCD was generally good. Day-to-day variabilities were typically larger in the reanal-
 759 ysis by a small factor (1-2 times), likely associated with the stimulation of fluctuations
 760 on timescales comparable with the interval between successive calls to the data assim-
 761 ilation procedure (15 minutes). We also found some uncertainty in the position of the
 762 edge of the polar caps in the MCD, since the main differences in surface temperature arise
 763 along those latitudes. The reanalysis also showed significantly warmer poles than in the
 764 MCD at most times of year, along with weaker mid-latitude jets, and a stronger cross-
 765 equatorial flow during the MY28 global dust storm.

766 The combined CIDO-LIDO dust reanalysis significantly improves the estimation
 767 of Martian horizontal and vertical dust distributions over a free-running model and over
 768 CIDO or LIDO alone. The reanalysis provides a solution generally consistent with the
 769 available Martian dust observations. Assimilation has considerable potential as a tool
 770 for studying individual dust lifting events and for mapping Mars' three-dimensional dust
 771 distribution over time. Elsewhere, we will report on two case studies using the scheme,
 772 investigating a southward-moving regional dust storm during MY29, and the global dust
 773 storm during MY28.

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